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CONCLUSION

We have compared the power requirements for pumping erbium-doped fiber amplifiers at 800, 980, and 1480 nm. The analysis used a quantitative amplifier model with measured emission and absorption cross-section spectra as inputs. It was shown that even under optimized conditions the pump power required to obtain similar amplifier performance is 7–8 dB higher when pumping in the 800-nm band. For small-signal amplification the required pump power for 980-nm pumping is 2 dB lower than for 1480-nm pumping. For a booster amplifier with a 1-mW input signal, the minimum required pump power for 1480-nm pumping is up to 1.7 dB lower than for 980-nm pumping, although the drive current will be lower for the latter laser. When the specified noise figure is lower than 4 dB the pump wavelength for the 800-nm band should be 806 nm (± 3 nm) and the pumping scheme should be codirectional. A detailed experimental investigation of a Ge/Al/P/Er-silica fiber pumped in the 800-nm band and at 980 nm showed that the pump power required to obtain similar small-signal gain for a 1551-nm

signal is ~ 7 dB higher for the 800-nm band, in good agreement with the model predictions.

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Optimum Design of Nd-Doped Fiber Optical Amplifiers

Thomas Rasmussen, Anders Bjarklev, Ole Lumholt, Mads Øbro, Bo Pedersen, Jørn Hedegaard Povlsen, and Karsten Rottwitz

Abstract—The waveguide parameters for a Nd-doped fluoride fiber amplifier have been optimized for small-signal and booster operation using an accurate numerical model. The optimum cutoff wavelength is shown to be 800 nm and the numerical aperture should be made as large as possible. Around 80% booster quantum conversion efficiency can be reached for an input power of 10 dBm and a pump power of 100 mW by the use of one filter.

INTRODUCTION

MOST of the optical communication systems installed operate at a signal wavelength of 1.3 μ m. Thus, optical amplifiers for this wavelength are attractive components. The recent demonstration of gain around 1.3 μ m in a

Nd-ZBLANP fiber [1] suggests a possible fiber amplifier, which will have the advantage over semiconductor amplifiers of being polarization independent and having very low insertion loss. So far the small-signal gain of the Nd-ZBLANP fiber amplifiers have been limited by saturation caused by a high level of amplified spontaneous emission (ASE) around 1050 nm. However, significant improvement has been predicted [2] and demonstrated [3] by effectively filtering of the 1050 nm ASE. The objective of this letter is to improve the efficiency both for small-signal and large-signal operation by optimizing the waveguide geometry. The analysis, which is based on a full-scale numerical model, is performed both with and without a 1050 nm filter. Optimizations of Er-doped fiber amplifiers have shown that confinement of the active ions increases the gain [4]. Thus, we have included this in the analysis of the Nd-amplifier.

MODEL

The model used is based on a numerical solution of the rate equations for a four-level laser system, with inclusion of excited state absorption (ESA) at the signal wavelength, as described in [2]. Furthermore, the model is extended to

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include the spectra of the 900, 1050, and 1300 nm ASE. The Nd-doped fiber has a step-index profile, and the pump and signal wavelengths are 0.795 and 1.34 μm , respectively. The spontaneous emission lifetime is 0.72 ms, the pump absorption cross section is $20.5 \cdot 10^{-25} \text{ m}^2$, the ESA (at 1.34 μm) and signal cross sections are $0.35 \cdot 10^{-25} \text{ m}^2$ and $5.15 \cdot 10^{-25} \text{ m}^2$, respectively, and the peak cross section value in the 1050 nm band is $27.2 \cdot 10^{-25} \text{ m}^2$.

In [2] it was assumed that the Nd^{3+} ion, after excitation to a higher energy level by the ESA, decays instantaneously to the ground state. This worst-case assumption is investigated in [5] where it is proposed that a fraction of the ESA decays to the upper laser level. To investigate the influence of this on the fiber design we have done calculations for both ESA decay to the ground state and the upper laser level. These calculations have shown no more than 5% differences in gain. In the following the worst case assumption is adopted.

SMALL-SIGNAL GAIN OPTIMIZATION

The optimum cutoff wavelength versus pump power (P_p) is shown in Fig. 1 for two numerical apertures (NA's). The dashed curves refers to a situation where an optical filter has been inserted at the optimum distance from the input end [6], which is determined as the distance where 50% of the pump power is absorbed. This filter is assumed to absorb the 1050 nm ASE totally, and leave all other wavelength bands unaffected. The insets shows the small-signal gain versus cutoff wavelength (λ_c) for NA = 0.1 and NA = 0.4. The figure indicates that one should choose NA as high as possible, and calculations for other numerical apertures confirm this. For small pump powers the cutoff wavelength should be chosen to 800 nm, but for larger pump powers, the cutoff wavelength is a function of both NA and the pump power. To assure single-mode operation of the signal, λ_c should be chosen less than or at 1300 nm. The figure also shows that if a filter is inserted the optimum cutoff wavelength is lowered towards 800 nm. A further analysis has shown that if the ASE at 1050 nm was totally removed throughout the fiber, the optimum cutoff wavelength would be 800 nm for all pump powers.

BOOSTER OPTIMIZATION

Fig. 2 shows the gain as a function of input signal power (P_s) for two different designs [(a) NA = 0.4, $\lambda_c = 0.8 \mu\text{m}$, (b) NA = 0.4, $\lambda_c = 1.3 \mu\text{m}$]. The pump power is 100 mW. In the small-signal region (the flat part of the curves) there is a significant difference between the two designs as discussed above, but in the booster-region (the decreasing part of the curves) no significant difference can be seen. This is also evident on the inset on Fig. 2, which shows the quantum conversion efficiency defined as $(G-1)P_s\lambda_s/P_p\lambda_p$, against λ_c for NA = 0.1 and NA = 0.4. G is the gain and λ_s and λ_p the signal and pump wavelengths, respectively. There is a maximum around 800 nm, but this is not significant. Notice that with one filter inserted one can reach a quantum conversion efficiency of nearly 80% with an input power of 10 dBm and a pump power of 100 mW.

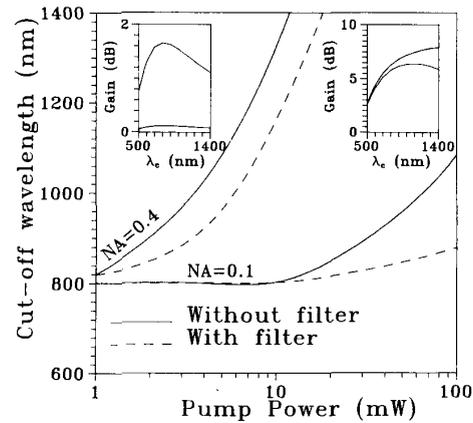


Fig. 1. Optimum cutoff wavelength versus pump power. The input signal power is -40 dBm. The left inset shows gain versus cutoff wavelength for a pump power of 1 mW, and the right for 100 mW. The upper lines are for NA = 0.4 and the lower for NA = 0.1. No filter is used in the insets.

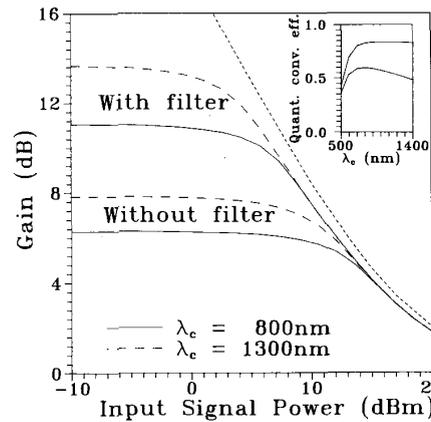


Fig. 2. Gain versus signal power for a pump power of 100 mW. The dotted line to the right is the theoretical limit, when the quantum conversion efficiency is 100%. The inset shows the quantum conversion efficiency versus cutoff wavelength.

CONFINEMENT

In order to investigate the influence which confinement of the Nd-doping has on small-signal gain and booster quantum conversion efficiency, we have calculated the curves shown in Fig. 3. A confinement of 50% improves the small-signal gain when the pump power is modest, but degrades the gain for larger pump powers. To explain this one has to consider the radial distributions of the involved intensities (all LP_{01} modes). Since the signal wavelength is 1300 nm and the dominant ASE wavelength is 1050 nm, the signal has the broadest field distribution. Therefore, confining the active ions will be more favorable to the ASE, than the signal. When the pump power level is small, the ASE power is also small, and confinement increases the signal gain.

In a booster the ASE is negligible since the signal power is large. Thus, confining should increase the quantum conversion efficiency, which is also the case for small numerical

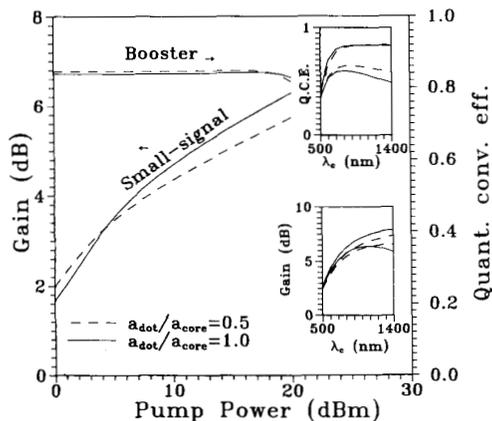


Fig. 3. Small-signal gain ($P_s = -40$ dBm) and booster quantum conversion efficiency ($P_s = 15$ dBm) versus pump power with 100% and 50% confinement of the Nd-doping. NA = 0.4, $\lambda_c = 800$ nm. The upper inset shows booster quantum conversion efficiency versus λ_c , the lower inset shows small-signal gain versus λ_c , both for 100 mW of pump power, and NA = 0.4 and NA = 0.1.

apertures as seen on the upper inset on Fig. 3. However, for the fairly large NA of 0.4, which is used in the main figure, the signal is well confined in the core so there is no significant improvement by confining the active ions in this case.

CONCLUSION

The waveguide parameters for a Nd-ZBLANP optical fiber amplifier has been optimized for the small signal and booster operation. In conclusion, the numerical aperture should always be chosen as high as possible, and the optimum cutoff wavelength should be chosen to 800 nm provided that the 1050 nm amplified spontaneous emission can be removed in

the small-signal case. A booster quantum conversion efficiency of nearly 80% can be reached for an input power of 10 dBm and pump power of 100 mW by the use of one filter.

By investigating two extreme cases of excited state absorption decay, we have shown that this does not influence the fiber design.

Finally, the Nd-dopant radius should be chosen equal to the core radius for the small-signal case. In the booster case there is an improvement in quantum conversion efficiency of 2.5 and 0.1% for numerical apertures of 0.1 and 0.4, respectively, when the Nd-ion distribution is confined to 50% of the core radius.

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