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Fraunhofer regime of operation for superconducting quantum interference filters


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Series arrays of superconducting quantum interference devices (SQUIDs) with incommensurate loop areas, so-called superconducting quantum interference filters (SQIFs), are investigated in the kilohertz and the gigahertz frequency range. In SQIFs made of high-\textit{T}_{c} bicrystal junctions the flux-to-voltage response \(\partial V/\partial \Phi\) is dominated by the variation in the critical current in the individual junctions (Fraunhofer-type) rather than by the SQUIDs interference. For a SQIF with 20 SQUID loops we find \(\partial V/\partial \Phi=40\,\text{mV}/\Phi_0\) and a dynamic range of more than 60 dB in the kilohertz range. In the 1–2 GHz range the estimated power gain is 20 dB and the magnetic flux noise level is as low as \(10^{-4}\Phi_0\). © 2008 American Institute of Physics. [DOI: 10.1063/1.3058759]

Microwave amplifiers based on superconducting quantum interference effects in two parallel connected Josephson junctions (JJs) (SQUID) (SQUID denotes superconducting quantum interference device) are characterized by a noise temperature close to the quantum limit (see, e.g., Refs. 1 and 2). SQUID amplifiers, as well as other Josephson devices without feedback, possess low saturation power, which is proportional to the characteristic voltage \(V_0=I_cR_n\), where \(I_c\) is the critical current and \(R_n\) is the normal state resistance of the JJ. For resistively shunted tunnel JJs made of low-\(T_c\) superconductors (LTS), \(V_0\) (here \(R_n\) is the shunt resistance needed to avoid hysteresis) only reaches 200–300 \(\mu\text{V}\) at \(T=4.2\,\text{K}\) (Ref. 2), while the bicrystal JJs made of high-\(T_c\) superconductors (HTS) can give \(V_0=1\,\text{mV}\) at \(T=77\,\text{K}\). 3

The output signal and saturation power can be increased by using an array of JJs or SQUIDs. An increase in the output signal proportional to the number \(N\) of SQUIDs was demonstrated in an amplifier based on a series-connected array of LTS SQUIDs. 4 Generally, however, the spread in \(I_c\) and \(R_n\) parameters of the HTS JJs is a pertinent problem that restricts the use of series-connected JJs or SQUID arrays. Superconducting quantum interference filters (SQIFs)—arrays of SQUIDs (series or parallel connected) with incommensurate SQUID-loop areas—accept much wider margins in parameter spread. 5, 7 For a SQIF with small normalized array loop inductances \(L_i<1\) (where \(L_i=2\pi L_i/\Phi_0\), \(\Phi_0=h/2e\) is the magnetic flux quantum, \(I_c\) is the critical current, and \(L_i\) is the SQUID-loop inductance) and a suitably chosen distribution of loop sizes, the magnetic field to voltage response \(V(H)\) is a nonperiodic function with a single narrow global minimum at \(H=0\). The contributions from the different SQUID loops are washed out, and the width of the \(V(H)\) minimum is determined by the number of loops and the size of the largest SQUID loop. 7 For short we will call this operational regime of the SQIF with a single \(V(H)\) minimum at \(H=0\) the S-mode. Note that in the S-mode also the influence of the Fraunhofer-like \(I_c(H)\) dependence in the individual JJs is visible but may be neglected. Recently, HTS SQIFs have been studied for microwave amplification. 8, 9

Sensitive magnetometers based on one-dimensional series arrays of HTS bicrystal JJs have been reported. 10 Here both the Fraunhofer dependence of the critical current in the JJs, \(I_c(H)=I_c(0)|\sin(\pi \Phi_j/\Phi_0)|/(\pi \Phi_j/\Phi_0)|\) (where \(\Phi_j=\mu_0 H \alpha_j\) is the magnetic flux in a JJ with effective area \(\alpha_j\)), and the flux focusing due to the geometry of the bicrystal JJ have to be taken into account. We call this operational regime the \(F\)-mode. A field-voltage transfer function \(\partial V/\partial \mu_0 H=7500\,\text{V/T}\) (\(\mu_0\)−vacuum permeability) was obtained for an array with \(N=105\) series-connected JJs. 10 This compares well with the sensitivity of a LTS SQUID array with small spread of parameters. 4

In this paper we investigate series-connected SQIF consisting of 20 HTS bicrystal SQUIDs with SQUID-loop inductances in the range of 15–300 \(\mu\text{H}\). The width (\(w\) \(=10\,\mu\text{m}\)) of the film forming both the JJs and the SQUID loops is much larger than the London penetration depth \(\lambda_L\) in \(\text{YBa}_2\text{Cu}_3\text{O}_7\) (YBCO). The external magnetic flux produced by the input signal coil induces circulating currents in the superconducting SQUID loops. Since \(w\gg\lambda_L\), the current circulating along the inside edge of the SQUID loop is compensated by the current in the outer edge of the SQUID loop. The presence of the circulating current in the JJs provides better inductive coupling between the JJs and the input circuit and in turn increases the microwave amplification.

Bicrystal \(\text{NdGaO}_3\) (NGO) substrates were used in the sample fabrication. NGO is characterized by a tolerable permittivity (\(\varepsilon=25\)) and fairly low microwave losses (\(tg\delta\) \(<10^{-3}\)). Devices were formed by ion-plasma and chemical etching of YBCO film deposited by dc sputtering at high oxygen pressure. The fabrication details of bicrystal JJs have been described elsewhere. 3 Single loop SQUID and series-connected SQUID arrays 8 were fabricated for comparison. Figure 1 shows the topology of a SQIF designed for use as a microwave amplifier. The input line (Au film) was deposited over a \(\text{SiO}_2\) insulator layer; the bottom layer is the YBCO film that forms the SQIF located inside the input line. The

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output voltage signal is taken directly from the SQIF.

We measured dc J-V and V(H) curves at low frequency as well as the output noise power $P_n=k_B T_n \Delta f$ in the frequency range $f=1-2$ GHz. For the noise power measurements we used a cryogenic preamplifier with noise temperature $T_{A1}= (8 \pm 2)$ K and gain $G=21$ dB, followed by a room temperature amplifier with $T_{A2}=130$ K and $G=40$ dB. The output voltage signal was simultaneously recorded by a spectrum analyzer and a quadratic detector integrated into the room amplifier.

From the $I_c(H)$ pattern measured on the reference SQUID with loop area $S=35$ $\mu m^2$ we estimate a loop inductance $L=(15 \pm 5)$ pH. The estimated spread in critical currents of JJs in the SQIF was $\delta I_c/I_c = 30\%$. Using the mean value of $I_c=100$ $\mu A$ for JJs with width $w=10$ $\mu m$, we calculate the normalized inductances of the loops in the SQIF to be in the range $I_c=4-90$.

Numerical simulations (using the PSCAN program) of the $V(H)$ curve for SQIFs with different numbers of loops $N$ are shown in Fig. 2. Experimental values of loop inductances as well as the Fraunhofer dependence of the JJ critical currents are taken into account. This leads to an aperiodic $V(H)$ response even for a single SQUID. The modulation of the side lobes is suppressed with increasing $N$. Dips due to both Fraunhofer modulations in junctions and SQUID interference can be seen for $N>10$.

Figure 3 shows a family of $V(H)$ responses of a SQIF with 20 loops plotted at different dc bias currents $I_b$. A single $F$-mode dip with small side voltage modulation is seen for all $I_b$. The inset shows a zoom-in of the central part. Besides the single $S$-mode dip predicted by numerical simulations (Fig. 2), additional dips produced by the SQUIDs response interference are observed. The width of the $V(H)$ dip is considerably larger than the dip in a SQIF operated in the $S$-mode. The width of the $F$-mode dip is well fitted using an effective JJ area $S_{eff}=30$ $\mu m^2$. This is larger than the evaluated area of JJ $W_A=1.5$ $\mu m^2$, where $w$ is the JJ width and $\lambda_L=0.15$ $\mu m$ is the London penetration depth in YBCO. This observation is in agreement with a strong flux focusing effect in a bicrystal JJ (Ref. 13) when the magnetic field is applied perpendicular to the film. The expected value of the $V(H)$ response, $V_S=\Sigma V_{0i}$, should increase with the number $N$ of SQIF loops, each contributing $V_{0i}=I_c R_{ni}$. However, the experimental $V_S=12$ mV turned out smaller than the estimated $\Sigma V_{0i}=20$ mV for a SQIF with $N=20$. For the SQIF
operating in the F-mode we obtain $dV/d\Phi=40$ mV/$\Phi_0$ (and $dV/d\mu_0H=270$ V/T), while for the reference SQUID $dV/d\Phi=1$ mV/$\Phi_0$.\(^5\) Thus for a SQIF made of HTS bicrystal JJs, the F-mode plays a predominant role in the V(H) response.

The output noise temperature $T_n(H)$ measured in the frequency range $f=1–2$ GHz for a 20 loop SQIF is presented in Fig. 4. The output noise power is up to 15 dB above the background noise level. The measured output noise signal from the SQIF can be interpreted as an incoherent superposition of voltage spectral density of thermal fluctuations $S'_V = \Sigma(8kTR_d/R_0)[1+1/(2J_0^2)]^2$ and voltage spectral density corresponding to contribution of magnetic flux conversion $S''_V = \Sigma(2kTL_d/R_0)(dV_i/d\Phi_0)^2$.\(^6\) For a bias current $I_B=1.25I_c$ and $R_2=30\ \Omega$, at the $T_n(H)$ peak we get $S''_V/S'_V$ proportional to max($I_c$/N)>1 [neglecting $R_2(H)$ and $I_c$ variation in the SQIF]. It indicates that the observed output noise signal is dominated by the magnetic flux conversion $S''_V$ rather than the thermal fluctuations $S'_V$.

With $dV/d\Phi=40$ mV/$\Phi_0$, the SQIF is characterized by a flux sensitivity $d\Phi=10^{-3}$ $\Phi_0$. Both the $dV/dH(H)$ and the $T_n(H)$ functions have complex dependences with several maxima. However, some of the $T_n(H)$ peaks are correlated with the $dV/dH$ function. The side lobe modulation of the $dV/dH$ [weakly seen in the $V(H)$ response] and the corresponding behavior of the $T_n(H)$ function may be caused by residual contributions from the SQUID-loop interference that could be increased by the asymmetry in the critical currents in large inductive SQUIDs ($I \gg 1$).\(^6\) Finally the complex behavior of the $T_n(H)$ function could be the result of the $R_2(H)$ dependence due to the scattering of the $I_i$ parameters.

In order to estimate the power gain $G=M^2\left(dV/d\Phi\right)^2/R_nR_s$,\(^7\) we take $R_s=50\ \Omega$ as the microwave source resistance, the measured dynamic resistance $R_d=30\ \Omega$, and an estimated mutual coupling inductance $M=2.4 \times 10^{-11}$ H, assuming a geometrical coupling coefficient $\alpha=0.2$ (for the SQIF layout shown in Fig. 1). Using the experimental value $dV/d\Phi=2 \times 10^{13}$ s$^{-1}$ obtained from measurements at low frequencies, we get a power gain $G=20$ dB.

In order to estimate the saturation power of the SQIF array, we measured the output signal at low frequencies $f=49$ Hz–90 kHz. The expected saturation power of non-coherently operating SQIF loops increases as the square root of $N$, $P_s \propto \sqrt{N}$.\(^5\) An analysis of the experimental data for the first and the second harmonic of the output signal response with an applied 900 Hz signal to the 20 loop SQIF biased at $I/I_c=1,1$ shows quasilinear dependence over 60 dB (see the inset in Fig. 4). We observed a signal distortion with the second harmonic amplitude of about 1% relative to that of the first harmonic amplitude. Note that a similar harmonic distortion was reported for a LTS SQUID amplifier.\(^8\)

Summarizing, we have fabricated and studied HTS SQIFs operating in the F-mode where the magnetic field-to-voltage response is mostly determined by a Fraunhofer dependence of the critical current of the individual JJs. The flux-to-voltage conversion factor in the F-mode is apparently lower than the factor expected in the originally suggested S-mode due to the smaller effective areas of the JJs than of the SQUID loops. Nevertheless our SQIF showed a power gain $G>1$ in the F-mode and a significant increase in saturation power and dynamic range.

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