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Published in:
Optics Express

Link to article, DOI:
10.1364/OE.22.011267

Publication date:
2014

Document Version
Publisher's PDF; also known as Version of record

Citation (APA):
Hansen, K. R., & Lægsgaard, J. (2014). Impact of gain saturation on the mode instability threshold in high-power fiber amplifiers. Optics Express, 22(9), 11267-11278. DOI: 10.1364/OE.22.011267
Impact of gain saturation on the mode instability threshold in high-power fiber amplifiers

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Abstract: We present a coupled-mode model of transverse mode instability in high-power fiber amplifiers, which takes the effect of gain saturation into account. The model provides simple semi-analytical formulas for the mode instability threshold, which are valid also for highly saturated amplifiers. The model is compared to recently published detailed numerical simulations of mode instability, and we find reasonably good agreement with our simplified coupled-mode model.

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OCIS codes: (060.2320) Fiber optics amplifiers and oscillators; (060.4370) Nonlinear optics, fibers; (140.6810) Thermal effects; (190.3100) Instabilities and chaos; (190.4370) Nonlinear optics, fibers.

References and links


1. Introduction

The recent discovery of transverse mode instability (TMI) in high-power fiber amplifiers [1] has motivated experimental [2–4] and theoretical [5–13] efforts to characterize and explain this phenomenon. TMI manifests itself as an instability in which the signal power fluctuates on a ms timescale between the fundamental mode (FM) and the first few higher-order modes (HOMs) when the average signal power is increased above a certain threshold. The experimental and theoretical research into this phenomenon has provided strong evidence that the underlying physical mechanism responsible for TMI is a form of stimulated thermal Rayleigh scattering (STRS) [7], caused by quantum defect heating of the active fiber. Essentially, the presence of light in the FM and HOM of a multimode active fiber results in a mode beating pattern in the signal intensity, which leads to a corresponding spatially and temporally oscillating thermal profile. Since the refractive index of silica is temperature dependent, a thermally induced long period grating (LPG) is generated, which has the correct period to couple the two modes. Due to the fact that this thermal nonlinearity is extremely slow, with a typical timescale of several ms, a phase-lag between this LPG and the mode beating pattern will exist if the latter changes due to noise or transients in the signal, which is necessary to induce strong coupling between the modes [6].

Several numerical models of TMI have been based on variants of the beam propagation method (BPM) [5–10]. While such numerical simulations can potentially include all relevant physical effects, the need to include the temporal dynamics, either by performing the simulation in the time domain or by including several frequency components in the frequency domain, leads to a numerically very intensive problem. Indeed, these numerical codes have so far employed various approximations, such as limiting the number of frequency components in the simulation [6, 7], expanding the electric field in a small number of transverse modes [8, 10] or limiting the simulation time in time domain algorithms [9].

In previous publications, we have proposed an alternative model of TMI in terms of a coupled-mode model [11, 12]. While this model has the advantage of numerical efficiency, and even allows for semi-analytical solutions, it does rely on a number of approximations, the most limiting of which is the neglect of spatiotemporal oscillations in the gain coefficient due to gain saturation. This approximation leads to an overestimate of the strength of the thermally induced LPG, which in turn leads to an underestimate of the TMI threshold power. In this paper, we show how this important effect can be included to lowest order in a coupled-mode formulation of TMI in fiber amplifiers.

2. Coupled-mode model

We consider the signal field in the scalar approximation written as a superposition of the FM and first HOM of the fiber, since typically only the first HOM is involved in TMI [3, 12]

\[ E(r, t) = a_1(z)\psi_1(r_\perp)e^{-i(\beta_1z-\omega_1t)} + a_2(z)\psi_2(r_\perp)e^{-i(\beta_2z-\omega_2t)}. \] (1)

Here, \( a_{1,2} \) are the amplitudes of the FM and HOM, respectively, with normalized mode profiles \( \psi_{1,2} \) and propagation constants \( \beta_{1,2} \). Note that the frequency of the light in the HOM \( \omega_2 \) differs from the light in the FM \( \omega_1 \). We consider here that the frequency shifted light in the HOM is due to either intensity or quantum noise, and we shall later integrate over \( \omega_2 \) to obtain the total power in the HOM. The electric field obeys the scalar wave equation

\[ \frac{\partial^2 E}{\partial z^2} + \nabla_z^2 E - \frac{v(r, t)}{c^2} \frac{\partial^2 E}{\partial t^2} = 0, \] (2)
where the relative permittivity is given by
\[ \varepsilon(r, t) = \varepsilon_f(r, t) + i \frac{g(r) \sqrt{\varepsilon_f}}{k} + \Delta \varepsilon(r, t). \] (3)

The relative permittivity of the unperturbed fiber is given by \( \varepsilon_f \), the perturbation due to the thermal nonlinearity is given by \( \Delta \varepsilon = \eta \Delta T \) and the gain coefficient is denoted by \( g \).

The change in temperature relative to the cooling fluid \( \Delta T \) obeys the heat equation
\[
\rho C \frac{\partial \Delta T}{\partial t} - \kappa \nabla^2 \Delta T(r, t) = Q(r, t),
\] (4)

subject to the convection cooling boundary condition
\[
\kappa \frac{\partial \Delta T}{\partial r} + h_q \Delta T(r, t) = 0,
\] (5)

where \( \rho \) is the density, \( C \) is the specific heat capacity, and \( \kappa \) is the thermal conductivity of the fiber material. These are assumed to be constant throughout the fiber cross section and independent of temperature. The convection coefficient of the coolant is denoted by \( h_q \). Note that we have assumed that the longitudinal heat diffusion is negligible compared to the transverse heat diffusion in the heat equation. This approximation was tested numerically in [8], where it was found that TMI is not strongly dependent on the longitudinal heat diffusion. In fibers with a mode beat length comparable to the core radius, this approximation may break down and longitudinal heat diffusion would need to be taken into account. Assuming a steady-state solution to the rate equations of the gain medium, the heat load is given by
\[
Q(r, t) = \left( \frac{\lambda_s}{\lambda_p} - 1 \right) g(I) I(r, t),
\] (6)

where \( \lambda_s \) is the signal wavelength, \( \lambda_p \) is the pump wavelength, \( I \) is the signal intensity and the gain coefficient \( g \) is given by
\[
g(I) = \frac{g_0}{1 + I/I_{sat}},
\] (7)

where \( g_0 \) is the small signal gain and \( I_{sat} \) is the saturation intensity, both of which are dependent on the pump power. The intensity of the signal can be written
\[
I(r, t) = I_0(r) + I(r, t),
\] (8)

where \( I_0 \) is a slowly varying component given by
\[
I_0(r) = |p_1(z)|^2 |\psi_1(r, t)|^2 + |p_2(z)|^2 |\psi_2(r, t)|^2,
\] (9)

and \( I \) is a spatially and temporally oscillating component given by
\[
I(r, t) = p_1(z) p_2(z)^* \psi_1(r, t) \psi_2(r, t)^* e^{-i(\Delta \beta z - \Delta \omega t)} + \text{c.c.}
\] (10)

We have introduced the difference in propagation constant \( \Delta \beta = \beta_1 - \beta_2 \) and frequency \( \Delta \omega = \omega_1 - \omega_2 \), as well as the scaled mode amplitudes
\[
p_i(z) = \sqrt{\frac{1}{2} n_c e_0 c a_i(z)},
\] (11)

where \( n_c \) is the core refractive index of the fiber. It is clear from Eq. (6) that the presence of light in the HOM leads to an oscillating heat load \( Q \), which results in a thermally induced LPG.
However, it is also clear from Eq. (7) that the strength of this grating is reduced if the gain is highly saturated, since the maxima in the mode beating intensity pattern coincide with minima in the gain coefficient under the assumption of fixed $g_0$ and $I_{sat}$. In the extreme case where $I >> I_{sat}$, we see that the oscillations in the heat load disappear, and no LPG is induced.

Assuming that the fiber amplifier is operating below the TMI threshold, where the excitation of the HOM is small compared to the FM, we can expand Eq. (6) in powers of $I/I_0 + I_{sat}$ which yields

$$Q(r, t) \approx \left( \frac{\lambda_s}{\lambda_p} - 1 \right) \frac{g_0}{1 + I_0/I_{sat}} \left[ I_0 + I_{sat} \left( \frac{I}{I_0 + I_{sat}} - \left( \frac{I}{I_0 + I_{sat}} \right)^2 \right) \right].$$  \tag{12}$$

Provided that the HOM excitation is weak compared to the FM excitation, which is the case when the amplifier is not operating above the TMI threshold power, we can neglect higher order terms in the expansion and keep only the first order term. The approximate heat load is thus given to first order as

$$Q(r, t) \approx \left( \frac{\lambda_s}{\lambda_p} - 1 \right) \frac{g_0}{1 + I_0/I_{sat}} \left( I_0 + \frac{I}{I_0 + I_{sat}} \right).$$  \tag{13}$$

A similar expansion of the heat load was considered in [10] to show that TMI does not occur under steady-state conditions, as also pointed out in [6].

As in [11] we can now derive coupled-mode equations, which are valid also for non-steady-state conditions, for the mode amplitudes $p_i$ by solving Eq. (4) in the frequency domain using an appropriate Green’s function

$$\Delta T(r, \omega) = \frac{1}{\kappa} \int \int G(r_\perp, r'_\perp, \omega) Q(r'_\perp, \omega) d^2 r'_\perp.$$  \tag{14}$$

In this paper, we shall consider only step-index fiber (SIF) models, and we assume that the thermal conductivity $\kappa$ and thermo-optic coefficient $\eta$ are constant throughout the fiber cross section. In this case, the appropriate Green’s function is the same as given in [12]. The resulting coupled-mode equations in terms of mode power $P_i = |p_i|^2$ are

$$\frac{\partial P_1}{\partial z} = -g_0(z) \chi(z, \Delta \omega) P_2 P_1 + \Gamma_1 \bar{g}(z) P_1,$$  \tag{15a}$$

$$\frac{\partial P_2}{\partial z} = g_0(z) \chi(z, \Delta \omega) P_2 P_1 + \Gamma_2 \bar{g}(z) P_2.$$  \tag{15b}$$

We have introduced an approximate gain coefficient $\bar{g}$ in the last term on the rhs. of Eqns. (15), in which spatial hole burning is only approximately taken into account. It is given by

$$\bar{g} = \frac{g_0}{1 + \Gamma_1 P_1/P_{sat}},$$  \tag{16}$$

where $P_{sat} = I_{sat} A_d$ is the saturation power, $A_d$ is the area of the doped region, and $\Gamma_i$ are overlap integrals with the doped region given by

$$\Gamma_i = \int_{A_d} |\psi_i(r_\perp)|^2 d^2 r_\perp.$$  \tag{17}$$

This approximation is valid provided that the FM profile does not vary much over the doped region of the fiber core, and is a good approximation for weakly-guiding, nearly single-mode fibers.
The nonlinear gain coefficient $\chi$ is given by

$$\chi(z, \Delta\omega) = \eta k^2 \left( 1 - \frac{\lambda_p}{\lambda_p} \right) \text{Im} [A(z, \Delta\omega)], \quad (18)$$

where $k$ is the vacuum wavenumber of the signal. The quantity $A$ is given by the integral

$$A(z, \Delta\omega) = \int \psi_1(\mathbf{r}_\perp) \psi_2(\mathbf{r}_\perp) \int_{A_d} G(\mathbf{r}_\perp, \mathbf{r}'_\perp, \Delta\omega) \frac{\psi_1(\mathbf{r}'_\perp) \psi_2(\mathbf{r}'_\perp)^*}{(1 + P(z)|\psi_1(\mathbf{r}'_\perp)|^2/\text{Isat})^2} d^2\mathbf{r}'_\perp d^2\mathbf{r}_\perp, \quad (19)$$

which can be calculated using standard numerical quadrature methods for any fiber for which the thermal Green’s function and mode functions are known. It is clear from Eq. (19) that the nonlinear gain coefficient $\chi$ is reduced as the signal power becomes comparable to the saturation power. This is due to the above mentioned reduction in the strength of the thermally induced LPG. While it is possible to calculate the integral in Eq. (19) for each $z$ in a numerical solution of the coupled-mode equations, we shall show that a more efficient approach is to approximate the inner integral in Eq. (19) so that the integral only need to be calculated once. We assume that the mode function of the FM $\psi_1$ in the denominator of the integrand can be approximated to be constant over the doped region, such that

$$|\psi_1(\mathbf{r}_\perp)|^2 \approx \frac{1}{A_d} \int \int_{A_d} |\psi_1(\mathbf{r}_\perp)|^2 d^2\mathbf{r}_\perp$$

in the denominator of Eq. (19). We hereby obtain the approximate expression

$$A(z, \Delta\omega) \approx \left( \frac{\tilde{g}(z)}{\tilde{g}_0(z)} \right)^2 \int \psi_1(\mathbf{r}_\perp) \psi_2(\mathbf{r}_\perp) \int_{A_d} G(\mathbf{r}_\perp, \mathbf{r}'_\perp, \Delta\omega) \psi_1(\mathbf{r}'_\perp) \psi_2(\mathbf{r}'_\perp)^* d^2\mathbf{r}'_\perp d^2\mathbf{r}_\perp. \quad (21)$$

With this approximation, the coupled-mode equations become

$$\frac{\partial P_1}{\partial z} = -\tilde{g}(z) \tilde{\chi}(z, \Delta\omega) P_2 P_1 + \Gamma_1 \tilde{g}(z) P_1, \quad (22a)$$

$$\frac{\partial P_2}{\partial z} = \tilde{g}(z) \tilde{\chi}(z, \Delta\omega) P_2 P_1 + \Gamma_2 \tilde{g}(z) P_2, \quad (22b)$$

where the approximate nonlinear gain coefficient $\tilde{\chi}$ is given by

$$\tilde{\chi}(z, \Delta\omega) = \frac{\tilde{g}(z)}{\tilde{g}_0(z)} \chi_0(\Delta\omega) \quad (23)$$

and $\chi_0$ is the same nonlinear gain coefficient derived in [11, 12]

$$\chi_0 = \frac{\eta k^2}{\kappa \beta_1} \left( 1 - \frac{\lambda_p}{\lambda_p} \right) \text{Im} [A_0(z, \Delta\omega)], \quad (24)$$

with

$$A_0(\Delta\omega) = \int \psi_1(\mathbf{r}_\perp) \psi_2(\mathbf{r}_\perp) \int_{A_d} G(\mathbf{r}_\perp, \mathbf{r}'_\perp, \Delta\omega) \psi_1(\mathbf{r}'_\perp) \psi_2(\mathbf{r}'_\perp)^* d^2\mathbf{r}'_\perp d^2\mathbf{r}_\perp. \quad (25)$$

We note that the approximate coupled-mode equations given by Eqns. (22) are similar to the ones derived in [11], except for the fact that the nonlinear gain coefficient is now $z$-dependent and has a reduced value due to the effect of gain saturation.

In order to check the approximation used in obtaining the approximate expression for $A$ given in Eq. (21), we compare the value of $\chi$ calculated using the exact expression in Eq. (19) and the
Table 1. Fiber parameters for Fiber A and B.

<table>
<thead>
<tr>
<th>Parameter</th>
<th>Value</th>
</tr>
</thead>
<tbody>
<tr>
<td>Core radius $R_{core}$</td>
<td>20 μm</td>
</tr>
<tr>
<td>Cladding radius $R_{clad}$</td>
<td>100 μm</td>
</tr>
<tr>
<td>Doped radius $R_{doped}$</td>
<td>20 μm / 15 μm</td>
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<td>Outer radius $R_{outer}$</td>
<td>500 μm</td>
</tr>
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<td>Core index $n_{core}$</td>
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<tr>
<td>V parameter $V$</td>
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<tr>
<td>Dopant concentration $\rho_{Yb}$</td>
<td>$3 \times 10^7 \text{ μm}^{-3}$</td>
</tr>
<tr>
<td>Fiber length $L$</td>
<td>1.5 m / 3 m</td>
</tr>
<tr>
<td>Signal wavelength $\lambda_s$</td>
<td>1032 nm</td>
</tr>
<tr>
<td>Pump wavelength $\lambda_p$</td>
<td>975 nm</td>
</tr>
<tr>
<td>Convection coefficient $h_q$</td>
<td>1000 W/(m²K)</td>
</tr>
<tr>
<td>Thermo-optic coefficient $\eta$</td>
<td>$3.5 \times 10^{-5} \text{ K}^{-1}$</td>
</tr>
<tr>
<td>Thermal conductivity $\kappa$</td>
<td>1.4 W/(Km)</td>
</tr>
<tr>
<td>Volumetric heat capacity $\rho C$</td>
<td>$1.67 \times 10^6 \text{ J/(Km³)}$</td>
</tr>
</tbody>
</table>

Fig. 1. Nonlinear gain coefficient $\chi$ as a function of $P/P_{sat}$ for a fixed frequency difference $\Delta \nu = 1 \text{ kHz}$. The solid curves are calculated using the exact expression for $A$, while the dashed curves are calculated using the approximate expression for $A$. The red curves are the results for Fiber A with a doped radius of 20 μm, while the blue curves are for Fiber B with a doped radius of 15 μm. As expected, the nonlinear gain coefficient decreases with increasing signal power due to the effect of gain saturation. The approximate calculation of $\chi$ remains reasonable as the signal power is increased, especially for Fiber B with the reduced doped area.

approximate expression given in Eq. (21) for a typical double-clad SIF and a range of values of $P_1/P_{sat}$. Two SIFs, called Fiber A and Fiber B, with different radii of the doped region are considered to investigate the effect of a reduced doped region. The relevant fiber parameters are given in Table 1 and the result is shown in Fig. 1. It is clear that the nonlinear gain coefficient decreases significantly as the signal power becomes comparable to the saturation power, which
where $S_{Acl}$ and counter-pumped operation, respectively. The pump gain coefficient $\tilde{\chi}$, where $
abla \Omega$ gives the following expression for the output HOM content

$$P_2(L) = \int_{-\infty}^{\infty} S_2(0, \Delta \omega) \left( \frac{P_1(L)}{P_1(0)} \right)^{\Gamma_2/\Gamma_1} \exp \left( \frac{\chi_0(\Delta \omega)}{1} \right) \int_0^L \frac{\tilde{g}^2}{\tilde{g}_0} P_1 dz,$$

where $S_2(0, \Delta \omega)$ is the input power spectral density of the light in the HOM and $L$ is the fiber length. It was argued in [12] that a likely source for seeding the TMI is intensity noise in the input signal. In this case the input power spectral density is

$$S_2(0, \Delta \omega) = P_2(0) \delta(\Delta \omega) + \frac{1}{4} R_N(\Delta \omega) P_2(0),$$

which gives the following expression for the output HOM content $\xi = P_2(L)/[P_1(L) + P_2(L)]$

$$\xi(L) = \xi(0) \left( \frac{P_1(L)}{P_1(0)} \right)^{1-\Gamma_2/\Gamma_1} \left[ 1 + \frac{R_N(\Omega)}{4} \sqrt{\frac{2\pi}{\chi_0''(\Omega)M}} \exp \left( \chi_0(\Omega)M \right) \right],$$

where $\Omega$ is the value of $\Delta \omega$ for which $\chi_0$ has its maximum value, $\chi_0''$ is the second derivative of $\chi_0$ with respect to $\Delta \omega$ and $M$ is given by the integral

$$M = \int_0^L \frac{\tilde{g}^2}{\tilde{g}_0} P_1 dz.$$

This integral is evaluated by first solving Eqns. (26a) and (26b) numerically using a standard numerical ODE solver followed by numerical quadrature. Likewise, the integral in Eq. (25)
required for calculating \( \chi_0 \) is evaluated for a range of frequencies using standard numerical quadrature techniques.

We calculated the HOM content \( \xi(L) \) as a function of output signal power for the fiber given in Table 1. The result is shown in Fig. 2 for both co- and counter-pumped operation, and is also compared to the result obtained when the effect of the gain oscillations are neglected as in [11, 12]. The input signal power was taken to be 10 W and the input HOM content \( \xi(0) = 0.01 \), which should correspond to nearly optimal coupling of the signal into the fiber core. The relative intensity noise was assumed to be -100 dBc/Hz, which we believe is representative of realistic experimental conditions. Defining the TMI threshold power \( P_{th} \) as the output signal power at which \( \xi(L) = 0.1 \), we find \( P_{th} = 723 \) W for counter-pumped operation and \( P_{th} = 672 \) W for co-pumped operation. Compared to the corresponding value of \( P_{th} = 260 \) W obtained when gain oscillations are neglected, it is clear that this effect is very important in this case. Interestingly, there is little difference in TMI threshold between the co- and counter-pumped cases. When gain saturation is low and the effect of the associated gain oscillations can be ignored, this is expected, as shown in [11]. Apparently, including the effects of gain saturation in the model approximately preserves the symmetry between co- and counter-pumping.

We have also performed the same calculation for Fiber B. The result is shown in Fig. 3. The TMI thresholds for co- and counter-pumped operation are 1047 W and 1286 W, respectively, which is substantially higher than for Fiber A. This is due to the combination of the reduced overlap between the modes with the doped region and the longer fiber length. From the dashed curve in Fig. 3 we see that the TMI threshold, calculated under the assumption of negligible gain oscillations, is 331 W for Fiber B. This increase of 27% compared to Fiber A is due to the reduced mode overlap with the doped region. However, the increase in TMI threshold when taking gain oscillations into account is much larger, 78% for counter-pumped and 56% for co-pumped. This increase in TMI threshold thus cannot solely be explained by the reduced mode overlap, but is also caused by the higher gain saturation. We have verified this conclusion by simulating a shorter version of Fiber B, which has the same length as Fiber A. In this case, the
TMI threshold for co-pumped operation drops to 767 W compared to 1047 W for the original Fiber B. This is due to the fact that the gain is much less saturated in this case, which in turn implies that the gain oscillations are weaker and the thermally induced LPG is stronger. However, this shorter version of Fiber B has a very poor efficiency due to the lower pump absorption, and is thus not considered a realistic case.

Fig. 3. HOM content $\xi$ as a function of signal output power $P_s$ for Fiber B for counter-pumped (blue curve) and co-pumped (red curve) operation. The dashed curve shows the result when the gain oscillations are neglected as in [11, 12].

4. Comparison to BPM calculations by Smith and Smith

In a recent paper [13], Smith and Smith investigated the effect of gain saturation on TMI numerically using a BPM code. In this section we compare our calculation of the nonlinear gain coefficient $\chi$ with a calculation of the same quantity obtained by the BPM presented in Fig. 5 in [13]. The fiber parameters for the two SIFs used for the comparison are given in Table 2. The first fiber, here called Fiber C, is a short fiber with a small pump cladding, which provides high pump intensity and hence lower gain saturation, while the second fiber, called Fiber D, is a longer fiber with a much larger pump cladding, which therefore operates with much higher gain saturation. We have calculated the nonlinear gain coefficient $\chi$ using the approximate expression given in Eq. (21) as a function of $z$ for a fixed frequency offset $\Delta \nu = 1.1$ kHz and the specific pump powers given in [13]. The result for Fiber C is shown in Fig. 4, while the result for Fiber D is shown in Fig. 5. Compared to Fig. 5 in [13] we find a reasonably good agreement with our results, but with some interesting differences. For Fiber C, we find an initially lower value of $\chi$ for the counter-pumped case, which rises to a slightly lower value than in [13]. For the co-pumped case, $\chi$ falls off more quickly with $z$ compared to [13] and also reaches a lower value. The same is true for the co-pumped case for Fiber D, while the counter-pumped case shows an initial value of $\chi$ slightly higher than in [13], which then falls off with $z$ to a slightly lower value. Interestingly, the counter-pumped case in the lower plot of Fig. 5 in [13] shows a slightly increasing nonlinear gain $\chi$ with $z$, which according to our analysis would seem to indicate a decreasing gain saturation with $z$. However from Fig. 2 in [13], it appears that the
Table 2. Fiber parameters for Fiber C and D used for comparison with [13].

<table>
<thead>
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<tr>
<td>Pump wavelength</td>
<td>$\lambda_p$</td>
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Gain saturation is in fact increasing with $z$. A possible explanation of this apparent discrepancy could be thermal lensing, since the larger heat load toward the output end of the fiber will cause the overlap between the HOM and FM to increase, which will increase the nonlinear gain.

To test the accuracy of the approximate calculation of $\chi$ for Fiber C and D, we again compare the exact and approximate expressions for a fixed frequency offset $\Delta \nu = 1.1$ kHz and for a range of values of $P_1/P_{sat}$. The result is shown in Fig. 6 and shows good agreement between the exact and approximate calculations even when gain saturation is substantial. The discrepancy between the calculation in [13] and our model is thus unlikely to be due to this approximation, but is rather due to other approximations, such as the fact that we do not include thermal lensing, and the fact that we treat the gain oscillations to first order only.

Fig. 4. Nonlinear gain coefficient $\chi$ as a function of $z$ for Fiber C. The blue curve is for counter-pumped operation with a pump power of 493 W, and the red curve is for co-pumped operation with a pump power of 525 W. The dashed curve shows the value of $\chi_0$. 

Received 24 Jan 2014; revised 4 Apr 2014; accepted 25 Apr 2014; published 2 May 2014

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5 May 2014 | Vol. 22, No. 9 | DOI:10.1364/OE.22.011267 | OPTICS EXPRESS  11276
5. Conclusion

In summary we have shown how the effects of gain saturation can be included approximately in a coupled-mode model of TMI. The resulting model is numerically very efficient and easy
to implement, and can be used to obtain quantitative estimates of the TMI threshold for a
given fiber amplifier under conditions where gain saturation is substantial. Applying our model
to particular SIF designs (Fiber A & B), we found that the estimate for the TMI threshold
increased by approximately a factor of 3 compared to a model which does not include the effect
of gain saturation. Only in the case of a short, rod-type fiber (Fiber C) did we find a nonlinear
gain comparable to the prediction of the model when neglecting gain saturation effects. We
also found that the benefit of reducing the doped area of the core is significantly greater than
predicted by the simpler model, due to the increased gain saturation. Including gain saturation
in the model breaks the symmetry between co- and counter-pumping, with the latter having the
higher TMI threshold.

Comparing our model to BPM simulations presented in [13], we found reasonably good
agreement between the nonlinear gain coefficient calculated with our model and the BPM re-
sults, although our model results in slightly lower values for the nonlinear gain coefficient. This
is likely due to the various approximations employed by our model, most importantly the ne-
glect of thermal lensing and the fact that gain saturation is only included to lowest order of
approximation. We thus expect our model to be able to provide good quantitative estimates of
the TMI thresholds of fiber amplifier designs under conditions where gain saturation is signifi-
cant.