Two level undercut-profile substrate-based filamentary coated conductors produced using metal organic chemical vapor deposition

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Abstract—The two level undercut-profile substrate (2LUPS) has been introduced as a concept for subdividing rare-earth-Ba$_2$Cu$_3$O$_7$ coated conductors (CC) into narrow filaments that effectively reduces the AC losses and improves field stability for DC magnets. The 2LUPS consists of two levels of plateaus connected by a wall with an undercut-profile, which enables a physical separation of the superconducting layer between the plateaus without reducing the effective width of the superconducting layer. 

In this study we report for the first time the results of fabrication and characterization of a filamentary CC produced in an industrial setup by SuperPower Inc. using ion beam assisted deposition and metal organic chemical vapor deposition (IBAD-MOCVD) on a 2LUPS realized at the Technical University of Denmark (DTU), whereas previous studies discussed the fabrication using alternating beam assisted deposition and pulsed laser deposition (ABAD-PLD).

We also present Hall probe scanning measurements performed using a standard THEVA TAPESTAR™ XL machine that is routinely employed for industrial critical current characterization of long length CCs. From these results it is clear that additional analysis of the measured field profiles are required when characterizing filamentary 2UPS CC using a standard TAPESTAR™ setting. Using a model representation of the 2LUPS we calculated the expected magnetization response by means of finite element methods simulations and we find a good agreement with the experimentally observed magnetic profiles.

Index Terms—Multifilamentary superconductors, High-temperature superconductors, Magnetic variables measurement, Finite element analysis.

I. INTRODUCTION

HIGH-TEMPERATURE SUPERCONDUCTORS (HTS) are relevant to many scientific and technological applications [1]–[3] such as electric motors and generators, electric power distribution systems, and high-field electromagnets.

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Superconducting electromagnets are used in magnetic resonance imaging (MRI) and nuclear magnetic resonance (NMR) devices, particle accelerators, and many other scientific instruments and in general made using low temperature superconductors (LTS). HTS present several advantages over LTS which include higher irreversibility fields [4] and higher operation temperatures. Rare-Earth Barium-Copper-Oxides, such as YBa$_2$Cu$_3$O$_7$ and Gd$_{1.2}$Y$_{0.8}$Ba$_2$Cu$_3$O$_{7}$ (GYBCO), are a well-documented [1]–[4] class of HTS materials in the 2nd generation HTS wires, i.e. called coated conductors (CC). The width of these tapes is an important design parameter, which affects hysteretic losses in AC applications as well as the magnetic field drift in strong DC magnets [5]–[7]. These effects are caused by the screening current, which arises inside the superconductor tape to prevent the magnetic field from penetrating the material. This current and its variation over time are responsible for AC-losses and may also reduce the precision of the required magnetic field distribution, which is critical for many superconducting magnet applications. It has previously been shown that these detrimental effects can be mitigated by transposing the tape and subdividing the tape along its width into many decoupled filaments [7], [8]. Filamentation methods employed to achieve this subdivision include mechanical striation, laser striation and ink jet printing of filaments [9]–[12]. The challenge with these approaches comes from the reduction of the effective width of the superconducting (SC) layer that decreases the engineering critical current density, in addition to potential critical damage over the long length. Filamentization can potentially be achieved without reduction of the effective width of the tape using the two level undercut profile substrate (2LUPS) concept [13], [14] (see Fig. 1), consisting of two levels of plateaus in the metal substrate surface separated by a vertical displacement with an undercut profile.
The undercut profile can be realized by surface modification of the areas of the substrate corresponding to the lower plateaus by means of electrochemical etching. In this study we present, for the first time, experimental characterization of a 2LUPS CC produced by SuperPower Inc. using ion beam assisted deposition and metal organic chemical vapor deposition (IBAD-MOCVD) in a commercial system on a 2LUPS prepared at the Technical University of Denmark. No post-treatment was conducted after layer depositions. The performance of the 2LUPS CC was evaluated by means of Hall probe scanning using a Theva TAPESTAR™ XL machine that is routinely employed for analysis of long length samples at SuperPower Inc. and compared to a CC produced without surface profiles. The 2LUPS CC is evaluated in this system since filamentary long length samples would eventually also require such industrial characterization. However, data obtained using a TAPESTAR™ machine does not take into account geometrical effects and further interpretation is therefore required to evaluate the critical current. In order to correctly calculate the critical current of the filamentary 2LUPS CC, we compare the experimental data with computations performed using FEM analysis.

II. METHODS

A. Samples fabrication

Starting from a commercial Hastelloy C276 tape (L×W×T = 300 mm × 10 mm × 0.89 mm), wet electrochemical etching is employed in order to create the desired 2LUPS geometrical profile and to reduce the surface roughness [14]. Protective masking tape has been applied and cut in a reel-to-reel system after which 3 strips were peeled off. Lower plateaus were then formed during subsequent etching. The widths of each of the two lateral filaments are 2.8 ± 0.1 mm, while the five central plateaus has a width of 0.9 ± 0.1 mm. The reference sample was produced using the standard etching procedure [14]. Following the 2LUPS fabrication, the buffer and superconducting GYBCO layers are deposited on the substrate in a commercial production system at SuperPower Inc. by means of IBAD-MOCVD [15]. The total thickness of the buffer layers is around 0.2 μm and the superconducting Gd1.9Y1.6Ba2Cu3O7−δ layer has a thickness of 1.1 μm. Finally, a 0.7 μm thick protective silver layer is deposited on the GYBCO layer by means of sputtering.

B. Scanning electron microscopy

The cross-section of the 2LUPS sample has been analyzed using focused ion beam scanning electron microscopy (FIB-SEM) in a Carl Zeiss 1540 XB electron microscope. Fig. 2 shows the cross-section of the 2LUPS and the different layers of the CC. A gap is present in the GYBCO layer just below the edge of the upper plateau (marked by “A”) that verifies the desired physical decoupling between the superconducting filaments. It was observed that the undercut length \( L = 1.2 \pm 0.5 \, \mu m \), which is significantly greater, compared to a previously reported Hastelloy-based 2LUPS sample \( L=0.3 \, \mu m \) [14]. In the curved region of the lower plateau (see Fig.2) the GYBCO layer is observed to be both porous and irregular, and as such considered non-functional. In contrast, a very dense and homogeneous microstructure is observed on the regions of the plateaus not located in the immediate vicinity of the edge of the filaments, which is typical for CC production at SuperPower Inc. Measuring several cross-sections, it was observed that the extension of the damaged region, marked by “B”, is equal to 18 ± 2 μm, which is significantly smaller than that obtained using other filamentization techniques.

C. Hall probe scanning

The remanent magnetic field in different positions along the \( x \) direction of the 2LUPS and the reference samples were measured at 77 K using a TAPESTAR™ XL machine at SuperPower Inc. The apparatus employs an array of 21 Hall sensors spanning a width of 12 mm [16]. Fig.3(a) displays the field profile plotted as function of \( y \). The data corresponding to the 2LUPS and the reference sample are plotted as red and blue dots, respectively. The error bars shown in the figure correspond to the variation of the field profile along the length \( x \) of the tape. The solid lines of matching colors correspond to FEM calculations discussed in the next section. The vertical gray lines indicate the filaments edges of the 2LUPS tape.
The TAPESTAR™ software utilizes conventional inversion techniques to compute $I_C$ from the magnetic field it generates in the immediate vicinity of the surface of the tape [18]. The resulting plots of the critical current $I_C$ as function of the position along the length of the tape, i.e. $x$ direction, are shown in Fig. 4 as dashed lines. It is observed that the critical current of the 2LUPS is substantially lower than that of the reference sample. However, if the spatial resolution is too coarse or the distance of the Hall sensor from the tape is too large to resolve the small signal generated by the narrow filaments, one would expect a similar reduction in the calculated $I_C$, which does not correspond to the true value. It is clear that the conventional inversion requires special attention due to the filament geometry. In addition, the setup is not ideal for a complete characterization of filamentized samples and potentially in future studies one would modify the TAPESTAR™ machine. Taking into consideration the great effect of the distance on the obtained field profiles [19]–[23], we calculated a set of typical field profiles from analytical formulas [17], which can be seen in Fig. 5. The formula, reported in the reference, is derived from the two-dimensional Biot-Savart law applied to a mono-dimensional current distribution. The different lines shown in the figure correspond to different values of the distance $\Delta z$ between the Hall sensors and the SC layer of the tapes. The distance $\Delta z$ increases from 0.5 mm to 1.5 mm with steps of 0.5 mm, as shown in the figure. As can be noticed, at a distance of 1 mm from the SC layer the field oscillations generated by the narrow filaments of the central region of the tape are completely smoothed out and would not be resolved by the measurement. The black dashed line corresponds to the case where the narrow filaments are not working. As can be noticed, a significantly higher level in the center would have been expected if the filaments were not working. We performed additional investigations employing finite element method analysis. As discussed in the next section, the field profiles calculated by means of FEM simulation are more realistic than the ones provided by the analytical formula since the FEM simulation solves the partial differential equation governing the magnetic field distribution, while imposing the constraint on the time-dependence of the applied field and the relevant constitutive relations.

D. Finite Element Method analysis

The experimental data set has been compared to corresponding results computed by FEM analysis. The computations have been performed using the magnetic field formulation interface of the commercial software COMSOL Multiphysics. The model solves Maxwell’s equations and the relevant constitutive relation between the magnetic field $\mathbf{H}$ and magnetic flux density $\mathbf{B}$. For this study we assumed the following constitutive relation: $\mathbf{B} = \mu_0 \mathbf{H}$. A power law $\mathbf{J} - \mathbf{E}$ constitutive relation is assumed between the electrical field $\mathbf{E}$ and the electrical current density $\mathbf{J}$ [25], [28]. The relation is parametrized by the critical current density $J_c$:

$$E = \frac{E_c}{\|J/J_c\|} \mathbf{J} / \|\mathbf{J}\|$$

(1)

For large values of the exponent $n$, the highly non-linear $\mathbf{J} - \mathbf{E}$ relation practically constrains the current density in each point of space to be either zero or equal to $J_c$, thus reproduc-
ing the behavior described by Bean’s critical state model [26], [27]. The critical current density $J_C$ depends on the norm $B$ of the magnetic flux density. We assume the $J_C(B)$ dependence reported in Ref. [24], which has been measured from CCs produced at SuperPower Inc. using the same procedure as for the samples studied in this work and the $J_C(B)$ relation can be described by:

$$J_C(B) = J_{C0}/(1 + (B/B_0)^a)$$ \hspace{1cm} (2)

We used the method of least squares to fit this model function to the data reported in Ref. [24], and we obtained the following parameters: $J_{C0} = 3.52$ MA/cm$^2$, $B_0 = 3.54$ T, $a = 3.62$, assuming that the critical current density $J_C$ depends only on the magnitude of $B$. The simulations have been performed using a time-dependent two-dimensional model, which does not resolve the $x$ direction. This simplification is justified since both the samples (2LUPS and flat reference) have length greater than 20 cm, and width of 1 cm. The intensity of the magnetic field applied in the $z$ direction is ramped from 0 mT to 0.037 mT, simulating the TAPESTAR measurement. The geometrical characteristics of the reference and 2LUPS samples have been used for the simulation, where the latter also includes the exact width of each of the filaments. Since the actual distance between the SC layer and the Hall sensor is not known with high precision, we calculated the distance by fitting the FEM results to the experimental data. Employing this procedure we obtained the value of $\Delta z = 1.19$ mm for the reference sample and $\Delta z = 0.85$ mm for the 2LUPS sample, which are values of distance within the expected range.

The results of the FEM computations are shown in Fig. 3 as red and blue solid lines for the 2LUPS and reference sample, respectively. The dashed red line corresponds to the FEM simulation for a tape where the narrow filaments are not working (labelled as “off”). As can be seen, the match between the computed field profiles and the experimental data sets are very good. The relative error $\delta$ between the measured field profile, $B_{Exp}$, and the one calculated from the simulation, $B_{Sim}$, has been quantified using the normalized sum of the squared residuals, expressed by the following formula:

$$\delta = \frac{\sqrt{\sum_k (B_{Exp} - B_{Sim})^2}}{\sum_k (B_{Exp})^2}$$ \hspace{1cm} (3)

where the index $k$ runs over the 21 experimental data-points. Once expressed in percentage, this formula gives the following result for the reference and 2LUPS samples:

$$\delta = 0.8\% \quad \text{(Reference)}$$ \hspace{1cm} (4)

$$\delta = 6.8\% \quad \text{(2LUPS)}$$ \hspace{1cm} (5)

$$\delta = 18.8\% \quad \text{(2LUPS – central filaments “Off”)}$$ \hspace{1cm} (6)

This agreement indicates that the critical current at zero field $J_C = 3.52$ MA/cm$^2$, obtained from the data reported in Ref. [24], is consistent with the experimental observation for both the samples. Moreover, the experimental data show a significantly better agreement ($\sim 2.5\times$ lower relative error $\delta$) with the simulation which includes the narrow filaments of the central region. For the geometry of the samples under consideration the current density of 3.52 MA/cm$^2$ corresponds to the critical current $I_C = 387$ A, which matches well with the calculation performed with the conventional approach when applied to the reference tape (375 ± 27 A). Measurement of $I_C$ for the 2LUPS sample using the standard Hall scan setup and employing the conventional conversion, i.e. using the same parameters as for the reference sample, results in much lower $I_C$ values (dashed red line in Fig. 4). It is already known that a more careful Hall scan analysis is necessary for CCs presenting narrow filaments as shown in [29]. For this reason we recalculated the value of $I_C$ for each position $x$ along the length of the tapes, by considering the scale-factor which would give the best match with the field profile computed by FEM simulation. The results of this calculation are shown in Fig. 4 as solid lines. It can be noticed that our method leads to approximately the same result when applied to the reference tape. For the 2LUPS tape our approach is in much better agreement with the expected value of critical current than the conventional calculation. This indicates that the TAPESTAR™ setup can be used to interpret functionality of the 2LUPS filaments though a lower working distance between sample and sensor should, and if possible more sensors, be applied for future characterization to enable a more precise $I_C$ evaluation.

III. CONCLUSION

A new type of 2LUPS-based filamentary CCs was produced using Ion Beam Assisted Deposition and Metal Organic Vapor Deposition in commercial systems at SuperPower Inc. Microstructural characterization of the GYBCO superconducting layer cross-sections revealed that narrow non-functional regions are present between the upper and lower plateaus. In addition, individual filaments were also found to be physically decoupled when producing the superconducting layer using the large-scale MOCVD technique.

Hall-probe scanning using a commercial Theva TAPESTAR™ system, suitable for long length characterization of standard CCs, was employed for critical current measurements of samples. The standard parameters and conventional (magnetic field/current) conversion used in this commercial characterization system are not ideal for evaluation of filamentary CCs and significantly lower $I_C$ values were observed for the 2LUPS CC compared to a standard flat sample. We performed FEM computations taking into consideration the 2LUPS geometry and the expected $J_C(B)$ dependence, and obtained a good agreement between the numerical model and the measurement for both the reference and 2LUPS samples.

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