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Effects of Solvent Etching on PMMA Microstructured Optical Fiber Bragg Grating

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Abstract—Polymer Optical Fiber Bragg Gratings (POFBGs) can be solvent etched to achieve higher sensitivity sensors. To analyze the influence of etch on the mechanical and thermal performance of PMMA based microstructured POFBGs (mPOFBGs), we etched annealed and unannealed samples in pure acetone, while at the same time monitoring the induced Bragg wavelength shifts. Optical-thermodynamical interpretation of these shifts together with thermal and mechanical testing of etched fibers, reveal that solvent etching can lead to trapping of residual plasticizing solvent in the polymer matrix and consequently to change of Young’s modulus and thermal instability of the etched mPOFBG. Removal of these solvent residuals by thermal treatment results in a stable linear response with changes in strain sensitivity and Young’s modulus introduced by the etching process in this case do not have a big effect.

Index Terms—Fiber optics, Fiber Bragg gratings, Optical polymers, Etching, Optical sensors.

I. INTRODUCTION

Fiber Bragg gratings (FBGs) inscribed in polymer optical fibers (POFs) is a promising technology for a wide range of sensing applications [1]–[3]. They respond to external stimuli shift in the reflected Bragg wave length $\lambda_B$ according to eq. (1), where $n_{eff}$ is the effective refractive index of the fiber at the FBG location and $\Lambda$ is the periodicity of the grating.

$$\lambda_B = 2n_{eff}\Lambda$$

(1)

Compared with FBGs inscribed in conventional silica optical fibers, polymer optical fiber Bragg gratings (POFBGs) additionally offer several interesting properties such as increased sensitivity to axial stress due to the low Young’s modulus $E$ of polymers and increased strain sensing range due to the high elastic limit. This can be seen from eq. (2) showing the $\lambda_B$ shift, $\Delta\lambda_B$, of a POFBG with cross sectional area $A$ in response to an applied axial force $F$.

$$\Delta\lambda_B = \lambda_B(1 - P_c)\frac{F}{EA} = \frac{S_c}{EA}F$$

(2)

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$S_c = \lambda_B(1 - P_c)$, where $P_c$ is the strain-optic or photoelastic coefficient, is the strain sensitivity of the FBG depending on changes in $\Lambda$ and $n_{eff}$ when the fiber is stretched [4]. In several applications, it is of interest to further increase the sensitivity of FBGs to tensile forces, i.e. decrease the amount of axial force needed to produce a certain value of strain. For instance, this is highly useful when using the FBG in combination with a transducer for chemical sensing and biosensing applications [5]–[8]. This can be achieved by reducing $A$ in eq. (2). One way of doing so is by chemical etching, which has the advantage that it can be precisely controlled and is easily scalable [9]. Etching of PMMA-based POFs using acetone or acetone/methanol mixtures has previously been reported [10]–[14]. Rajan et al. [10] etched a PMMA-based POFBG humidity sensor using an acetone/methanol mixture to decrease the response time of the sensor. Together, Bhowmik et al. [14], Leal-Junior et al., [15] and Pospori et al. [16] have demonstrated that annealing and reduction of the PMMA fiber diameter by etching with acetone or a acetone/methanol mixture results in a reduced Young’s modulus and increased sensitivity towards strain, stress, force, temperature and pressure. The process of annealing in order to release residual stress of POFs has previously been demonstrated to be of great importance in ensuring stable sensing properties without hysteresis [17], [18]. Fasano et al. [19] demonstrated, that mixtures of methanol and water could be used to introduce stress relaxation of drawn POFs. In this paper, we present in greater detail, the effects of solvent etching on the thermo-mechanical properties of microstructured polymer optical fibers (mPOFs).

II. THEORY, EXPERIMENTS AND RESULTS

A. PMMA etching

Using Hansen Solubility Parameter (HSP) theory [20], it can be shown how well different solvents are able to etch thermoplastic polymers like PMMA. The polymer solubility is determined by the ratio of two distances in the three-dimensional HSP space. The coordinates of solvents and polymers in this “solubility space” are Dispersive ($\delta_D$), Polar ($\delta_P$) and Hydrogen bonding ($\delta_H$) cohesive energy density parameters like those given in Table I, where the values for the PMMA grade (GEHR PMMA®) used here have been determined experimentally from unannealed cane pieces. The other listed solvent HSP and molecular volume $MVol$ values are from the database included in the Hansen Solubility Parameters in Practice (HSPiP) software version 5.0.09. The HSP unit is $\sqrt{\text{MPa}}$ or sometimes $\sqrt{\text{cal} \cdot \text{cm}^{-3}} = 0.49 \cdot \sqrt{\text{cal} \cdot \text{cm}^{-3}}$. The
MVol unit is $\text{cm}^3\text{mol}^{-1}$. The ratio of the distance between polymer and solvent in HSP space ($R_0$), to that of the interaction radius ($R_0'$) of the polymer (also determined experimentally), the Relative Energy Difference ($RED$) number, tells how well a solvent dissolves the polymer. $RED$ below 1 means that the HSPs for polymer and solvent are all sufficiently similar for the polymer to be dissolved by the solvent cf. the rule of thumb "like dissolves like" [20]. The smaller the $RED$ value, the better the solvent is for dissolving the polymer. For $RED \approx 1$ the polymer swells. Far away from the polymers solubility sphere determined by $R_0$, $RED \gg 1$ and the solvents here do not affect the polymer. It is important to note, that it depends on the polymer and solvent how far away from $RED = 1$ dissolution or no interaction is observed.

HSPs and $R_0$ for a polymer depend strongly on the specific grade and treatment history. They are determined by testing the polymer in different solvents and fitting a solubility sphere to the observed dissolution, swelling and no interaction data. Important test parameters determining the HSP, $R_0$ and hence $RED$ number accuracies are e.g. the number of test solvents close to the solubility sphere determined by $R_0$, $RED \gg 1$ and the solvents here do not affect the polymer. It is important to note, that it depends on the polymer and solvent how far away from $RED = 1$ dissolution or no interaction is observed.

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In our case of unannealed PMMA POF etching, swelling takes place relatively close to $RED = 1$. Acetone is not a particularly good solvent, as it has a $RED$ of 0.97, however, since etching of a polymer fiber is actually a “controlled dissolution”, good etching solvents are expected to be somewhat bad solvents, such that the outer layers of the polymer are “slowly” removed. For the etch to be slow, a relatively large $MVol$ is also an important parameter, because large volume means slower diffusion into the polymer [20]. Trichloroethylene has different $\delta_D$, $\delta_P$ and $\delta_H$ values but is still a candidate for PMMA etching, since it has comparable $RED$ and $MVol$ values to that of acetone, see Table I. Experiments also show that trichloroethylene works well as an etching solvent. Chloroform is another solvent which dissolves PMMA well, but because it has both lower $RED$ and $MVol$ values, experiments confirm that it is difficult to produce a controlled POF etch. The $RED$ and $MVol$ values for methanol also clearly confirm that this solvent which is widely used for etching, when mixed with acetone, and doping of POFs, only swells the polymer. HSPs are additive and by mixing acetone with e.g. methanol it is possible slow down the etch by getting HSPs closer to but still below $R_0$ for a specific PMMA grade of interest [10–14], [22–25]. However, for our PMMA POFs (unannealed), there is no need to slow down the acetone etch and increasing the $RED$ just a little from 0.97 (acetone etchant), to 1.08 by instead utilizing the common 50/50 v/v% acetone/methanol mixture, gives crazing (microscopic cracks in the surface) in addition to slow etching. We believe this is probably also to some extent due to clustering properties of methanol. It should also be mentioned that we still observe this effect for annealed PMMA POFs, even though this treatment leads to residual stress relaxation and $R_0$ lowering, and is known to reduce the risk of Environmental Stress Cracking (ESC) and crazing [20].

### B. PMMA sorption and desorption of acetone

Originally, this work was initiated by a suspicion that solvent sorption could lead to long term plasticization and therefore influence on etched PMMA based POFBG performance. This was based on earlier observed pronounced long term plasticization when etching low $T_g$ (78°C) TOPAS POFs with cyclohexane. This softening could contribute to the explanations by Bhowmik et al. [14] why they observe that Young’s modulus of PMMA based POFs are lowered by etching.

In order to estimate the amount of residual acetone left in the material during the etching process, a 0.6851 g piece of PMMA fiber cane (of the same PMMA-grade as the mPOF used in this work) was placed in a petri dish, weighed on a high precision scale (0.1 mg resolution) and placed in saturated acetone vapor. The weight increase was then monitored at different times during the acetone sorption process. The PMMA piece was exposed to the solvent vapor instead of liquid because small drops of liquid on the sample surface are difficult to get off before weighing. Besides, exposure to saturated vapor, i.e. a Relative Humidity ($RH$) of 100 %, corresponds to liquid exposure like during etching cf. equation (3), where $a_s^v$ and $a_s^l$ are the solvent’s ($s$) activities in the vapor ($v$) and liquid ($l$) phases respectively.

$$RH = a_s^v100\% \leftrightarrow a_s^v = a_s^l$$

PMMA cannot be saturated with acetone because $RED < 1$ i.e. PMMA is dissolved and equation (4), where $\mu_s^v$, $\mu_s^l$, $\mu_s^g$ are the chemical potentials of the solvent in the liquid ($l$), vapor ($v$) and polymer ($p$) phases, will hold as long as liquid solvent is available to maintain a saturated vapor around the polymer.

$$\mu_s^v > \mu_s^l > \mu_s^g$$

After five days, the petri dish containing the cane piece was moved to room conditions (fume hood). At this time the shape of the piece had changed completely to a flat, round puddle at the bottom of the petri dish. The weight change was repeatedly monitored during desorption. The results are shown in Fig. 1. The desorption was fast in the beginning where most of the sorbed acetone was diffusing out of the structure, but then slowed down. After 35 days there was still more than 6 wt% acetone left in the PMMA. The petri dish containing the PMMA piece was then moved from room conditions to

<table>
<thead>
<tr>
<th>Material</th>
<th>$\delta_D$</th>
<th>$\delta_P$</th>
<th>$\delta_H$</th>
<th>$R_0$</th>
<th>$MVol$</th>
<th>$RED$</th>
</tr>
</thead>
<tbody>
<tr>
<td>PMMA</td>
<td>18.7</td>
<td>8.6</td>
<td>8.9</td>
<td>6.7</td>
<td>NA</td>
<td>NA</td>
</tr>
<tr>
<td>Acetone</td>
<td>15.5</td>
<td>10.4</td>
<td>7</td>
<td>NA</td>
<td>73.8</td>
<td>0.97</td>
</tr>
<tr>
<td>Trichloroethylene</td>
<td>18</td>
<td>3.1</td>
<td>5.3</td>
<td>90.1</td>
<td>0.98</td>
<td></td>
</tr>
<tr>
<td>Chloroform</td>
<td>17.8</td>
<td>3.1</td>
<td>5.7</td>
<td>NA</td>
<td>80.5</td>
<td>0.89</td>
</tr>
<tr>
<td>Methanol (cluster)</td>
<td>14.7</td>
<td>5</td>
<td>10</td>
<td>NA</td>
<td>100</td>
<td>1.36</td>
</tr>
<tr>
<td>Acetone/methanol 50/50 v/v%</td>
<td>15.1</td>
<td>7.7</td>
<td>8.5</td>
<td>NA</td>
<td>NA</td>
<td>1.08</td>
</tr>
</tbody>
</table>
C. FBG fabrication

The POF used in this work was a microstructured, 2-ring PMMA fiber fabricated in-house using the drill-and-draw method, see Fig. 2. It has an average diameter of 150 μm with hole diameters of 2 μm and a hole-to-pitch ratio of 0.40 making it endlessly single mode [26].

10 pieces of mPOF (40 cm each) from the same draw were used for grating inscription. 5 of the pieces were used as is. To study the influence of annealing on mPOFBG etching, the remaining 5 pieces were annealed in water at 65 °C for 5 days. This process releases residual stress originating from the drawing process, ensuring thermal stability of the grating below the annealing temperature [18].

The FBGs with \( \lambda_B \approx 850 \, \text{nm} \) were inscribed using the phase mask technique [2], [28], see the setup in Fig. 3. A 325 nm CW HeCd UV-laser (IK57511-G, Kimmon) with the output power attenuated to 18 mW was focused using a cylindrical lens through a phase mask with a grating pitch of 572.4 nm (Ibsen Photonics A/S) on the optical fiber while monitoring the peak growth using a reflection mode interrogation setup. Each FBG was inscribed in less than 7 minutes using this technique [29]. It should be mentioned, that using a pulsed KrF laser, FBGs in POFs can be inscribed in less than 30 seconds [30], [31]. The interrogation setup consisted of a supercontinuum laser (SuperK Versa, NKT Photonics A/S) launched into one of the arms of a 3 dB coupler (TW850R5A2, Thorlabs). One output arm was butt-coupled to the POFBG and the second arm terminated. The reflected light was then led back through the fiber coupler and into a CCD-spectrometer (CCS175M, Thorlabs) connected to a PC. The typical reflection spectrum for an unannealed PMMA mPOFBG is shown in Fig. 4. A dynamic gate algorithm [32] was applied on the acquired spectra to achieve a sub-pixel resolution of ~10 pm of the determined Bragg-wavelength. The \( \lambda_B \) of the resulting mPOFBGs was ~849 nm in the unannealed fibers and ~847 nm in the annealed fibers. The difference of ~2 nm in \( \lambda_B \) is due to the annealing induced changes in refractive index from stress relaxation and perhaps birefringence when the PMMA molecular structure gets less ordered.

D. Etching of mPOFBGs

We believe the changes in POFs during etch are due to combinations of five more or less interrelated, simultaneous, reversible (rev.) and irreversible (irrev.) mechanical and optical effects of varying magnitude, leading to both positive (+) and negative (−) \( \lambda_B \) shifts. Which effects dominate regarding induced \( \lambda_B \) shift at a certain time depends on the POF material, processing, solvent used for etching etc. and can be estimated by following the \( \lambda_B \) shift during the etch. Solvent sorption and etch leads to:

a) Expansion/initial swelling due to added volume of solvent (+, rev.).

b) Refractive index change from mixing with solvent having different refractive index than the fiber (+, −, rev.).

c) Contraction from stress relaxation/plasticization by polymer molecular rearrangement and etch removal (−, irrev.).

d) Refractive index change from stress relaxation by polymer molecular rearrangement and etch removal of tensile surface stress leading to densification (+, irrev.).
Reflection Intensity (dB)

Wavelength (nm)

Fig. 4. Reflection spectrum for a typical unannealed 150 µm diameter PMMA 2-ring mPOFBG. FWHM is approximately 2 nm.

e) Refractive index change purely from plasticization (+, rev.).

Ueberrreiter [33] divides the dissolution zone between pure solid polymer and pure solvent in four separate layers: Infiltration, solid swollen, gel and liquid. It is the changes going on during etch in the inward moving infiltration, solid swollen and perhaps to some extent gel layers, which induce changes in \( \lambda_B \). Swelling starts when the free volume between the polymer chains in the infiltration layer has filled up and initially, when there has been no significant stress relaxation, it is strong enough to expand the fiber along its axis \( \alpha \) and will contribute to a positive shift of \( \lambda_B \), cf. eqs. (1) and (2). It has even been shown, that polymers (hydrogels) which swell/de-swell in specific chemical environments, can be placed on etched mPOFBGs outside the cladding and in this way induce \( \lambda_B \) shifts for chemical sensing [34]. However, if the polymer is made too soft, this transducer principle is not working [34], [35]. The \( n_{eff} \) related \( \lambda_B \) shift due to presence of solvent in the fiber \( \lambda_B \), is, according to the Lorentz-Lorenz equation for binary mixtures, eq. (5), [36], negative if the solvent refractive index is lower than that of the polymer \( (n_{PMMA} = 1.4906) \) [37]) and vise versa. In eq. (5), \( n_{12} \) is the refractive index of mixture of component 1 and 2, \( \alpha = \frac{n_1^2 - 1}{n_1^2 + 2}, \beta = \frac{n_2^2 - 1}{n_2^2 + 2}, \)

\( n_1, \) and \( n_2 \) are the refractive indices of the pure component, \( \phi_i = \frac{w_i \rho_{12}}{\rho_i} \), with weight fraction \( w_i = \frac{m_i}{m_1 + m_2}, \rho_i \) the density of the pure components, \( i = 1, 2 \) and \( \rho_{12} \) the mixture density.

\[
n_{12} = \sqrt{\frac{1 + 2\phi_1 \alpha + 2\phi_2 \beta}{1 - \phi_1 \alpha - \phi_2 \beta}}
\]  

Free energy \( \Delta G = \Delta H - T \Delta S \), where \( T \) is the absolute temperature, is stored as residual stress in the fiber during the drawing process by increase in latent heat \( \Delta H \) and entropy \( \Delta S \). During stress relaxation the stored energy is released again in form of heat and molecular disordering. Outer stressed layers removed by the etch or plasticized by sorption leads to stress redistribution in the remaining fiber. At equilibrium, the axial stress distribution \( \sigma_{ax}(r) \), where \( r \) is the distance from the fiber center at some cross section with area \( A \), is balanced and equation (6) holds [39].

\[
\int_A \sigma_{ax}(r)dA = 0
\]  

So, if the fiber has been drawn as typical for optical fibers [40] and wires [39], such that, at equilibrium, axial residual stress is mainly compressive around the center and tensile closer to the surface, plasticization or etch will result in less tension stress counteracting the compressive and the fiber will shrink, which will cause contraction e) of the FBG and therefore negative shift of \( \lambda_B \). When the fiber tensile strain \( \varepsilon \) is reduced, its volume \( V \) decreases with \( \Delta V \) according to eq. (7), [41]:

\[
\Delta V = V(1 - 2\nu)\varepsilon
\]  

where \( \nu \) is Poisson’s ratio \( (\nu_{PMMA} = 0.4) \), [42]). At the same time, this means that the fiber density \( \rho \) and hence also the refractive index \( n \) increase according to the Lorentz-Lorenz equation for pure substances [43], eq. (8),

\[
n = \frac{M + 2\rho R}{M - \rho R}
\]  

where \( M \) is the molecular weight, \( R \) the molar refraction and \( \rho \) the density. In eq. (2), this FBG \( n_{eff} \) dependency on strain induced volume change is represented by the photoelastic coefficient \( P_c = \frac{\Delta n_{eff}(\varepsilon)}{\Delta V} = \frac{1}{2}B_2 - \nu(B_1 + P_2) \) \( \Delta V = V(1 - 2\nu)\varepsilon \), where \( B_1, P_2 \) are Pockel’s coefficients of the strain-optic tensor [4], [44]. In this way d), the shrinkage also contain a \( \lambda_0 \) red-shift component by increasing \( n_{eff} \). Here, it should also be mentioned, that the more crystalline a polymer is, the higher are it’s density and therefore refractive index, cf. eq. (8). Thus, when the fiber is relaxed, it looses some degree of long range "crystalline" ordering or anisotropy, which might oppositely lead to density and \( n_{eff} \) lowering \( \lambda_B \) blue-shift. However, we believe this effect is small, because inherently, PMMA is not a semi-crystalline thermoplastic polymer. Plasticization in itself, e), will also affect \( n_{eff} \) in eq. (1). For PMMA, Lamble et al. [45] have demonstrated that the refractive index and the stress-optic coefficient increase linearly for increasing content of the plasticizer dibutyl phthalate (DBP). By using DBP with \( n_{DBP} = 1.4925 \approx n_{PMMA} = 1.4909 \), they ensured that this effect was mainly due to plasticization and not a refractive index binary mixture effect. The maximum amount of DBP was 20 wt% giving a refractive index change of 0.0048, which
in our study would correspond to a $\lambda_B$ change of $\sim$3 nm. So, based on this work, we anticipate a $\lambda_B$ red-shift component due to plasticization related $n_{eff}$ increase.

The stronger the solvent interacts with the polymer, the faster are the processes a)-e). For instance, the speed of POF etching is strongly influenced by the initial fiber stress relaxation level, see Fig. 5. Unannealed POFs etch much faster than annealed because residual stress lowers the free energy barrier for dissolution. In terms of HSP thermodynamics, for unannealed, unrelaxed POFs, $R_0$ is larger and therefore RED lower for the same solvent. When the polymer solubility is increased by better solvent HSP match, so are swelling, dissolution and solvent diffusivity in the polymer [20], and consequently etch speed. Detailed modeling of polymer dissolution kinetics has been described by e.g. Miller Chou et al. [46] and Narasimhan et al. [47], [48]. For small molecules diffusion is often much faster than polymer chain stress relaxation movement. This is seen from the study presented in Fig. 6. When an unannealed POFBG is immersed in water, expansion diffusion initially dominates over $\lambda_B$ blue-shift due to any total refractive index lowering b), d), e), $n_{\text{water}} = 1.3324$ and residual axial stress relaxation shrinkage c), i.e. fiber length reduction. The three contributions to refractive index change are small at the PMMA water saturation (1.5 wt%) and of opposite sign and together do probably not contribute much to the total $\lambda_B$ shift. When the fiber has become saturated with water, it is sufficiently soft for $\lambda_B$ stress relaxation blue-shift c) to become dominating. For the annealed fiber there is no residual stress in the fiber and when it is immersed in water only $\lambda_B$ red-shift a) is observed. Again, some minor negative b) and positive d), e) contributions to the shift are expected to come from the refractive index changes. For unannealed PMMA based POFs we have earlier seen a similar behavior on a much longer time scale in methanol/water 50/50 v/v% [19]. The concentration of small water molecules in this mixture is high enough to, initially, induce a fiber length increase a) dominating over any total refractive index lowering b), d), e) and residual axial stress relaxation shrinkage c), i.e. we observe a $\lambda_B$ red-shift. After some time this change to dominating POF length reduction c), and thus $\lambda_B$ blue-shift due to swelling induced stress relaxation. However, in this case with methanol ($n_{\text{methanol}} = 1.3145$, [37]) swelling, the saturation level is much higher and the refractive index negative contribution b) to the shift might be larger. When the RED number is lowered by increasing the methanol concentration to 60 v/v% or 70 v/v%, $\lambda_B$ blue-shift b), c) dominates at all times over the total red-shift a), d), e) probably mainly due to increased swelling induced stress relaxation speed c). Likewise, stress relaxation c) $\lambda_B$ blue-shift dominates at all times during annealing of the fibers in this study at 65 °C in water, see Fig. 7. The observed annealing blue-shift is 81 nm, which is similar to the shifts observed by Fasano et al. [19] for the methanol/water solvent relaxation mixtures. The initial dominance over red-shift contributions a), d), e) is mainly due to much faster temperature than water equilibrium. Later, like for the methanol/water mixtures with high methanol content, the elevated temperature also contributes to the dominance through lowering of the PMMA-water RED number, which leads to higher water saturation concentration b), c).

Fig. 5. Diameter etching rate in pure acetone of unannealed and pre-annealed PMMA mPOFs with initial diameters of approximately 150 µm.

![Graph of Diameter etching rate in pure acetone of unannealed and pre-annealed PMMA mPOFs with initial diameters of approximately 150 µm.](image)

Fig. 6. Bragg-wavelength response of two 150 µm diameter PMMA mPOFs after being immersed in distilled water at room temperature. The annealed mPOFBG absorbs water and saturates with a total shift of $\sim$2.7 nm. The unannealed mPOFBG starts annealing at room temperature as water is absorbed.

![Graph of Bragg-wavelength response of two 150 µm diameter PMMA mPOFs after being immersed in distilled water at room temperature.](image)

a) of the fiber and therefore red-shift of $\lambda_B$ due to fast water

Fig. 7. Bragg-wavelength response for a 150 µm diameter 2-ring mPOFBG during annealing in water at 65 °C. After 32 hours, the Bragg-wavelength had shifted from 849 nm to 786 nm (a shift of 81 nm).

In this work, to study the influence of etching on the mechanical and thermal mPOFBG properties, the end tips of the mPOFs were sealed using a UV curable adhesive (NOA78, Norland Products, Inc.) to avoid acetone from entering the

![Graph of Bragg-wavelength response for a 150 µm diameter 2-ring mPOFBG during annealing in water at 65 °C.](image)
air channels of the fiber. The mPOFBGs were etched in pure acetone (99.9 %, Sigma Aldrich). In works done by others [10]–[14], [22]–[24], mixtures of acetone and methanol have been used to improve surface quality, but as explained, in our case the best surface quality was obtained using acetone only. The etching process was done while monitoring the Bragg-wavelength by immersing ~8 cm of the fiber (including the FBG) in a test tube filled with acetone placed in a water bath at 23.4 °C for thermal stabilization. Fig. 8 shows the behavior of three mPOFs annealed before FBG inscription and etched from 150 µm diameter to 125 µm, 92 µm and 56 µm. Note the much smaller time scale and \( \lambda_B \) shifts when etching, compared to the earlier presented work. Since the POF is stress free, the initial negative \( \lambda_B \) shift cannot be attributed to stress relaxation \( c, d \), which is irreversible, but only negative change in effective refractive index due to high concentrations of acetone \( b \) \( n_P = 1.47, n_A = 1.35 \) \([37]\). This effect largely dominates over any \( \lambda_B \) red-shift due to sorption induced POF expansion \( a \) and plasticization refractive index increase \( e \). Importantly, the shift is very reversible, i.e. most of the shifts have disappeared after a desorption period. However, for the fiber etched to 56 µm there is still a small blue-shift after 20 mins. desorption, which indicates that there is still some acetone left in the material. Evidently, also at this time, the negative effect of \( b \) is larger than the positive of \( a \) and \( e \). The fast and small \( \lambda_B \) red-shifts observed when taking the POFBGs out of solution are due to cooling when the liquid acetone at the surface evaporates. After cooling, the negative shift continues for as long as there is solvent at high concentration near the fiber surface, which can diffuse to lower concentrations closer to the sensing FBG in the fiber core. Fig. 9 shows the behavior of three POFs unannealed before FBG inscription and etched from 150 µm diameter to 120 µm, 100 µm and 74 µm. Note again the faster etch due to lower \( \text{RED} \) number induced by residual stress. During etch, these POFs behave much like the described similar fibers dipped in water, cf. Fig. 6, and methanol/water 50/50 v/v\% \([19]\). Initially, for all three fibers, \( \lambda_B \) red-shift effects \( a \), \( d \), \( e \) dominates over \( b \), \( c \), due to faster diffusion than stress relaxing polymer chain rearrangement.

![Fig. 8. Response of FBGs inscribed in pre-annealed mPOFs during etching in pure acetone. Vertical dashed lines represent the time at which each sensor was removed from the solvent and left to stabilize under room conditions.](image1)

![Fig. 9. Response of FBGs inscribed in unannealed mPOFs during etching in pure acetone. Vertical dashed lines represent the time at which, each sensor was removed from the solvent and left to stabilize under room conditions.](image2)

Note the much smaller time scale and \( \lambda_B \) shifts when etching.

Due to the hygroscopic nature of PMMA, the thermal response of the pre-annealed mPOFBGs was measured with the FBG immersed in water. A test tube filled with demineralized water was placed in a water bath on a feedback controlled heating plate while stirring. A thermometer was placed in the test tube to monitor the temperature. The temperature was stabilized at 24.5 °C ± 0.1 °C and the mPOFBG was left in the test tube overnight to completely saturate the FBG with...
water, causing a positive shift in the Bragg-wavelength of $\approx 2.7$ nm, see Fig. 6. In this way, it is ensured that the measured thermal response of the FBG is not influenced by changes in the relative humidity but only small changes in PMMA saturation level as the temperature is changed. The thermal response was only measured on the pre-annealed fibers below the annealing temperature of 65 °C to avoid any influence from thermally induced relaxation of the mPOFBGs.

The thermal responses of an un-etched and etched (diameters: 150 µm, 102 µm and 56 µm) mPOFBGs are shown in Fig. 10 and Fig. 11, where the zero-point of the Bragg-wavelength shift has been chosen, such that $\Delta \lambda_B = 0$ at 24.5 °C for all the FBGs. As seen in Fig. 10, showing the temperature. The etched mPOFBGs, on the other hand, exhibit a nonlinear response as the temperature is increased but linear response as temperature is decreased. This behavior is associated with desorption of the residual acetone originating from the etching process as the fiber is heated, causing a positive shift of the Bragg-wavelength, i.e., the reverse of $\mathbf{b})$ in the previous section. In other words, acetone with a lower refractive index than that of the fiber is driven out, which leads to increase of $n_{eff}$. Note also that, the gain in $\lambda_B$ correspond well with that lost at the end of etch, Fig. 8. As the mPOFBG is subsequently cooled, no further desorption is happening, and the thermal response becomes linear as in the un-etched case. In Fig. 11, the thermal response of the etched mPOFBGs related to the second cycle of temperature variation is shown and compared with the un-etched mPOFBG, showing very similar responses with sensitivities of $-31.9$ pm/°C ± 0.8 pm/°C (56 µm) and $-31.5$ pm/°C ± 0.5 pm/°C (102 µm). This demonstrates, that reducing the diameter of the fiber by solvent etching, does not have a significant impact on the thermal response of the mPOFBG. This result is in contrast to the results presented by Bhowmik et al. [11], [12], [14], who found linear diameter correlated thermal responses with significantly increased sensitivity of etched step-index PMMA FBGs. The cause of this discrepancy may be that the thermal variation in these references was done without humidity control. As the temperature is increased the relative humidity decreases, which means that the surrounding air can hold more water. This induces fiber desorption of water and an additional blueshift of the Bragg-wavelength, which could dominate the redshift caused by solvent desorption. The humidity induced shift becomes even more apparent as the fiber diameter is reduced due to a rapidly lowered response time [10]. A nonlinear, diameter-independent thermal response of etched step-index PMMA-PFBGs was also observed by Hu et al. [24] at a constant relative humidity of 50 %, however, the grating response for decreasing temperature was not shown, and no further elaboration on the observed behavior was made. Furthermore, it is not mentioned in these papers whether their fibers were pre-annealed and up-down cycles are missing.

\[ F. \text{ Mechanical sensing characteristics of etched mPOFBGs} \]

The strain sensitivity and force-sensitivity of both the etched and unetched mPOFBGs were measured using the setup shown in Fig. 12. The mPOFBG was mounted in clamps between a fixed force meter (Thorlabs FSC102, 1 mN resolution) and
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Young’s modulus of etched mPOFs. The method of calculation and related uncertainties are described in Appendix A.

The strain sensitivity was determined as the slope of the acquired ($\Delta \varepsilon$, $\Delta \lambda_B$) data for each sensor. Young’s modulus $E$ was calculated from eq. (9),

$$E = \frac{4kL}{\pi D^2} \tag{9}$$

where $k$ is the spring constant in Hooke’s law, $\Delta F = k \Delta L$, $L$ is the distance between the fixing points holding the mPOFBG and $D$ is the fiber diameter. Thus, by evaluating $k$ from the experiments described above, Young’s modulus of the mPOFs could be determined. The results are shown in Fig. 13 and Fig. 14 and demonstrates that both the strain sensitivity and Young’s modulus are affected by the etching process, the most significant changes being observed for the unannealed mPOFBGs. The obtained strain sensitivities are shown in Fig. 13 and are seen to increase, as the diameter of the fiber decreases due to solvent etching. The magnitude of change, however, is quite modest. For the unannealed mPOFBGs a change from 0.68 pm/µε ± 0.01 pm/µε (150 µm diameter) to 0.80 pm/µε ± 0.01 pm/µε (57 µm diameter) corresponding to 18 %, is observed. For the pre-annealed mPOFBGs, it changes from 0.73 pm/µε ± 0.02 pm/µε (150 µm diameter) to 0.79 pm/µε ± 0.01 pm/µε (56 µm diameter) corresponding to a 8 % change. It is also observed, that there is a significant difference in $E$ between the unannealed and annealed mPOFBGs, see Fig. 14. The modulus of the unannealed mPOFBG was measured to be 3.21 GPa ± 0.07 GPa whereas it for the annealed mPOFBG was 2.75 GPa ± 0.10 GPa. This is expected, as the polymer chains in unannealed fibers have a larger degree of orientation along the fiber axis [9], [38] and thus carry a greater part of the load on the stronger covalent bonds [50]. The Young’s modulus of the unannealed fiber is seen to decrease during the etching process. The modulus for the fiber etched to 57 µm was measured to be 2.66 GPa ± 0.15 GPa, which is a 17 % decrease compared to the modulus of the unetched fiber. Only a small decrease in the measured Young’s modulus of the annealed fibers due to solvent etching is observed, and the measurement uncertainties makes it difficult to make a conclusion in this case. However, since this decrease is also observed by Leal-Junior et al. [15], we are confident it is due to plasticization. The large decrease in $E$ for the unannealed mPOFBG during etch can be accounted for with three components:

1) Stress relaxation by polymer molecular rearrangement in the FBG. As discussed earlier, this was not observed to become dominating regarding $\lambda_B$ blue-shift during the short etching process, but this does not mean it was not significant.

2) Tensile stress near the surface will also lead to more orientation along the fiber here and thus high contribution to the total $E$ [50]. This will be etched away during etch and could explain some of the the $E$ lowering.

3) Plasticization. Note, that the unannealed mPOFBG when etched reach the same low $E$ level as the annealed in Fig. 14.

In order to further investigate this, an unannealed mPOFBG was prepared and it’s Young’s modulus $E$ measured. Subsequently, while monitoring $\lambda_B$, it was submerged into distilled water at 65°C, which starts the annealing of residual stress.

Fig. 13. Measured strain sensitivity for mPOFBGs etched in acetone. The error bars denote the standard error of the slope obtained by linear fitting multiplied by a coverage factor of 2.11 taken as the inverse T-distribution at a 95 % confidence level with 17 degrees of freedom, which gives a measure of the repeatability of the measurements.

Fig. 14. Measured Young’s modulus $E$ of etched mPOFs.
causing $\lambda_B$ blue-shift. After observing small wavelength shifts the fiber was removed from the water and allowed to stabilize in room conditions for one hour after which $E$ was measured again. This process was repeated four times, each time increasing the annealing duration up to a total annealing time of approximately 30 minutes. The results are shown in Fig. 15, where the strain sensitivity is given as the fractional value $\frac{\Delta \lambda}{\Delta \varepsilon}$, due to the general scaling of the sensitivity with wavelength, see eq. (2). It is seen, that Young’s modulus is decreasing and strain sensitivity is increasing as the fiber is annealed due to stress relaxation c), however, it is also seen that the $\lambda_B$ shift “required” (and thus the degree of annealing) for Young’s modulus and strain sensitivity to become close to that of the etched unannealed mPOFBG (57 µm diameter) in Fig. 13 and Fig. 14 is approximately 9 nm. Since this large shift is more than compensated during etch, see Fig. 9, the expansion a) of the mPOFBG must be very large. Further, since approximately the same level of $E$ is reached by etching to a diameter of 57 µm as in this annealing, the etched fibers must be almost fully $E$ relaxed long before full stress relaxation, which requires a blue-shift of at least 81 nm as shown in Fig. 7 and by Pasano et al. [19]. Since etching and annealing relax the surface first, this may indicate, that the high $E$ of unannealed fibers originates from and is strongly dependent on the presence of molecules within a distance of ~47 µm to the surface, oriented with length direction frozen-in highly parallel to the fiber axis during drawing, cf. the $E$ influence component 2) above.

The force sensitivity of the mPOFBGs were found as the slope of $(\Delta F, \Delta \lambda_B)$ data, see Fig. 16, showing the force calibration curves of three pre-annealed mPOFBGs. The results of both unannealed and pre-annealed mPOFBGs are shown in Fig. 17 and compared to the calculated values using eq. (2) and fixed values for $E$ and $S_{\varepsilon}$ (2.6 GPa and 0.75 pm/µε, respectively). From this figure, it can be seen that the force sensitivity increases greatly with reduced diameter as expected and that the relatively small changes in strain sensitivity and Young’s modulus introduced by the etching process do not have a big effect. The dominating effect is the reduced cross sectional area of the fiber. These findings are also in good agreement with similar studies [15], [16].

Fig. 15. Measured Young’s modulus (red) and fractional strain sensitivity (blue) for an initially unannealed 150 µm mPOFBG for different degrees of annealing, which is represented by a blue-shift of the Bragg-wavelength.

Fig. 16. Response of mPOFBGs (pre-annealed) towards tensile force.

Fig. 17. Force sensitivity $S_F$ of etched mPOFBGs compared with calculated value $S_F = \frac{S_{\varepsilon}}{E A}$, see eq. (2) - using $S_{\varepsilon} = 0.75$ pm/µε and $E = 2.6$ GPa.

### III. Conclusion

For our PMMA based mPOFBGs etched in acetone we can conclude:

- We have found theoretical, thermal and mechanical evidence for trapped, residual solvent in the polymer matrix after etch: According to HSP theory, solvents good for etching tend to stay inside the polymer. For annealed mPOFBGs, thermal response is only nonlinear during a first temperature rise where the solvent, which act as a plasticizer, is driven out. Due to etch we have found that Young’s modulus is reduced for both unannealed and annealed mPOFBGs to a level below that of the annealed before etch, which indicates plasticization.
- Removal of the acetone plasticizing residuals by the thermal treatment results in a stable linear response with a thermal sensitivity indistinguishable of that of the unetched POFBG.
- Our unannealed mPOFBGs get fully Young’s modulus $E$ relaxed by polymer molecular rearrangement and removal when using acetone etchant. They do probably not get fully stress relaxed.
- The high Young’s modulus $E$ of unannealed fibers is a surface (first ~47 µm) effect. It seems to originate from and be strongly dependent on the presence of molecule chains near the surface with length direction oriented highly parallel to the fiber axis.
• Good solvents for POE etching are characterized by RED $\leq 1$ and relatively high $M_W$.

• Due to the differences in residual stresses, unannealed and annealed mPOFBGs etch very differently. We have demonstrated that these differences are well explained in terms of HSP theory.

• The mPOFBG force sensitivity increases greatly with reduced diameter and the relatively small changes in strain sensitivity and Young’s modulus introduced by the etching process in this case do not have a big effect. The dominating effect in this acetone PMMA mPOFBG etch, is the reduced cross sectional area of the fiber.

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