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Measuring the stress field around an evolving crack in tensile deformed Mg AZ31 using 3DXRD

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Abstract
The stress field around a notch in a coarse-grained Mg AZ31 sample has been measured under tensile load by using the individual grains as probes in an in situ high-energy synchrotron diffraction experiment. The experimental setup, a variant of Three-Dimensional X-Ray Diffraction microscopy, allows the position, orientation and full stress tensor of each illuminated grain to be determined, and hence enables the study of evolving stress fields in coarse grained materials with a spatial resolution equal to the grain size. Grain resolved information like this is vital for understanding what happens when the traditional continuum mechanics approach breaks down and fracture is governed by local heterogeneities (e.g. phase or stress differences) between grains. As a first approximation the obtained results were averaged through the thickness of the sample and compared with an elastic-plastic continuum finite element simulation. It was found that a full three-dimensional simulation was required to account for the measured transition from the overall plane stress case away from the notch to the essentially plane strain case observed near the notch tip. The measured and simulated stress contours were shown to be in good agreement except at the highest applied load where a stress relaxation at the notch tip was observed in the experimental data. This stress relaxation is attributed to the initiation and propagation of a crack. Finally it was demonstrated that the measured lattice rotations could be used as a qualitative measure of the shape and extent of the plastic deformation zone.

1. Introduction
Within the last decade the stress fields around different types of cracks have been studied using a variety of in situ diffraction techniques. To get sufficient penetration power either neutrons [1,2] or high-energy synchrotron X-rays [3-7] have been employed. The latter case allows the unique combination of imaging and diffraction, also known as 3D crack tip microscopy [8], to probe the local conditions of the crack tip region. Here it is possible to combine tomographic evidence of crack growth, closure (in case of cyclic fatigue) or the spatial distribution of reinforcements and voids with the evolution in the stress field and plastic deformation zone measured by diffraction [9-11]. Common to these diffraction experiments is that the sample is scanned with respect to the beam, and that the dimensions of the beam determine the spatial resolution of the stress map. The analysis procedure is in all cases a generalization of the classical powder diffraction technique [12], where the resulting local strains and stresses are obtained as averages over the grains in the gauge volume. In order to obtain
diffraction data of sufficient quality the gauge volume must comprise at least of order 1000 grains. Consequently this procedure is limited to relatively fine-grained materials [13].

For more coarse-grained materials 3-Dimensional X-Ray Diffraction (3DXRD) microscopy is a candidate technique [14]. With 3DXRD up to 1000 grains are illuminated simultaneously and the diffracted signals from these grains are monitored on one or more detectors, while rotating the sample around an axis perpendicular to the beam. In the variant called Diffraction Contrast Tomography (DCT) only one near field-detector is used [15-17]. DCT has been shown to be a powerful tool for studies of short cracks, as space-filling 3D maps of grains and grain orientations can be provided and directly compared to 3D movies of the crack front as it penetrates through the sample [18,19]. However, it is difficult from DCT to derive strain and stress information with sufficient accuracy, and only samples deformed less than 1 % can be studied.

In the following we focus on a different variant of 3DXRD, where only a far-field detector is used. Based on such a set-up the centre-of-mass positions, relative volumes, mean orientations and full stress tensors for each grain within the illuminated volume can be monitored (but not the exact 3D grain morphologies). In previous of these so-called grain centre mapping experiments the number of mapped grains were rather limited [20-24], but recent progress has made it possible to map the stress field in a representative volume within the bulk of a polycrystalline sample by using the individual grains as probes [25-27]. Notably, the spatial resolution of the stress map is in this case determined by the grain size and the detector resolution (in terms of pixel size and spatial distortions) rather than by the dimensions of the incoming X-ray beam.

In this paper we present a far-field type 3DXRD study of the individual grains around a notch in a 0.8 mm thick coarse-grained specimen of the magnesium alloy AZ31. We determine the evolution of the stress field and the plastic zone (quantified by the rotations of the individual grains) at the notch tip during in situ tensile loading. The spatial resolution is 30 µm. The 3DXRD stress mapping is complemented by phase contrast tomography to visualize the crack growth. To our knowledge this is the first example of a stress map around a crack tip in a coarse grained material where both the sample
size and the spatial resolution are of sufficient quality to be technologically relevant, and in fact are similar to the very best strain scanning experiments on fine-grained materials [5].

2. Experimental

2.1. Sample preparation and characterisation

As sample material the Mg alloy AZ31 was chosen. A continuously cast billet was homogenized for 12 h at a temperature of $T = 350 \, ^\circ\text{C}$. Then the billet was extruded by indirect extrusion with a billet temperature of $T_B = 350 \, ^\circ\text{C}$ using a press ratio of $R = 25$ and a ramp speed of $v_R = 0.5 \, \text{mm s}^{-1}$ resulting in an extrusion profile with a diameter of 25 mm. From this rod flat samples with 0.8 mm thickness were machined with the longitudinal axis in the extrusion direction. After machining they were ground and polished. Subsequently a notch was applied by laser beam cutting. The final geometry of the sample with a notch is shown in Fig. 1(a).

![Fig. 1 Geometry of the sample (a) with distances in mm, and sketch of the mesh in the notch region (b) used in the elastic-plastic continuum finite element model.](image)

Details of sample material characterisation are given elsewhere [28]. For the purpose of the present study it is important to mention that the extruded profile shows a $(10\bar{1}0)\langle11\bar{2}0\rangle$ double fiber texture,
and that the microstructure of the profile shows some elongated grains which stem from the cast structure and some very small dynamically recrystallised grains, whereas the majority of the grains are statically recrystallised grains that have grown up to a grain size of 50 µm. The macroscopic tensile stress-strain curve for the sample material is shown in Fig. 2 along with the stress-strain curve from the synchrotron experiment. The tensile yield strength was determined to 180±3 MPa, the ultimate tensile strength to 244±3 MPa and the fracture strain to 20±2 %.

![Stress-strain curve](image)

**Fig. 2** Stress-strain curves for continuous loading and for the actual experiment. Arrows mark the target loads for 3DXRD measurements. Above the yield point the 3DXRD measurements were started after roughly half of the relaxation.

### 2.2. Synchrotron experiment

The experiment was carried out at beamline 1-IDC at the Advanced Photon Source at Argonne National Laboratory using optics and a bandwidth of $\Delta E/E = 1.2 \times 10^{-3}$ [29]. The X-ray energy was calibrated to the Tb K-edge (~52 keV) and continuously monitored by measuring the transmission through a Tb foil to enable corrections for temporal variations during the data acquisition. Two detectors were used alternately:

- For 3DXRD a MAR165 (with an active area of 165×165 mm$^2$ and a pixel size of 78.64×78.64 µm$^2$) was placed at a sample-to-detector distance of 41 cm to enable the detection of the first 10 full Debye-Scherrer rings of Mg.
- For tomography a high-resolution detector with 2048×2048 pixels of 1.4×1.4 µm$^2$ was placed on the same rail as the 3DXRD detector.

The sample was mounted in a tension frame with a 445 N (100 lbs) load cell to monitor the macroscopic stress state. A resistive strain gauge was glued onto the sample face to monitor the
macroscopic strain. The tension frame was driven in position control, and every time a target load for stress mapping was reached (5 MPa, 65 MPa, 110 MPa, 170 MPa, 190 MPa, 205 MPa and unload to 10 MPa – see Fig. 2 for the actual stress-strain data where arrows mark the target loads), the elongation was paused for around 12 h. During this time the sample was realigned relative to the X-ray beam, the stress states of the individual grains in the volume of interest were mapped by 3DXRD and the evolution of the crack was monitored by a tomography experiment. The 3DXRD experiments were started after roughly half of the observed stress relaxations above the yield point, and the stresses continued to relax during the measurements. The stress versus time dependence resulting from this relaxation is of the order 15 MPa over 12 h at the highest load, thus considerably less than the average grain resolved stress error bars (Table 2). Note that the loading history from the end of the 190 MPa synchrotron measurements to the beginning of the measurement at 205 MPa is missing. A dashed line indicates the approximate behaviour based on the maximum applied load and the corresponding relaxation observed at lower loads.

The 3DXRD experiment was carried out by illuminating the volume of interest with a planar beam, which was defined by a slit to originally 500 µm in the horizontal direction and 25 µm in the vertical direction. By translating the sample with respect to the beam, the sample was probed in layers perpendicular to the tensile axis. To probe all grains in the illuminated volume, and not just those that happen to fulfil the Bragg condition, the sample was rotated around the tensile axis and diffraction images were acquired in the angular range of $[-150^\circ ; -30^\circ] \cup [30^\circ ; 150^\circ]$ in steps of 0.5° for each layer. As the load and hence the amount of plastic deformation around the notch increased, the diffraction spots smeared out. It was therefore necessary to slit down the beam and increase the exposure time per diffraction image in order to get non-overlapping diffraction spots with a good signal-to-noise ratio. The beam dimensions, number of mapped layers and exposure time per diffraction image are summarized in Table 1.

After each 3DXRD experiment a set of radiographs for tomographic reconstruction were collected by opening the beam to 2×2 mm², rotating the sample $[0^\circ ; 180^\circ]$ in steps of 0.5° and making 2.5 s/radiograph exposures on the high resolution detector. The radiographs were corrected for detector-specific noise and background and then fed into a filtered-back projection algorithm for reconstruction.
The reconstructed images were subsequently corrected for ring artefacts [30]. In Fig. 3 the edge-on radiographs measured at 5 MPa and 205 MPa, which give the best view of the crack propagation, and the tomographic reconstruction of the non-straight crack front at 205 MPa are shown. From these the notch opening along $z$ was measured to be 210 $\mu$m up until 110 MPa, 230 $\mu$m at 170 MPa, 235 $\mu$m at 190 MPa, and finally 265 $\mu$m at the maximum load of 205 MPa. At 190 MPa the progression of the crack front into the sample (towards $-x$, c.f. the coordinate system in Fig. 1) has just started. The estimated movement is 30 $\mu$m. At 205 MPa the length of the crack is estimated to be 100 $\mu$m, thus 2-3 grain diameters.

**Fig. 3** Edge-on radiographs measured at (a) 5 MPa and (b) 205 MPa showing the crack progression in mm along with the tomographic reconstruction of the irregular crack front at 205 MPa (c).

<table>
<thead>
<tr>
<th>Load [MPa]</th>
<th>Beam width [\mu m]</th>
<th>Beam height [\mu m]</th>
<th>Number of layers</th>
<th>Exposure time [s / image]</th>
</tr>
</thead>
<tbody>
<tr>
<td>5 (d$_0$ reference)</td>
<td>500</td>
<td>25</td>
<td>15</td>
<td>0.5</td>
</tr>
<tr>
<td>65</td>
<td>500</td>
<td>25</td>
<td>15</td>
<td>0.5</td>
</tr>
<tr>
<td>110</td>
<td>500</td>
<td>25</td>
<td>15</td>
<td>0.5</td>
</tr>
<tr>
<td>170</td>
<td>400</td>
<td>20</td>
<td>17</td>
<td>0.5</td>
</tr>
<tr>
<td>190</td>
<td>200</td>
<td>10</td>
<td>15</td>
<td>1.5</td>
</tr>
<tr>
<td>205</td>
<td>200</td>
<td>10</td>
<td>15</td>
<td>1.8</td>
</tr>
<tr>
<td>10 (unload)</td>
<td>200</td>
<td>10</td>
<td>15</td>
<td>1.8</td>
</tr>
</tbody>
</table>
2.3. Analysis of diffraction data

First, the positions of the diffraction peaks were located. This step was performed using the 2D peaksearch module from the FABLE suite of programs [31] for the loads 5 MPa, 65 MPa, 110 MPa and 170 MPa. At the higher loads the overlaps due to plastic deformation could only be discerned by employing the more sophisticated and time consuming 3D peaksearch routine DIGIgrain [32]. The program GrainSpotter [33] was then used to index the grains, i.e. assign the diffraction spots to oriented grains, in each illuminated layer. The refinement of the centre-of-mass positions, orientations and full elastic strain tensors of the indexed grains was performed as a simultaneous 12-parameter-per-grain fit to the assigned reflections using the FitAllB module of FABLE [25]. The average number of reflections per grain after the built-in outlier rejections and the convergence of the refinement can be seen in Table 2. The relative grain volumes were estimated based on the intensities of the assigned reflections. The grain resolved stress tensors were calculated from the strain tensors using the following values for the stiffness tensor: \(C_{11}=58.5\) GPa, \(C_{12}=25.0\) GPa, \(C_{13}=20.8\) GPa, \(C_{33}=61.0\) GPa, and \(C_{44}=16.6\) GPa [34].

<table>
<thead>
<tr>
<th>Load [MPa]</th>
<th>Number of grains</th>
<th>spots/grain</th>
<th>3D position [µm]</th>
<th>Relative volume [%]</th>
<th>Orientation [°]</th>
<th>Stress components</th>
</tr>
</thead>
<tbody>
<tr>
<td>5</td>
<td>3268</td>
<td>85</td>
<td>18</td>
<td>19</td>
<td>0.09</td>
<td>11</td>
</tr>
<tr>
<td>65</td>
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<td>85</td>
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<td>110</td>
<td>2383</td>
<td>85</td>
<td>17</td>
<td>19</td>
<td>0.08</td>
<td>10</td>
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<tr>
<td>170</td>
<td>1750</td>
<td>75</td>
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<td>20</td>
<td>0.07</td>
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<tr>
<td>190</td>
<td>207</td>
<td>55</td>
<td>69</td>
<td>14</td>
<td>0.21</td>
<td>45</td>
</tr>
<tr>
<td>205</td>
<td>99</td>
<td>40</td>
<td>104</td>
<td>20</td>
<td>0.31</td>
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<td>40</td>
<td>112</td>
<td>28</td>
<td>0.34</td>
<td>73</td>
</tr>
</tbody>
</table>

The quality of the grain resolved results obviously depends on the calibration of the global parameters describing the experimental setup, i.e. the sample-to-detector distance, the position of the beam centre on the detector, the tilts of the detector face and of the sample rotation stage, the wavelength of the incident beam and the lattice constants of the undeformed material. All of these parameters were refined using FitAllB on the indexed grains from the data set collected at a macroscopic load of 5 MPa.
assuming no residual stress as described previously [25]. 5 MPa was chosen as the reference point to ensure that the sample was properly gripped. Besides a drift in incident energy of the order $\Delta E/E = 10^{-4}$ over the 4 days it took to do the experiment, the beam centre on the detector was also found to drift during the entire experiment, presumably because the diffraction detector was moved out of the beam to make space for the tomography detector occasionally. Both of these effects were corrected for.

Since the average grain size (30 µm) is larger than the height of the planar beam (25-10 µm) many grains were illuminated in more than one layer. This 2D grain information was collected into 3D by stacking the illuminated layers and matching up the grains with similar positions and orientations taking into account the relative volumes. At the 5 MPa reference load 3268 such 3D grains were identified around the initial notch. This corresponds to around 90% of the grains in the illuminated volume judging from the number of assigned to observed reflections when taking into account the beam size relative to the sample thickness. The number of 3D grains that can be followed at the subsequent loads can be seen in Table 2. Here it should be emphasised that between 170 and 190 MPa the beam dimensions were halved in both directions, thus the number of grains in the gauge volume goes down by a factor eight. Grains in the top and bottom layers were kept even though some of these will be grain fractions rather than entire grains c.f. Fig.4. The error bars on the centre-of-mass grain positions, volumes, orientations and strain tensors were estimated by FitAllB as described previously [25]. The average values of these are also summarized in Table 2.

In Fig. 4 the reconstructed grains in the volume of interest in front of the crack are displayed. The grain maps from different loads have been aligned by minimising the distance between grains found in both maps, and the initial notch position $(x,z)=(0,0)$ and opening angle have been marked on each map to illustrate the notch opening and crack propagation. The size of the spheres representing the individual grains reflects the relative grain volumes. The grains are colour coded according to how much the crystallographic lattice of each has rotated relative to the initial orientations found at the 5 MPa reference load. For small deformations like the present 1% the lattice rotation are roughly proportional to the uptake of plastic deformation in the grain, especially when taking into account the strong extrusion texture of the sample that minimises the orientation dependence of the lattice rotations [35,27].
2.4. Modelling

A continuum elastic-plastic finite element model of the central part of the test specimen (Fig. 1) has been built in the commercial finite element code Abaqus [36]. A full 3D model is used in order to account for the measured transition from the overall plane stress case away from the notch to the essentially plane strain case near the notch tip. Using the symmetry planes in the transverse direction through the notch and along the mid plane in the thickness direction, only one quarter of the central part with the dimension 8x4x0.4mm³ has been modelled using 8-noded brick elements. A uniform mesh was used in the thickness direction with eight elements meshing half of the thickness. As seen in Fig. 1(b), a structured concentrated mesh was used in the notch region, while the overall mesh was unstructured. The notch has through the thickness been modelled using a constant tip radius of 17 µm. The magnesium material is modelled using a continuum J2-flow theory with a hardening curve following the monotonic loading curve in Fig. 2. The hardening curve has been measured up to a strain
of 8 %, beyond which the hardening curve has been extended using a power-law fit. The elastic part has been modelled with $E = 40$ GPa and $\nu = 0.3$, while plastic deformation occurs for stresses exceeding the initial yield stress $\sigma_y = 180$ MPa. The finite element simulation does not include time dependency. This is why the load drop together with continuing elongation in the notched region observed during each of the approximately 12 hour long 3DXRD mapping is not captured. Instead the strain level corresponding to the midpoint of each strain range monitored during the 3DXRD mapping was used. This corresponds to the second half of each of the six load drops starting at the points indicated by arrows in Fig. 2. The stress and effective plastic strain contours, which are later compared with the measured values, have each been extracted relative to the undeformed reference state.

3. Results

3.1. Grain resolved stresses

The principal results of the investigation, the grain resolved stresses, show significant grain-to-grain variations as expected from the different orientations and thus yielding criteria of the individual grains. The differences in axial stress between neighbouring grains increase with the applied stress and are roughly three times the average error bars given in Table 2 far away from the notch. Before macroscopic yielding the differences increase roughly 10% for grains closer than 0.10 mm to the notch tip. At an applied load of 190 MPa a 50% increase is observed for grains closer than 0.05 mm to the tip, and at 205 MPa grains with a distance of 0.05-0.10 mm from the original position of the notch show stress differences to their neighbouring grains that are twice as large as for grains in other parts of the mapped volume. This trend is in agreement with the triaxiality of the macroscopic stresses at the crack front and the observed stress relaxation attributed to a crack propagation of the order 0.10 mm at the maximal applied load.

Grain resolved information like that measured in the present study is vital for understanding what happens when the traditional continuum mechanics approach breaks down and fracture is governed by local heterogeneities (e.g. phase or stress differences between grains). In the following section 2D projections of the measured grain resolved stresses are compared with results from an elastic-plastic continuum finite element model. Not surprisingly the correspondence decreases as the crack initiation approaches and the local grain resolved conditions at the very crack tip become increasingly important.
3.2. Stress fields

2D projections of the measured grain resolved stresses collapsing the sample thickness onto the $xz$-plane were made and compared with the simulated contours. The 2D projections were produced on a $10 \times 10 \ \mu m^2$ grid by Laguerre tessellation [37], thus for each grid point the contribution was taken from the grain for which the distance from grid point to the projected grain centre divided by the grain radius was smallest. At 65 MPa the volume fraction of contributing grains was 0.85, while this fraction was increased to 0.97 at 205 MPa where the measured volume was thinner in the $y$-direction perpendicular to the projection plane. After the tessellation a $3 \times 3$-pixel smoothing was applied. The 2D projections of selected measured stress tensor components are visualized by colour scaling. The limits are chosen such as to best display the evolution over the whole load cycle. At higher loads, local extreme stress values are outside of the chosen limits. The corresponding simulated contours are shown as white lines.

**Fig. 5** Measured 2D projections of the axial stress component ($\sigma_{33}$ in colour) overlaid by the contour levels from the simulation.
The 2D projection of the stress along the tensile z-axis ($\sigma_{33}$) is given in Fig. 5, while Fig. 6 shows the stress in the direction of crack growth along the x-axis ($\sigma_{11}$). For the four load levels before the onset of crack growth, the measured normal and axial stress fields in Fig. 5 and Fig. 6 have been overlaid with the stress contours obtained from the elastic-plastic continuum finite element model. Even though the model was built for a homogeneous material, while the measured stress contours are extracted from a polycrystalline material measuring the deformation state in each of the individual grains, a reasonable match can be observed. In particular the simulations support the observation of normal stresses ($\sigma_{11}$) building up in front of the notch, while axial stresses ($\sigma_{33}$) build up in a butterfly-shaped contour around the tip. In the plastic regime the simulated stresses tend to be higher than the measured stresses in accordance with the stress relaxation observed during the diffraction experiment which results in reduced stresses relative to the hardening curve used in the simulation (Fig. 2). In the elastic regime the measured stresses vary more or less continuously, whereas in the plastic regime the fluctuations increase significantly. This is consistent with the increase in the estimated errors reported in Table 2 and the increased role played by the local grain resolved conditions as the crack initiates. Furthermore, the y-dimension of the mapped volume is roughly equal to the x-dimension, thus at applied loads of 190 MPa and beyond the grain statistics leading to the projected stress contours on the xz-plane (Fig. 5 and Fig. 6) are poorer.

At a macroscopic load of 190 MPa all the grains in the mapped volume have yielded and show large stresses in the x and z-directions. However, at 205 MPa some of the grains immediately in front of the notch tip display much lower, sometimes even compressive, stresses. This stress relaxation and accompanying shift in the position of the axial stress ($\sigma_{33}$) maximum to be further ahead of the notch tip is a result of the crack propagation. Upon unloading the stress reversal zone with compressive stresses can be seen immediately in front of the crack where the tensile stresses were maximum at the maximum load.

While the simulations predict the position of maximum axial stress ($\sigma_{33}$) to move ahead of the notch as the load is increased in the elastic regime, this shift cannot be measured by 3DXRD. This has two different plausible explanations. Firstly, at the lower loads the expected shift in stress maximum away from the notch tip is smaller than the average grain diameter, and the measured stress maximum is
therefore likely to be smeared out due to the presence of intra-granular stress gradients. Sub-grain stress gradients as well as domain size effects are reflected in the peak broadening of the diffraction peaks along 2θ. Separating out the effects of domain size, the present experiment shows significantly larger peak broadening for the grains in the vicinity of the notch tip in good accordance with the expected intra-granular stress gradients within these grains. Secondly, at more moderate loads the grains of maximum stress have deformed plastically, and as discussed later the method tends to miss a number of the plastically deformed grains, see e.g. Fig. 4 (170 MPa). This could therefore very well be the reason that the expected shift in the axial stress (σ_{33}) maximum in front of the notch tip does not show up in the diffraction measurement before macroscopic yielding. After macroscopic yielding a shift is indeed observed and attributed to crack propagation.

**Fig. 6** Measured 2D projections of the normal stress component (σ_{11} in colour) overlaid by the contour levels from the simulation.
In Fig. 7(a) the experimentally extracted out-of-plane normal stress component ($\sigma_{22}$) is plotted together with the corresponding model prediction. Both the prediction and the measurement show significant variations near the notch tip indicating divergence from the plane stress case in this region. This observation implies that a full three-dimensional simulation is required to account for the measured stresses. The corresponding comparison of the in-the-plane shear stress component ($\sigma_{13}$) is shown in Fig. 7(b). Especially an agreement of the positive and negative region above and below the notch can be observed.

![Fig. 7 2D projections of the (a) out-of-plane normal ($\sigma_{22}$) and (b) in-plane shear ($\sigma_{13}$) stress components measured at 170 MPa (in colour) overlaid with the simulated contour levels.](image)

3.3. Lattice rotations

Whereas diffraction peak broadening is related to plastic deformation on the sub-grain scale, lattice rotations reflect the plastic deformation on the grain scale. The shape and extent of the plastically deformed zone can thus be estimated from the data in Fig. 4. In Fig. 8 the 2D projections of the lattice rotations from Fig. 4 are shown overlaid with the contours of the effective plastic strain ($\varepsilon^p = \sqrt{\frac{3\varepsilon_{ij}^P \varepsilon_{ij}^P}{2}}$, the plastic strain analogy to the von Mises stress) from the simulation. Here it can be seen that there is a good correlation in both the shape and the extent of the plastic zone between the simulation and the measurement based on lattice rotations. However, in the part of the sample immediately above and below the notch the observed lattice rotations can of course not be directly used as a measure of the
plastic deformation. In this region a correction for the effect of notch opening, which makes the grains move apart and rotate, must be applied.

Fig. 8 2D projections of experimental lattice rotations in Fig. 4 overlaid with the contour levels of the effective plastic strain (values given in %) from the simulations.

4. Discussion

It is well known that the material microstructure often plays a pivotal role in dictating the modes of fracture and failure, and the macroscopic response of real materials. The grain morphology, lattice orientation, elastic modulus, and the toughness of the individual grains and interfaces are key parameters that control the failure mechanisms in coarse-grained materials. These effects could also be directly observed in the present data set where the crack front was seen not to move along a straight line, but rather in irregular ways, c.f. Fig. 3. Based on this insight, concepts such as grain boundary engineering have been developed to improve the fracture resistance of polycrystalline materials [38].

A number of 2D or 3D modelling approaches for describing fracture in polycrystals have been demonstrated. As examples:
- Yang and co-workers [39] used a lattice spring model, where a spring fails if the stored elastic energy in the spring exceeded a critical value.
- Zhai and Zhou [40] proposed a micromechanical FEM model in which the cohesive surface formulation of Xu and Needleman [41] was used.
Sukumar and co-workers [42] described brittle fracture with an extended finite element method, which uses the notion of partition of unity implying that re-meshing is not required. Finite element modelling has also been used in a number of cases to determine the stress field in both mono-phase [43,44] and multi-phase polycrystalline specimens [45] in order to predict crack growth. The current work is an attempt to provide relevant tools for validating such models.

In our view, the above results demonstrate the potential for stress mapping of materials where the grain size is comparable to the relevant crack dimensions. The stress resolution is adequate – at least at applied loads up to the yield point – and the sample dimensions are compatible with characterising short cracks. Furthermore, the size of the volume characterised provides reasonable grain statistics for averaging stresses in distances from the crack tip exceeding ~50 µm. The number of grains at the very crack tip is limited, in the present case to around 20 grains. This number is evidently insufficient for a detailed quantitative comparison with a continuum model, but it allows for the observation of symmetries, inhomogeneity and for an approximate comparison of properties such as the “size of the plastic zone”.

A promising route for future studies would be to include a near-field detector in the set-up, either in a DCT configuration or in a classical 3DXRD configuration. Similar to the work in [18,19] a space-filling 3D grain map can then be provided and directly compared to 3D movies of the crack front as it penetrates through the sample. In combination with the technique introduced in this article the result would be a very comprehensive characterisation of the influence of the grain scale heterogeneity on the fracture dynamics. Other relevant generalisations include:

- Simultaneous measurements of the plastic strain by tracking of the centre-of-mass positions of the individual grains. Alternatively by tracking of embedded marker particles, such as second phase particles, by tomography [46-51].
- Work on multiphase materials. The software introduced above can be run iteratively, one phase at the time to give the stress distribution in both matrix and individual reinforcing fibres as a coarse-grained analogy to e.g. work by the Manchester group [11,52].
- Combined diffraction and tomography studies where the grain resolved information can be complemented by information about crack debris, closure (for cyclic fatigue), crack deflection at
reinforcements or void coalescence for a complete 3D crack tip microscopy [8].

- A more detailed characterisation of the grains at the very crack tip, e.g. in terms of dislocation content [53] or intra-grain stress distribution [54].

The main limitation of the technique relates to the azimuthal broadening of the diffraction spots with increasing plastic deformation. This is the reason that the fraction of initially identified grains that are observed also at 170 MPa and at 205 MPa are 85% and 47%, respectively. More specifically we may distinguish between two effects: an increasing amount of overlapping spots and the fact that the signal-to-noise ratio deteriorates when the diffraction spot is spread over many more pixels on the detector. The latter (and more frequent) case either implies that the diffraction spot completely vanishes after setting an intensity threshold – and hence no reflection is observed – or that the centre-of-mass determination of the diffraction spot is flawed. There are several remedies to this problem. The spot overlap can be reduced by reducing the beam size and/or by placing optical elements like a conical [55] or spiral slit [56] between the sample and the (far-field) detector. The conical or spiral slit may define a 3D gauge volume of say 20x20x200 µm³ within the sample. More generally the centre-of-mass approach of this paper – where each diffraction spot is associated only with a centre-of-mass position – may be replaced by a description in terms of a high-resolution orientation distribution function [57]. This would not only enable tracking of grains to larger strains, but may also be a step in direction of measuring intra-grain stresses.

5. Conclusions
The present study represents an important step towards understanding the mechanisms governing crack growth in coarse-grained materials, namely the application of the recently matured far-field 3DXRD technique to measure the grain resolved stresses around a notch in a coarse-grained specimen during in situ deformation. Tomography and elastic-plastic continuum finite element simulations complemented the diffraction measurements. Substantial stress variations between neighbouring grains, an irregular crack front and local deviations from the continuum model after crack initiation was observed, illustrating the importance of taking the local grain information into account in a thorough mechanical description. The evolution of the stress field and plastic deformation zone (as described qualitatively in terms of lattice rotations of individual grains) was determined from the diffraction experiment with a
spatial resolution of 30 µm (the grain size). The measurement clearly showed a concentration of tensile stresses in the out-of-plane stress component near the notch tip indicating a divergence from the expected plane stress case in this region, which was consequently implemented in the simulations. The measured and simulated stress contours were shown to be in good agreement except at the highest applied load where a stress relaxation at the notch tip was observed in the experimental data. This stress relaxation is attributed to the initiation and propagation of a crack that was also observed with tomography. Future prospects in complementing the diffraction experiments by high-resolution tomography to study plastic flow or multiphase materials, by more localised studies of sub-grains at the very crack front and by crystal plasticity finite element simulations have been presented and discussed with the aim to obtain a more comprehensive understanding of crack growth in coarse-grained materials.

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