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Modelling dose rate to single grains of quartz in well-sorted sand samples: the dispersion arising from
the presence of potassium feldspars and implications for single grain OSL dating

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Abstract

Single grain OSL has become a widely used approach in Quaternary geochronology. However, the origins of $D_e$ distributions and the sources of variation in individual dose estimates are still poorly understood. The amount of scatter in these distributions on top of the known uncertainties in measurement and analysis is defined by overdispersion and this quantity is generally used for weighting individual $D_e$ values to calculate a central equivalent dose. In this study, we address the nature and amount of different sources of dispersion in quartz single grain $D_e$ estimates, by (i) using appropriate statistical tools to characterize $D_e$ populations and (ii) modelling, with a specifically designed GEANT4 code, dose rate distributions arising from the presence of potassium feldspar grains in well-sorted sands. The model uses Monte Carlo simulations of beta emissions and interactions in a random close packing of quartz and feldspar spheres representing a sand sample. Based on the simulation results, we explain the discrepancy between intrinsic and natural overdispersion values in a well-bleached sample, thus validating the model. The three parameters having the most influence on dispersion in dose rate distributions, and modelled in this study, appear to be grain size, potassium content and total dose rate.

Finally an analysis of measurement uncertainties and other sources of variations in equivalent dose estimates leads us to conclude that all age models (both logged and unlogged) which include an overdispersion value to weight individual $D_e$ values rely mainly on unknown parameters; this ignorance may lead to an inadvertent bias in $D_e$ estimates. Assuming counting statistics make a small contribution to dispersion (as is often the case), we suggest that in some cases it is most appropriate to use unweighted averages of equivalent doses when dividing by commonly measured average dose rates.

Keywords: Single grain OSL; dose rate distributions; age models; overdispersion; GEANT4 simulations
1. Introduction.

Quartz Optically Stimulated Luminescence (OSL) has become a widely used tool for establishing the chronology of sediment burial. The Single Aliquot Regenerative protocol (SAR: Murray and Wintle, 2000; 2003) allows the determination of individual equivalent dose estimates \( (D_e) \) from aliquots of arbitrary number of grains, including individual grains. Equivalent dose distributions derived from single-grain measurements are usually significantly dispersed, requiring some statistical treatment for their analysis; the choice of this statistical treatment can have a significant effect on the accuracy of the resulting OSL ages. For instance, post-depositional mixing of sediments (e.g., Tribolo et al., 2010) and/or insufficient resetting of the OSL signal before deposition (e.g., Jain et al., 2004; Olley et al., 2004) may lead to dose distributions where the central value is not representative of the sediment burial event. In single-grain equivalent dose analysis, the key concept of overdispersion (OD) is defined as the dispersion of results that cannot be explained by ‘within aliquot errors’, i.e. the measured or otherwise known uncertainties assigned to individual equivalent dose estimates (see Galbraith et al., 1999, for an introduction and discussion on its significance in OSL dating; see also Galbraith and Roberts, 2012).

Statistical models have been proposed to identify the \( D_e \) representative of the target event. For example the Minimum Age Model (MAM, Galbraith et al., 1999), the IEU (Thomsen et al., 2007; Jain et al., 2004) and the leading edge model (Lepper, 2001), have been suggested as tools to resolve the best-bleached component, and the Finite Mixture Model (FMM, Galbraith and Green, 1990; Roberts et al., 2000) has been suggested to identify individual dose components present in a mixture. These models require the input of an estimate of OD appropriate to the sample had it been well bleached; this can be either taken as a value presumed to be typical of well-bleached samples in general (i.e. <20 %, Jacobs et al., 2008a) or
experimentally determined from well-bleached samples with similar characteristics to those of the sample under investigation (Thomsen et al., 2007).

However, little is known about the nature and source(s) of overdispersion in single grain $D_e$ distributions. Thomsen et al. (2012) have demonstrated that overdispersion is dependent on dose in well-bleached samples irradiated with a known gamma dose; in two samples they found the overdispersion increased as the given dose increased. In naturally irradiated samples, different beta dose rates to different grains in sedimentary media are also expected to contribute to overdispersion in $D_e$ values (e.g., Mayya et al., 2006; Cunningham et al., 2012). These different dose rates arise because the range of beta particles is comparable to the size of sand grains, and to the inter-granular distance. In particular, the presence of hotspots – such as potassium feldspar grains, which generally represent an important source of dose rates in sands – generates skewed, wide dose rate distributions (Mayya et al., 2006; see also Brennan, 2006, for a discussion on the effect of hotspots on alpha dose rate distributions). Mayya et al. (2006) simulated beta dose rate distributions from individual potassium-rich feldspar grains to single 200 µm grains of quartz, and they showed that the dispersion in beta dose rates from potassium increases as the average potassium content (i.e. the number of feldspar grains) is decreased. Nathan et al. (2003) compared experimental and simulation results, using the Monte Carlo radiation transport code MCNP transport code, for different cases of heterogeneity in sedimentary environments. Despite weak agreement between experimental and numerical datasets, they showed that beta dose rate heterogeneity (either in the form of cold or hotspots) can influence single grain $D_e$ distributions. Cunningham et al. (2012) used MCNP to simulate dose rate distributions induced by NaOH grains containing artificially produced, short-lived $^{24}\text{Na}$ to mimic the effect of potassium feldspar grains. They were able to reproduce the shape of experimentally determined dose rate distributions, which can be fitted with log-normal distributions, but did not manage to get quantitative agreement between modelled and experimental data. Nevertheless, it is now clear that the presence of radioactive hotspots
induces positively skewed distributions of dose rates; conversely, the presence of coldspots such as calcareous blocks in ‘lumpy environments’ leads to negatively skewed distributions (see Brennan et al., 1997, for a study of gamma dose rates). These distributions are in contrast to those postulated by Jacobs et al. (2008b) who suggested that coldspots were the explanation for the two discrete modes in their dose distributions; both in view of the experimental and modelling results above, this seems unlikely (see also Guérin et al., 2013).

Despite this general understanding of the effect of hotspots in governing dose distributions, very few studies have compared experimental equivalent dose with simulated dose rate distributions. Recently Chauhan and Singhvi (2011) compared measured equivalent dose with modelled dose rate distributions, to assess whether the measured dispersion in $D_e$ values from multi-grain aliquots could be explained solely by dose rate distributions, or if an extra-source of dispersion such as poor bleaching was needed to explain the scatter in $D_e$ measurements. However, this study was not based on single grain $D_e$ measurements and it is not clear how many sensitive grains were present per aliquot. Moreover, the dispersion in $D_e$ values was taken as the standard deviation of individual estimates, and it did not account for the uncertainties on the individual $D_e$ values. In the absence of the knowledge of the effect of these uncertainties, it is difficult to interpret these results quantitatively.

2. Background

The purpose of this study is to study beta dose rate distributions from potassium feldspar grains to single grains of quartz in sand using the radiation transport toolkit GEANT4 (Agostinelli et al., 2003). In particular, parameters influencing these dose rate distributions are identified and the model has been tested on a well-bleached, well characterised sand sample. A statistical analysis of $D_e$ distributions from both natural and gamma dosed fractions of the sample are provided, and consequences regarding the use of various published age models is discussed.
Since $D_e$ estimates on individual grains have highly variable uncertainties, most OSL age models apply weighting factors to calculate representative equivalent doses. Moreover, most $D_e$ distributions reported in the literature exhibit overdispersion. In the most commonly used logged age models (such as for example the Central Age Model and the Minimum Age Model; Galbraith et al., 1999), the same relative OD (in%) is added in quadrature to individual relative $D_e$ uncertainties, assuming multiplicative error properties (i.e. absolute uncertainties proportional to doses); the weighted average of logged $D_e$ values (geometric mean) corresponds to the central dose. Conversely, in unlogged age models the same absolute OD (in Gy) is added in quadrature to individual absolute $D_e$ uncertainties, assuming additive error properties (i.e. constant absolute errors); the weighted average of $D_e$ values (arithmetic mean) corresponds to the central dose. In both cases the OD parameter is added in quadrature to each dose estimate in the weighted mean calculation of $D_e$. The choice between logged or unlogged models depends on the shape of measured $D_e$ distributions: multiplicative error properties lead to lognormal distributions (and to the choice of logged age models), whereas additive error properties lead to normal distributions (and to the choice of unlogged age models; for a discussion on this point, see Arnold et al., 2009).

Thomsen et al. (2012) tried to determine whether dose distributions from uniformly gamma-irradiated samples were normal or lognormal: they studied $D_e$ distributions of samples bleached in a solar simulator and then delivered a homogeneous well-known gamma dose, to study the nature of intrinsic overdispersion. They concluded that both logged and unlogged models provided reasonable, but not perfect fits to their $D_e$ distributions; in particular, they found no evidence for multiplicative error properties in equivalent dose measurements that could justify using logged age models.

For this study, a sand sample from a beach-ridge from Skagen (Denmark; see Buylaert et al., 2006; Nielsen et al., 2006; Guérin et al., 2012) was chosen for two reasons: firstly, because its OSL
properties satisfy the general criteria for acceptability of the SAR protocol (fast component, recycling, recuperation, dose recovery etc.) and in this area, the average OSL ages determined with large multi-grain aliquots of quartz are, for a number of sediment samples (n=20), in good agreement with radiocarbon data (Nielsen et al., 2006); secondly, the beta dose rate from potassium contributes a significant fraction (50 %) of the total dose rate to quartz; hence it is likely that, if dose rate distributions are affected by potassium and have implications regarding single-grain $D_e$ populations, such an effect will be observed in this sample. It thus is a good candidate to (i) model beta dose rate distributions from potassium and (ii) experimentally characterise the implications of such modelling for analysis of equivalent dose distributions. As a result, the effect of potassium feldspar grains on the dispersion of $D_e$ measurements from the natural distribution is presumed to be significant. Following Buylaert et al. (2006), this sample will be referred to as ‘the inter-comparison sample’.

3. Samples, material and methods

3.1. Sample preparation and characterization

Gamma spectrometry

Sediment was homogenised by crushing and sealed in a plastic box containing ~10 g of material. This sealed sample was then stored for at least three weeks to ensure radon build-up, before measurement using high resolution, low background gamma spectrometry, at the IRAMAT-CRP2A in Bordeaux. The potassium, uranium and thorium contents are given in Table 1. The corresponding dose rates have been calculated using dose rate conversion factors from Guérin et al. (2011) and using grain-size attenuation factors from Guérin et al. (2012). The accuracy in dose rate determination, using the infinite matrix assumption, has been questioned in general – and for this sample in particular – by Guérin et al. (2012), especially when it comes to grain-size attenuation factors for uranium and thorium. However, the exact value of the attenuation factors (constants) is not critical for our study since we are
only interested in comparing the equivalent dose and dose rate distributions in this sample; we therefore used attenuation factors for beta dose rates from uranium and thorium. The effect of moisture on gamma dose rate was taken into account following Guérin and Mercier (2012), using the mean grain size of the sample and using the cubic-centred packing model. For the effect of moisture on beta dose rates, we used the water correction factors from Nathan and Mauz (2008) in sediments containing no carbonates, which were indirectly confirmed by Guérin and Mercier (2012). Here it should be noted however, that these correction factors have not been adapted to sand samples (for which the geometry of energy emission and absorption has consequences on the effect of moisture on beta dose rate – see Guérin et al., 2012). For the potassium feldspar extracts, the internal dose rate was calculated using dose rate conversion factors for potassium (Guérin et al., 2011) and the self-dose values from Guérin et al. (2012), and assuming an internal potassium content equal to 12.5±0.5% (Huntley and Baril, 1997). Finally, the contribution from Rb was calculated according to Readhead (2002) and Huntley and Hancock (2001).

**Grain size analysis and element composition**

Grain size analysis and single grain element composition were obtained from Scanning Electron Microscope (SEM) image analysis and Energy Dispersive Spectrometry (EDS), respectively. Guérin et al. (2012) already modelled dose rates in this sample but their study focused on average dose rates to the different grain-size classes. Nevertheless, the sample characteristics were taken from this previous study: the grain size distribution can be found in their Fig. 1 (where the frequency corresponds to the actual number of grains rather than the most commonly used mass fraction). The sample is a well-sorted medium sand, with a mean grain size of 360 µm (geometric mean following Folk and Ward, 1957, calculated using the GRADISTAT program, Blott and Pye, 2001; in the following, all mean grain sizes are calculated accordingly). Based on EDS analysis, it is mainly (>99% by number of grains) made up of three
minerals: quartz (85% of the grains), potassium (7%) and sodium (8%) feldspar. Single grain EDS analysis further revealed that the grain-size distribution of potassium feldspar grains is similar to that of the sample taken as a whole. The potassium concentration, calculated from the abundance of potassium feldspar grains, and assuming a 12.5% K content of these feldspars (corresponding to the peak in the histogram of K concentration from single grains, Fig. 2 in Guérin et al., 2012) is ~1% by mass and compares very favourably with gamma spectrometry results (Table 1).

**Sample preparation**

Prior to mineral separation, the sample was wet sieved to isolate 180-250 µm sand grains. These grains were then treated with HCl (10%) to remove carbonates, and with hydrogen peroxide (H$_2$O$_2$) to remove organic contaminants; despite a weak reaction, both treatments were continued until no further reaction was visible. Two aqueous solutions of sodium heteropolytungstates (densities 2.58 and 2.62 g.cm$^{-3}$) were used to isolate K-rich feldspar fractions (<2.58 g.cm$^{-3}$) and quartz (>2.62 g.cm$^{-3}$). The quartz fraction was then etched with HF (40%) for 40 minutes to remove the outer portion of the grains affected by alpha irradiation. After etching, any fluoride contaminants were removed by rinsing with 10% HCl. This fraction was then re-sieved to >180 µm for further analysis, in particular for single grain measurements; this latter step removes any <180 µm grains resulting from the dissolution of residual feldspar in the quartz-rich fraction, or of small quartz grains.

**3.2. Luminescence instrumentation**

Grains were mounted in 9 mm base-diameter stainless steel cups using silicon oil. Aliquots of ~6 mm in diameter were measured for quartz, at the IRAMAT-CRP2A in Bordeaux, and of ~3 mm in diameter for feldspar extracts, at Risø. Luminescence measurements were made using Risø TL/OSL DA-15 and DA-20 readers (Bøtter-Jensen et al., 2003; 2010); for quartz multi-grain aliquots, blue (470 nm) light-emitting diodes (LED) were used with 7.5 mm Hoya U-340 detection filters; for feldspar, IR diodes
emitting at 875 nm were used in combination with coupled Schott BG39 and Corning 7-59 detection filters (transmission 320–460 nm). Each $^{90}\text{Sr}/^{90}\text{Y}$ source was calibrated during the measurement period by measuring several aliquots of calibration quartz irradiated with gamma rays (4.81 Gy; hereafter referred to as Risø calibration quartz) from a national secondary-standard $^{137}\text{Cs}$ source; this calibration has been independently confirmed by Bos et al. (2006).

Single grains of quartz were measured using an automated Risø TL/OSL reader (DA 20) fitted with a single grain attachment (Duller et al., 1999; Bøtter-Jensen et al., 2000). The grains were loaded into aluminium single-grain discs; each disc contains 100 holes 300 µm in diameter and 300 µm deep, on a 10x10 rectangular grid with 600 µm spacing between centres. A green laser (532 nm) was used to stimulate these grains individually, with light detection through a 7.5 mm Hoya U-340 glass filter. To confirm that only one grain was loaded into each hole, the single grain discs were visually inspected using a microscope before measurement. Radiochromic films allowed the determination of a coefficient of variation of 5.6% in dose rates to individual positions on the single-grain disc (Lapp et al., 2012). Correcting for this spatial variation in dose rates to single grains did not significantly change the measured $D_e$ distributions, so we used a single beta source dose rate for all grain positions.

3.3. Modelling: LSD algorithm and GEANT4

The model used in this study was already described in detail by Guérin et al. (2012) and a previous version of the GEANT4 code is available in Guérin (2011). Here GEANT4 (Agostinelli et al., 2003; Allison et al., 2006) is used to simulate the beta emission spectra from potassium feldspar grains (Fig. 1; such grains represent 7% of the total), and to track each primary (electron) and secondary (photon and electrons) particle transport individually in a random close packing of spherical grains. The random close packing is based on the Lubachevski-Stillinger-Donev (LSD) algorithm (Donev et al., 2005). The grain size distribution of the sample was determined experimentally by SEM image analysis (sample grains were
thinly spread on a glass plate to ensure no grain overlap). The equivalent radius of the grains was
determined assuming spherical grains (by equivalent radius of a grain we mean the radius of a circle
whose surface would correspond to apparent, generally irregular surface of the grain). The compactness
of the sediment obtained by random packing of the grains, using the LSD algorithm, is 0.635; as a result,
the density of the medium when air fills the pore space, is calculated to be 1.68 g. cm$^{-3}$.

The sample water content, as determined experimentally, is 12% - which corresponds to a
sediment density of 1.88 g.cm$^{-3}$. To obtain the same density for the wet sediment in our Monte Carlo
simulations, air is replaced by uniform, ‘light water’ (with a density of 0.55 g.cm$^{-3}$) in pore spaces; this
leads to a calculated wet density for the simulated sediment equal to the experimental value. Here it
should be noted that these dry and wet sediment density values corresponding to the simulations are
close to 'typical' sediment densities such as those given e.g. by Aitken (1985, Appendix H). The low
density, uniformly distributed ‘water’ is an approximation; in practice, surface tension effects alter the
spatial distribution of water (density: 1 g.cm$^{-3}$) in the pore spaces – water forms thin layers at the
surface of grains and tends to accumulate where grains touch each other. Such modelling goes beyond
the scope of this study, however, it is difficult to say if a more realistic distribution of water would
significantly affect the results of the simulations. For charged particles, the stopping power (unit: cm$^2$.g$^{-1}$)
determines the energy loss in the media, so for example, energy loss in 10 µm of water with a density
of 0.55 g.cm$^{-3}$ is equivalent to crossing 5.5 µm of identical water but with a density of 1 g.cm$^{-3}$; one can
ignore here 4.5 µm of air because of the negligible mass. As a result, in terms of energy loss in pore
space, the two scenarios are equivalent (light, uniformly distributed water, or dense, localised water and
air). However, some difference between the two cases will occur in terms of directional straggling; but
these are expected to even out on average.
Beta particles are emitted isotropically and their starting point is sampled homogeneously within the feldspar potassium grains. For simplicity, Guérin et al. (2012) simulated either pure potassium feldspar grains (with a K content of 14%, following stoichiometric values), or grains with zero potassium content. This assumption allows simplification of the simulations; however, the continuous distribution of K in the grains (cf. SEM-EDS analysis presented in Fig. 2 of Guérin et al., 2012) suggests that the actual potassium distribution may be somewhat less heterogeneous than in the model. Here, it should be noted that: (i) the potassium content of grains having a K content less than 6% are considered as zero potassium grains; this is considered acceptable since these grains represent only ~10-15 % of the total potassium in the sample; (ii) SEM-EDS analyses characterise only the surface of the grains, while the beta dose rate originates in the entire volume (so SEM-EDS values might not be representative of the content of the grains). We also observed low but non-zero values of K content from measurement of quartz grains, implying that at least some K is residing on the surface of all grains. Thus, the number of feldspar grains with intermediate K values is likely to be even lower than that observed in the data, suggesting that our assumed binary distribution of K should have little influence on the validity of the simulation results.

For tracking of both photons and electrons, Penelope physics datasets were used, as they are well-adapted to the simulation of low energy electromagnetic interactions (Salvat et al., 2001). Production cuts (i.e. range of secondary particles below which these secondary particles are not generated) and maximum step size were set to 20 µm to ensure accurate tracking down to one tenth of the diameter of the dosimeter grains of interest. In other words, the energy that would be carried away by a particle with a range of less than 20 µm was assumed to deposit locally, and the interaction probabilities were recalculated, by extrapolation of the provided Penelope datasets, every 20 µm along the particles tracks. To mimic infinite matrix conditions, a reflection algorithm was used (Nathan, 2011; Guérin et al., 2012).
Whereas in Guérin et al. (2012), the dose was only recorded in the grain-size classes of interest, in this study every quartz grain in the range from 180 to 250 µm in diameter is treated as an independent dosimeter; this allows us to obtain beta dose rate distributions from potassium feldspar to quartz grains. For each set of simulations (i.e. for each grain size distribution and potassium content), ten different random close packing configurations were used. For each configuration, the emission and tracking of 20,000,000 primary particles were simulated at the calculation centre of the French National Institute of Nuclear and Particle Physics (IN2P3). The uncertainties on the different numbers given in the following are obtained by taking the standard errors on individual values from the 10 different simulated configurations.

4. Results

4.1. Multi-grain aliquots OSL, IRSL and age control

For the inter-comparison sample studied here, the quartz OSL signal is dominated by the fast component. The SAR protocol (Murray and Wintle, 2000; 2003) was used with a preheat temperature of 200°C, held for ten seconds, and a cutheat temperature of 180°C before test dose measurements. The net signal intensity used in further calculations was derived from the sum of the OSL in the first 0.8 s of stimulation minus a background signal (calculated from the following 2.4 s of stimulation, i.e. early background subtraction). Nine aliquots were first exposed to a SOL 2 solar simulator for 3 hours and then given a dose of 5 Gy in the luminescence reader. The measured to given dose recovery ratio (0.97±0.05) showed that our SAR protocol was well-suited to measure equivalent doses for this sample. 21 equivalent doses were measured using multi-grain aliquots of quartz; the average recycling ratio was 0.99±0.07, and the resulting equivalent dose and age (4.73 ±0.23 ka) are shown in Table 2.

The IRSL from ~3 mm aliquots of K-rich feldspars was also measured (n=6); the corresponding equivalent dose is 6.90 ±0.30 Gy. A g-value of 2.8±0.2 %/decade was obtained from fading
measurements performed on the same aliquots. Using the fading correction from Huntley and Lamothe (2001), the resulting age of 4.28 ±0.27 ka is in good agreement with the quartz OSL age, which confirms that the quartz OSL signal was well reset at the time of deposition (cf. Murray et al., 2012). A post-IR IRSL at 290 °C (pIR-IR290; Thiel et al., 2011) dose of 13.7±0.6 Gy was obtained from six different aliquots, giving an apparent age of 6.69 ±0.36 ka. This age overestimation of ~2 ka is not surprising given the young age of the sample since it is well-known that residual, difficult-to-bleach doses affect post-IR IRSL D_e determination from young samples. It corresponds to a residual dose of ~6 Gy for this signal, which fits within the variability of observed residual doses for well-bleached samples (i.e., samples sufficiently exposed to sunlight to reset the quartz OSL signal; see, e.g., Buylaert et al., 2011). This further indicates that the quartz OSL from this sample is most likely unaffected by poor-bleaching.

4.2. Single grain OSL D_e and dose rate distributions.

The single grain D_e measurements were all made using the SAR protocol with a preheat at 260 °C for ten seconds, and a cutheat at 220 °C prior to test dose response measurement (note that thermal transfer is negligible for this sample, cf. Nielsen et al., 2006). The net signal used in D_e calculations was derived from the sum of the OSL in the first 0.05 s of stimulation minus a background signal (time average of the last 0.2 s; total stimulation time: 1s). Dose estimates from individual grains were accepted if they passed the following rejection criteria (derived from Thomsen et al., 2005; 2007; 2012): an error on the first test dose signal of less than 20% and a recycling ratio consistent with unity at two standard deviations. Recuperation was negligible for all samples. Note that the purity of the quartz extracts was examined on multi-grain aliquots using an IR-test (IRSL/BLSL ratios < 1%; Murray et al., submitted).

Fig. 2 shows the relationship between the first (‘natural’) test dose signal and measured equivalent dose for single grains (i) from the international calibration standard “Risø calibration quartz”
(batch 54, heated and then given a 4.81 Gy dose using a secondary-national standard $^{137}$Cs source in scatter free-geometry, Fig. 2a), (ii) from fractions of quartz from the inter-comparison sample exposed to a solar simulator for three hours and then given gamma doses of respectively 1.92, 4.81 and 9.62 Gy (Figs. 2b, c, d), and (iii) from natural quartz from the inter-comparison sample (Fig. 2e).

### 4.2.1. Single grain gamma dose distributions

**Table 3** lists a number of statistical characteristics of the equivalent dose distributions of figure 2, and resulting $D_e$ measurements derived using different statistical models: the Central Age Model (CAM; Galbraith *et al.*, 1999), the CAM$_{UL}$ (Arnold *et al.*, 2009) and a simple unweighted arithmetic mean; where relevant, dose recovery ratios are also given; all dose recovery ratios are within 10% of unity. Furthermore, they are all consistent with unity, within two standard errors (except for the 1.92 Gy dose recovery test, where the CAM$_{UL}$ gives a measured to given dose ratio equal to $0.93 \pm 0.03$).

The Risø calibration quartz and the inter-comparison sample show different average luminescence intensities in response to a fixed test dose of 2.2 Gy (first test dose signal). Furthermore, the average luminescence intensity of the signals induced by gamma irradiations in dose recovery experiments depends on the given dose. As a consequence, the average relative uncertainties on individual dose estimates vary between the different samples: 13% for the Risø calibration quartz (given dose: 4.81 Gy) and 27%, 21% and 13% for the inter-comparison sample for given doses of 1.92, 4.81 and 9.62 Gy, respectively (see **Table 3**). However, the relative overdispersion (OD) values from the CAM show little variation between Risø calibration quartz and the inter-comparison sample, or as a function of dose for the latter (16% on average; cf. **Table 3**); the different OD values for the gamma dose recovery experiments are statistically indistinguishable, which confirms the pattern seen by Thomsen *et al.* (2007; 2012) in the low dose region. Similar conclusions can be drawn for the CAM$_{UL}$, when the absolute OD (in Gy) is expressed as a fraction of the central dose. **Fig. 3** shows a standardised residual...
analysis in the form of quantile-quantile plots (see Galbraith and Roberts, 2012, for other examples of such plots and their discussion). Quantile-quantile plots can be used to visually assess the normality of the distribution of residuals from the models. The standardised residuals \((\delta_i - \delta)/\sigma_i\), where \(\delta_i\) is the \(i^{th}\) measurement of dose, \(\sigma_i\) its associated uncertainty – i.e., the quadratic sum of the analytical uncertainty and the overdispersion – and \(\delta\) is the central value determined with the model) are sorted and plotted against the estimates expected from a normed and centred Gaussian distribution. The 1:1 line indicates the expected fit to the data if residuals are normally distributed.

Interestingly, from Fig. 3 it can be seen that for the gamma dose distributions, the standardised residuals from both the CAM and the CAM_{UL} are consistent with a normal distribution, i.e. the observed residuals plotted against a normal distribution fall on a 1:1 line, despite a few outliers in the tail regions. In other words, it appears that the intrinsic overdispersion can be well described either by the same relative or the same absolute uncertainty; this makes the choice between normal and lognormal age models arbitrary at this stage.

**4.2.2. Dose rate distributions to single grains**

One of the differences between laboratory gamma dosed and the natural \(D_e\) distributions lies in the different dose rates to which individual quartz grains have been exposed in sedimentary media. Fig. 4 shows the results of the GEANT4 simulations of the single-grain beta dose rate distribution from potassium feldspar grains for the inter-comparison sample. This distribution is positively skewed (skewness: 1.07) and can be fitted by a lognormal distribution (red line), which is in agreement with previously published work (Mayya et al., 2006). The positive skewness can be understood as a result of few quartz grains being close to potassium feldspar grains (high dose rate tail of the distribution), whereas most quartz grains are at some distance – compared to the range of beta particles – from beta radioactive sources (mode of the distribution). The distribution has a relative standard deviation of 16
31.2±1.4 %. Note that the presence of hot-spots does not lead to any distinguishable bi-modality in the resulting dose rate distribution.

Grain to grain variations in gamma and cosmic dose rates can reasonably be assumed to be negligible in this sample, given the range of these radiations (>tens of cm). We assume that there is no other source of dispersion in beta dose rates to quartz grains – which is difficult to prove but supported by the absence of heavy minerals such as zircons, apatites, etc. from the SEM-EDS analysis (these could be potential sources of uranium and thorium). As a consequence, the relative standard deviation of the total dose rates to single grains of quartz (15.6±0.7 %) is obtained by multiplying the dispersion in beta dose rates from potassium by the relative contribution of this component to the total (50 %).

4.2.3. Over-dispersion in the inter-comparison natural sample

Equivalent doses measured for the natural portion of the inter-comparison sample are plotted against natural test dose responses in Fig. 2e. Standardised residual analyses from the CAM and the CAMUL are shown in Fig. 5. As for the gamma dosed populations, both models provide good fits to the experimental data, and the resulting equivalent doses are consistent with each other.

At least two contributions to the OD from the natural equivalent dose population have been quantified at this stage: (i) an intrinsic OD (i.e. the OD resulting from the measurement protocol; in other words we regard this intrinsic OD to originate from unrecognised/unquantified uncertainties inherent in the measurement, rather than as so-called ‘natural variations in the OSL properties’ giving rise to different true equivalent doses – cf. Galbraith et al., 2005). Our best estimate of this is determined from the gamma dose recovery tests in the dose range of interest (section 4.2.1); and (ii) an extrinsic OD (i.e. the OD resulting from all environmental factors external to the grains, such as the degree of light exposure before burial and grain-to-grain variations in dose rate). In this sample, we
consider the extrinsic OD to be dominated by the dispersion in dose rates (section 4.2.2), since we are confident that the sample was well bleached at deposition.

In this sample, we have determined the intrinsic OD from the CAM (15±3 % at ~5 Gy) and the standard deviation in dose rates (15.6±0.7 %); these can be summed quadratically to give a minimum estimate of OD that should be observed in the natural sample, of 22±3%. This compares very favourably with the measured OD (23±2 %); thus, it seems that the natural OD for the well-bleached inter-comparison sample can be fully explained by two contributions: the intrinsic OD and the dispersion in dose rates.

4.3. Factors influencing the dispersion in dose rates

Given that it appears that dose rate variations contribute about 50% to the total OD in our sample, it is now useful to investigate the factors influencing the dispersion in dose rate to single grains. This was done by varying several parameters of the GEANT4 model. Mayya et al. (2006) have already shown the effect of average potassium concentration on dose rate distributions in sands where potassium is located in potassium-rich feldspar grains: the skewness and dispersion of dose rate distributions increase as the number of potassium-rich grains is decreased (relative to the number of quartz grains). This can be understood by considering that the average distance between source and dosimeter grains is increased as the potassium content is decreased because of a reduction in the number of feldspar grains; as a result, fewer quartz grains are close to potassium sources and most are at a distance from any source. Similarly, one would then also expect that the average grain size of the sediments would have a similar effect on dose rate distributions: as the grain size is increased, the distances between source and dosimeter grains is also increased, which should lead to more skewed and more dispersed distributions.
Two parameters were thus varied in the simulations: firstly, the potassium content was varied by changing the fraction of potassium-rich feldspar grains, while keeping the grain size distributions similar for quartz and feldspar; secondly, the grain size distributions of the whole sample were multiplied by different scaling factors, so that the sorting of the sediments remained untouched but the simulated grain size distributions went from fine/very fine sands up to medium/coarse sands (Fig. 6).

Fig. 7 shows frequency histograms of beta dose rate distributions to quartz grains for various grain size distributions and average matrix (bulk) potassium concentrations. The relative standard deviation of these distributions is plotted in Fig. 8, for different grain sizes, as a function of average potassium content. As expected, the relative dispersion increases when the potassium content is decreased and/or when the grain size is increased, up to 135±19 % for a mean grain size of 637 µm with a K content of 0.14%.

We have compared our results, in terms of relative standard deviation in beta dose rates from potassium feldspar, with those presented by Mayya et al. (2006) in their Fig. 4. For 1% potassium and a unique grain size (200 µm), Mayya et al. (2006) found a relative standard deviation of ~28%; for a mean grain size of 255 µm, we found 20% and only 9% for 149 µm. It is not straightforward to understand these differences, partly because the two approaches are so different (in particular, Mayya et al. focused on determining the minimum dose due to the presence of hotspots: Mortheakai, Pers. Com.), and parameters may have different values. For example, the emission of beta particles in their paper is considered to be point-like, whereas in the Monte Carlo simulations the initial position within the emitting grains is sampled homogeneously. Straggling effects are taken into account in our Monte Carlo modelling, but not in Mayya et al. (2006). Furthermore, in their paper the minimum distance between a quartz grain and the closest hotspot – defined as the distance between the centres of the two corresponding grains (Mortheakai, Pers. Com.) – is 0; in other words, two grains can overlap, which is physically unrealistic. This may seem to be negligible, but it most likely explains the important high dose
tails in their Fig. 3. b (it should be emphasized that this did not affect the minimum dose due to the
presence of hotspots). Finally, there is an apparent peak in the dose distributions for very low doses (cf.
Fig. 3a), which according to the original authors is a numerical artefact (Morthekai, Pers. Com.); this
could contribute to the relative standard deviation in dose rates. It is very difficult to know at this stage
if one or more of these factors can explain the difference between our results and those from Mayya et
al. (2006). Nonetheless, despite these differences, the tendencies observed when parameters are varied
(in particular potassium content) remain the same. Because Monte Carlo simulations have fewer
approximations and closely mimic nature and because we model a more representative sediment
matrix, we tend to believe that the results of our simulations are better representative of the deviation
in dose rates due to the presence of hotspots compared to the approximate analytical treatment in
Mayya et al. (2006).

It should be emphasised here that, for a given potassium content and assuming that beta dose
rates from potassium are the only source of grain to grain dose rate variations, the dispersion on total
dose rates to quartz grains will decrease as the total dose rate is increased; this is because the relative
contribution to dose rate from potassium is decreased. In other words, the dispersion values from Fig. 8
should always be scaled by the relative contribution of beta dose rates from potassium to the total. In a
comprehensive study of more than 4,000 sediment samples from various contexts and geographical
locations, Ankjærgaard and Murray (2007) have shown that beta dose rates account on average for
~67% of the total dose rates (when working on sand-sized grains previously etched with concentrated
HF, i.e. not accounting for any alpha dose rate contribution). Moreover, for 95% of the samples, $^{40}$K
contributed between 40 and 92 % of the total beta dose rate. In other words, the contribution to the
total dose rate to quartz grains from the beta dose rate derived only from potassium ranged from 27 to
62% in almost all cases. Fig. 9 shows the modelled dispersion in total quartz dose rates as a function of
potassium content for three samples in each of which the total dose rate is fixed (at 1, 2 and 3 Gy.ka$^{-1}$).
The known likely range of potassium-derived beta dose rate contribution to the total (27 to 62%, derived above) are shown as dashed lines in Fig.9; these indicate the likely standard deviations to be expected for well-sorted sands of different mean grain sizes, for typical K concentration.

5. Discussion

5.1. Typical OD for well-bleached samples

Many age models (FMM, MAM, IEU) require the input of an OD value before the model can be used to identify a representative dose component(s). Determining an accurate OD representative of a well-bleached population is, therefore, at the heart of most single-grain studies. Sometimes, the intrinsic OD determined by a dose recovery experiment is used as a minimum value (e.g., Thomsen et al., 2007); in other studies the OD is allowed to vary between fixed values (e.g., Jacobs et al., 2008b).

The results from our Monte Carlo simulations show that a typical OD value for a well-bleached sample will depend on grain size, potassium content, and total dose rate. Considering the effect of potassium feldspar grains on dispersion in single-grain dose rates, there is no a priori limit on the OD of a natural sample and certainly no typical OD for well-bleached samples. This could explain the wide range of ODs observed in natural samples presumed to be well-bleached (but presumably affected by beta dose rate heterogeneities, see Fig. 1 of Thomsen et al., 2012, and references therein; see also, e.g., Jacobs, 2010; Jacobs et al., 2011; 2012; 2013; Gliganic et al., 2012).

5.2. When to use the dose rate model

It should be emphasised that the model presented here is expected to be used in cases where single grain dose rates need to be simulated to disentangle different sources of OD in single grain D_e measurements. In this study, the model successfully explains the discrepancy between the observed OD in the natural D_e distribution and the intrinsic OD resulting from the measurement protocol, for a single
sample. It is difficult to predict how well it might perform on a variety of samples. Nevertheless, our understanding of the processes involved allows us to be confident that dispersion in beta dose rates arising from the distribution of potassium will be most important when the average grain size is in the sand and gravel range (rather than silt or clay), the potassium content is low (<1%), and the total dose rate is small (<1 Gy.ka⁻¹).

5.3. Implications for the use of different age models

Interestingly, in the literature so-called age models (e.g., CAM, MAM, FMM) are actually dose determination models. Very few studies focus on dose rates during burial, even fewer consider dose rate distributions in the analysis of single-grain dose distributions. The simulation results presented in this study raise important questions concerning how luminescence ages are calculated. In particular, it is not clear how individual data should be weighted; each single grain equivalent dose estimates is not measured with the same precision, and each grain has received a different unknown dose rate; thus a dose distribution is not equivalent to an age distribution.

In the ideal case of a single ‘true dose’ (i.e., every grain has absorbed the same dose), the central dose (in this section, by central dose we mean the value most appropriate for use with an average dose rate to derive an age) is commonly derived using the logged or the unlogged central age model (see section 2). An alternative approach, commonly used in multi-grain analyses, is to use the unweighted arithmetic mean; this approach discards analytical uncertainties on individual dose estimates on the grounds that these uncertainties are trivially small compared to the variability in Dₑ measurements. The use of this latter approach inherently implies that the main source of dispersion is unknown and is much bigger than all known sources of analytical uncertainty.

The intrinsic dispersion in the Dₑ data can generally be equally well described by normal and lognormal distributions. However, the dose rate distribution from potassium feldspar grains is
positively skewed and can only be best described by a lognormal distribution; each individual dose
component of this distribution is sampled in the lab as a normal or a log normal distribution. In many
cases dose recovery distributions can be adequately fitted using the CAM (which assumes a log normal
distribution); this has been demonstrated here by the Gaussian distribution of the residuals from the
CAM (Fig.3 – right panel), using the inter-comparison sample. Given that log normal distributions can
describe both the natural dispersion arising from dose rates, and the measurement induced dispersion
in the data, suggests that the use of logged models best represent \( D_e \) distributions in natural samples,
and so provide the best estimates of OD.

The question is how best to determine the burial dose from a distribution of single grain \( D_e \)
values, in the absence of the knowledge of the underlying dose rate distribution. Almost all dose rates
based on high resolution gamma spectrometry, Neutron Activation Analysis (NAA), beta counting or any
other analytical technique are arithmetic means of repeated measurements of the spatially averaged
radioactivity in the sample. Thus, we typically know only an average dose rate in the sample. The dose
distribution on the other hand is known at the single grain level and different measures of central
tendency can be applied to derive a representative dose. If an age is derived by dividing a geometric
mean \( D_e \) by an arithmetic mean dose rate, then the age is likely to be underestimated to some degree
(since unweighted geometric means are systematically lower than unweighted arithmetic means). For
example, consider a sample in which the distribution of dose rates dominates the natural dispersion.
Suppose this to be a well-bleached, well-behaved sample in which the uncertainty on the measurement
of dose is negligibly small (e.g., 1%). All grains must by definition record the same age. First, consider a
10 ka old fine-grain sample with a uniform dose rate of 1 Gy/ka. We measure 3 grains, each with a dose
of 10 Gy: the average age is 10 ka and the CAM age is 10 ka with no overdispersion. Now let us consider
a less homogeneous (coarse grained) sample, of the same age and average dose rate, from which we
sample three representative grains which have experienced dose rates of 0.6, 0.9 and 1.5 Gy/ka
resulting in doses of 6, 9 and 15 Gy. We still observe the average dose rate (1 Gy/ka), the dose
distribution is positively skewed, and the individual grain ages of course remain at 10 ka; the average
age is 10 ka but the CAM age (using a geometric mean of D<sub>e</sub> values) is 9.3 ka. Thus, it appears that if the
scatter in the measured equivalent dose distribution arises primarily as the result of grain to grain
variability in the dose rates (which is much larger than the intrinsic variability due to measurement), and
one measures an average dose rate, then it is more appropriate to use the simple mean rather than the
dose rate (in any case, it is unlikely to be
explained by multiplicative error properties, because these should be entirely accounted for by the dose
recovery experiments). In such cases the central dose may be best estimated by using an unweighted
arithmetic average; the CAM will bias the results to give an equivalent dose inappropriate to the average
dose rate, and so an underestimate of the age. Obviously, this approach has the drawback that one
gives equal weight to individual D<sub>e</sub> values that are known with different degrees of precision.

To some extent, the above problem could be circumvented if a geometric mean dose rate was
available. However, only direct measurement of single-grain dose rate distributions would allow the
calculation of geometric mean dose rates; in general, such data is not obtainable experimentally. For the
inter-comparison sample, the geometric mean of the GEANT4 simulated single-grain beta dose rates
from potassium is ~5% lower than the arithmetic mean; in the absence of uncertainties on \( D_e \) measurements (or if they are negligible compared to other sources of dispersion), using a geometric mean to calculate a central dose (e.g., CAM) together with an average dose rate would thus result in an age underestimation of ~2% compared to that using the arithmetic mean for equivalent dose and dose rate (since beta dose rate from potassium contributes ~50% of the total). In a worst case scenario, the difference between geometric and arithmetic means of beta dose rates from potassium could be as much as 50% (based on a simulation assuming \([K] = 0.14\%\), mean grain size of 637 µm – cf. Fig. 8). Then the discrepancy arising from the division of a geometric mean equivalent dose estimate by an average dose rate would result in an age underestimation of 25%. Although this effect depends on potassium content, grain size and total dose rate, it should be noted that the consequence of using a geometric mean is a systematic age underestimation, not a random variation about some mean. In absence of modelled data such as presented in this study (or a better analytical approach), it appears that arithmetic means of \( D_e \) should give us the most accurate age.

### 5.4. A case study: the inter-comparison sample

In this sub-section, we discuss in practice the choice of an appropriate analysis model for \( D_e \) calculation, following the simulation results obtained for the inter-comparison sample. In the natural \( D_e \) distribution of the well-bleached inter-comparison sample, three main factors contribute to variations in \( D_e \) estimates, each with similar magnitudes: known measurement uncertainties, an intrinsic source of scatter in single-grain doses (measureable, but of unknown origin; Thomsen et al., 2005; Galbraith et al., 2005), and finally dose rate variations. Fig. 10 shows the relationship between the absolute (Fig. 10a) or relative (Fig. 10b) errors and individual dose estimates for the gamma dose recovery distribution obtained using the inter-comparison sample. The data are highly scattered, and no clear trend is observable; if anything, there may be a weak tendency for the
relative uncertainties to decrease with dose. Both the CAM and CAM_{ul} indicate that the data are overdispersed, with a magnitude similar to the dispersion of measurement uncertainties. Finally, Fig. 10 (c, d) shows the same plots of uncertainty as a function of dose for the distribution of D_{e} estimates from the natural inter-comparison sample. Again there is no unique trend, although overdispersion is increased presumably due to variations in external beta dose rates. In all four cases it is difficult to ascertain which age model, i.e. CAM, CAM_{ul} or unweighted arithmetic average, to use.

It is interesting to note that experimentally, for the equivalent dose populations investigated in this study, the CAM-based equivalent doses are greater than CAM_{ul} results (Table 3), because – even though there is no clear trend in the graphs from Fig. 10 – average relative uncertainties appear to decrease slightly with increasing dose (thus placing more weight on higher dose estimates when using CAM), and inversely for absolute doses (with the consequence of lowering CAM_{ul} central dose).

The fact that we measure non-zero intrinsic OD values tells us that we underestimate our uncertainties on individual D_{e} estimates, possibly due to variability in grain to grain natural OSL properties that is not accounted for by analytical uncertainties. When applying our GEANT4 based dose rate model, the only source of systematic variation in D_{e} distributions that we know of – albeit only through modelling – is the variation in dose rates to individual grains; it is difficult to justify using this source of dispersion in the weighting of De estimates (as is the case when using e.g. the CAM), as dose rate variability is independent of uncertainties of dose estimates. We may then decide to simply ignore our uncertainties in central D_{e} estimations and calculate unweighted averages of equivalent doses. Not only are the dose recovery tests satisfactory (cf. Table 3), but doing so we would also compare arithmetic means of equivalent doses with arithmetic means of dose rate. In the case of the inter-comparison sample, (i) the three identified sources of dispersion in D_{e} (analytical uncertainties, intrinsic and extrinsic OD) are of comparable sizes and (ii) the intrinsic OD is reasonably well-fitted by both
normal and lognormal distributions of corresponding uncertainties (cf. standardised residual analyses of
the CAM and the CAM_{ul}); it is not surprising that all three models (including an unweighted arithmetic
average) give a set of consistent ages. In general, we argue that only a careful analysis of the sources of
dispersion can lead to an informed decision regarding the most appropriate $D_e$ determination model on
a sample by sample basis.

Our study has not considered issues related to changes in OSL sensitivity of the grains between
nature and laboratory irradiations; any such changes can be a source of additional intrinsic over-
dispersion in $D_e$ populations (undetected by gamma dose recovery experiments; see, e.g., Stokes,
1994a, b; Singhvi et al., 2011). But given that there is a satisfactory agreement between the observed
OD and the predicted OD (from intrinsic OD and beta dose rate variation), we think this sensitivity
change effect must be negligible. From a dose rate perspective, the presence of highly radioactive
minerals such as, e.g., zircons will also induce additional extrinsic over-dispersion. Our Geant4 model
was designed to describe the effect of potassium feldspar grains on extrinsic over-dispersion, because
this was considered the most likely and most important source – at least for beta dose rates.

Nevertheless, similar models using the same architecture can be used to predict the effect of uranium
and thorium sources of any given geometry and size distributions.

6. Conclusion

We have developed a new model to quantify the effect of grain size and potassium
concentration (feldspar hotspots) on the grain-to-grain dose rate variations for well sorted sediments.
The model is successfully tested using experimental data obtained from a well characterised sediment
sample, and predictions are made for other sediments with similar sorting but different grain sizes and K
concentrations. The model provides estimates of minimum expected extrinsic overdispersion for various
grain size distributions, potassium contents and total dose rates. These estimates, together with an
analysis of over-dispersion in laboratory dose recovery data (intrinsic overdispersion), allow us to
investigate the sources of variability in equivalent dose measurements from individual grains in our test
sample. It is shown that consideration of beta dose rate variability has an important bearing on the use
of statistical models such as CAM, MAM, FMM, IEU, etc. for deriving the representative equivalent dose
that corresponds to the average dose rate estimate. Furthermore, our results imply that, for well
bleached samples, unweighted arithmetic mean dose together with the average dose rate may provide
a more accurate estimation of age, particularly in cases where the dispersion in measured $D_e$ values is
dominated by extrinsic over-dispersion rather than measurement uncertainties. This conclusion has
important implications for the analysis of more complicated dose distributions affected by incomplete
bleaching and post-depositional mixing.

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References


on Nuclear Sciences, 53, 270-278.

Ankjærgaard, C., Murray, A.S., 2007. Total beta and gamma dose rates in trapped charge dating based
on beta counting. Radiation Measurements 42, 352-359.

procedure for optical dating of young and modern-age sediments. Quaternary Geochronology 4, 306-
325.


**Figure captions**

Fig. 1: example of a Geant4 simulation of beta emission from potassium feldspars. The grains are randomly packed using the LSD algorithm (Donev et al., 2005). Blue spheres represent potassium-rich feldspar grains, whereas grey ones represent quartz grains. Electron tracks generated inside feldspar grains are shown in red, while secondary photon tracks are shown in green.

Fig. 2: single grain OSL natural test dose response as a function of dose for (a) Risø calibration quartz (annealed by heating and then given a 4.81 Gy gamma dose), quartz from the intercomparison sample bleached in a daylight simulator (Hönle SOL 2) at a distance of 80 cm for 3 hours and then given gamma doses of 1.92 Gy (b), 4.81 Gy (c) and 9.62 Gy (d), and natural (e).

Fig. 3: Q-Q plots of the standardised residuals from CAMUL (left) and CAM (right). Risø calibration quartz (a); quartz from the intercomparison sample: bleached and given gamma doses of 1.92 Gy (b), 4.81 Gy (c) and 9.62 Gy (d).

Fig. 4: Geant4-simulated beta dose rate distribution from potassium-rich feldspar grains to single grains of quartz in the intercomparison sample.

Fig. 5: Q-Q plots of the standardised residuals from CAMUL (left) and CAM (right), for the natural D_e distribution from the intercomparison sample.

Fig. 6: The different grain size distributions simulated with Geant4, corresponding to well-sorted sands. While the shape of the distributions is unchanged, the grain sizes are multiplied by different factors to investigate the effect of mean grain size on the single-grain dose rate distributions. The shaded bar indicates the 180-250 µm fraction, from which each grain is treated as an independent dosimeter.

Fig. 7: Examples of beta dose rate distributions from potassium feldspar in well-sorted sands for different potassium contents and mean grain sizes. The relative standard deviation is indicated in each case (RSD). Left: the potassium content is fixed (1.10%) but the mean grain size increases from top to bottom. Right: the mean grain size is fixed (360 µm) but the potassium content increases from top to bottom.

Fig. 8: Relative standard deviation, obtained with Geant4 simulations, of single grain beta dose rate distributions from potassium, as a function of potassium content for different grain sizes.

Fig. 9: Relative standard deviation, obtained with Geant4 simulations, of total dose rates to single grains of quartz, as a function of potassium content for different grain sizes. Total dose rate is 1 Gy.ka^{-1} (a), 2 Gy.ka^{-1} (b), 3 Gy.ka^{-1} (c). The dashed lines indicate the range of most likely values for potassium content in each case (see text for details).

Fig. 10: Single grain OSL absolute (left) and relative (right) uncertainties as a function of dose. (a), (b): intercomparion sample, 4.81 Gy gamma distribution. (c), (d): intercomparison sample, natural distribution. Inset in (c): without the 3 high uncertainty points.
<table>
<thead>
<tr>
<th>K (%)</th>
<th>U (ppm)</th>
<th>Th (ppm)</th>
<th>Water content (%)</th>
<th>Gamma dose-rate (Gy.ka(^{-1}))</th>
<th>Beta dose-rate (Gy.ka(^{-1}))</th>
<th>Cosmic (Gy.ka(^{-1}))</th>
<th>Total (Gy.ka(^{-1}))</th>
<th>Fraction contributed by beta from K</th>
</tr>
</thead>
<tbody>
<tr>
<td>1.06 ± 0.02</td>
<td>0.42 ± 0.02</td>
<td>1.38 ± 0.04</td>
<td>12</td>
<td>0.33 ± 0.01</td>
<td>0.74 ± 0.03</td>
<td>0.17</td>
<td>1.24 ± 0.06</td>
<td>0.50</td>
</tr>
</tbody>
</table>

Table 1: Radiometric and dose-rate data for the inter-comparison sample, as measured in IRAMAT-CRP2A.
Table 2: Ages obtained for the intercomparison sample, using quartz OSL, feldspar IRSL at 50 °C and post-IR IRSL at 290 °C.

<table>
<thead>
<tr>
<th></th>
<th>Quartz OSL</th>
<th>Feldspar IRSL at 50 °C</th>
<th>Feldspar postIR-IRSL at 290 °C</th>
</tr>
</thead>
<tbody>
<tr>
<td>( D_e ) (Gy)</td>
<td>4.73 ± 0.23</td>
<td>6.90 ± 0.30</td>
<td>13.7 ± 0.6</td>
</tr>
<tr>
<td>Age (ka)</td>
<td>3.81 ± 0.26</td>
<td>4.28 ± 0.27</td>
<td>6.69 ± 0.36</td>
</tr>
</tbody>
</table>

\(^a\) the Huntley and Lamothe fading correction model was applied.
<table>
<thead>
<tr>
<th>Given dose (Gy)</th>
<th>n(^b)</th>
<th>Avg. error (%)(^c)</th>
<th>CAM</th>
<th>CAM-UL</th>
<th>Unweighted average</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td>De (Gy)</td>
<td>OD (%)</td>
<td>Recovery ratio</td>
<td>De (Gy)</td>
<td>OD (% of $D_e$)</td>
</tr>
<tr>
<td>InterComp.</td>
<td>1.92 ± 0.03</td>
<td>85</td>
<td>27</td>
<td>1.86 ± 0.06</td>
<td>17 ± 3</td>
</tr>
<tr>
<td></td>
<td>4.81 ± 0.07</td>
<td>71</td>
<td>21</td>
<td>4.74 ± 0.16</td>
<td>15 ± 3</td>
</tr>
<tr>
<td></td>
<td>9.62 ± 0.14</td>
<td>108</td>
<td>13</td>
<td>9.63 ± 0.28</td>
<td>15 ± 2</td>
</tr>
<tr>
<td>Nat</td>
<td>4.51 ± 0.15</td>
<td>123</td>
<td>23</td>
<td>4.51 ± 0.15</td>
<td>23 ± 2</td>
</tr>
<tr>
<td>Risø Cal.(^a)</td>
<td>4.81 ± 0.07</td>
<td>369</td>
<td>13</td>
<td>108 ± 1</td>
<td>15 ± 1</td>
</tr>
</tbody>
</table>

**Table 3: Statistics of the different single-grain quartz $D_e$ distributions.**

\(^a\) for the Risø calibration quartz, all measured equivalent doses are given in seconds rather than in Gy, as the result is used for dose-rate calibration of the reader.

\(^b\) Number of accepted grains.

\(^c\) Average relative uncertainty on individual dose estimates.
Fig. 2:
Fig. 3: 

[Graphs showing standardised residuals against theoretical (normal) quantiles for four different scenarios labeled a, b, c, and d.]
Fig. 4:
Fig. 5:
Fig. 7
Fig. 8

Rel. std. dev. on beta dose rates from K (%) vs. Potassium content (%) for different thicknesses.

- $\mu = 149 \mu m$
- $\mu = 255 \mu m$
- $\mu = 360 \mu m$
- $\mu = 496 \mu m$
- $\mu = 637 \mu m$
Fig. 9

(a) Graph showing the relationship between rel. std. dev. on total dose rate (%) and potassium content (%). Different symbols represent different mean particle sizes: µ = 149 µm, µ = 255 µm, µ = 360 µm, µ = 496 µm, and µ = 637 µm.

(b) Similar graph as in (a) but with a higher potassium content range from 0.0 to 2.0%.

(c) Graph with a potassium content range from 0.0 to 3.0%.
Fig. 10