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Metamaterial Composite with an Ultra-Broadband Usable Range of over 25 Terahertz

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Abstract: Using a metamaterial composite, we demonstrate a bandpass filter that has only a single transmission mode from 0 to >25 THz. This usable bandwidth matches, or exceeds, that of currently used THz sources.

OCIS codes: (160.3918) Metamaterials; (230.7408) Wavelength filtering devices; (300.6495) Spectroscopy, terahertz.

1. Introduction

Since their introduction nearly fifteen years ago, metamaterials have provided a unique design paradigm [1]. However, their use is often limited to a small bandwidth around their design frequency due to the resonant nature of most metamaterials, appearance of unwanted higher order modes, and prevalence of the effective medium approximation where metamaterials are only considered to have a "designed" permittivity and permeability in the deep sub-wavelength regime. While this narrowband behavior can be easily integrated with continuous wave sources and applications, the two most common methods of THz generation, photoconductive antennas and nonlinear generation in either crystals or air plasma, both yield broad spectra that suggest there is an unfulfilled need for broadband THz components.

Here we reconcile the disconnect between traditional broadband sources and narrowband metamaterial components by demonstrating a metamaterial composite with an ultra-broadband usable range. As a proof of principle of this ultra-broadband concept, we have made a series of band pass filters that, to the best of our knowledge, display the largest usable bandwidth of any metamaterial device to date. The filters have a single pass band, with a central frequency scalable from 0.86 – 8.51 THz, while rejecting all other frequency components up to >25 THz.

2. Metamaterial Composite

The filter itself is a metamaterial composite of two distinct components. The first is a metamaterial cross structure, and the second, depending upon your perspective, is either a Babinet complement [2] or a metal mesh [3]. To clearly identify these constituents, we present an optical picture of a 2x2 array of unit cells in Figure 1(a). The unit cell in the bottom right, outlined in black, clearly identifies the metamaterial cross component at the center of the unit cell. However, this unit cell choice suggests that the cross is placed inside of its own, slightly larger, Babinet complement. If the unit cell is translated by (-1/2, 1/2) x P, to the grey outlined unit cell, a different structure is suggested. In this new unit cell, if the cross is ignored, it can be seen that the Babinet complement is also a metal mesh filter. Metal meshes, which have been widely used since their introduction by Ulrich [3], can be described as...
inductive and capacitive meshes. A similar structure was presented previously [4], although we have expanded upon the bandwidth by an order of magnitude.

The sample dimensions - except for $\varepsilon$, which was held constant - were all scaled by a dimensionless parameter, $\sigma$, and can be found in the caption for Figure 1. Ten different samples with varying values of $\sigma$ were fabricated, all on high resistivity Si, spanning the range of $0.86 - 8.51$ THz. Furthermore, the samples were fabricated in single- and double-sided configurations, as shown in Fig 1(b).

3. Results and discussion

The samples were measured with THz-time domain spectroscopy using a two-color air plasma for generation and air biased coherent detection (ABCD), resulting in an ultra-broadband frequency coverage, spanning $0.1 - 30$ THz. The computer simulations were conducted using CST Microwave Studios.

An aggregate comparison between simulation and experiment is shown in Figure 2. The two plots compare the central frequency ($f_0$) and full width at half maximum (FWHM) versus $\sigma$ for both the single- and double-sided samples. Figure 2(a) also has a line that is simply the $\sigma = 1$ central frequency scaled linearly with $\sigma$. If $\varepsilon$ was scaled with $\sigma$, instead of kept constant, every dimension would be scaled identically and this line would predict $f_0$ for all $\sigma$ due to the scale invariance of Maxwell’s equations. However, this would either reduce the filter performance or the fabrication would be practically unfeasible. Fits to the data reveal the actual scaling versus $\sigma$: $f_0 = 8.223\sigma^{-1.424} + 0.2774$; $\Delta f_1 = 5.511\sigma^{-1.557} + 0.2178$; $\Delta f_2 = 3.161\sigma^{-1.464} - 0.2098$. $\Delta f_1$ and $\Delta f_2$ are the FWHM for the single- and double-sided samples, respectively.

The effect of the single- versus double-sided structures can clearly be seen in Figure 2(b). The peak transmission through the bandpass structures is roughly equivalent to the Fresnel transmission coefficient for the air HR-Si interface. Subsequently, the additional layer on the backside of the Si wafer further reduces the transmitted bandwidth without significantly affecting its peak performance.

![Figure 2](STu2F.1.pdf)

Figure 2: (a) The aggregate agreement between $f_0$ in simulation and experiment. To highlight the effect of the constant minimum linewidth in all structures, the solid black line represents a linear scaling of $f_0$ vs $\sigma$. (b) The comparison of the FWHM for experimental and simulated results.

4. Conclusion

In conclusion, we have presented a metamaterial composite with an ultra-broadband usable bandwidth suitable for virtually any THz source. We have constructed a series of band pass filters that clearly demonstrate this concept, and provided simple equations that can be used to construct frequency filters from $0.86 - 8.51$ THz without need for simulation or design. It is our hope that this work will bring an easy to fabricate, functional THz component to the laboratory, and expand the reach of metamaterial based THz components towards broadband functional components.

5. References