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Trends in the Exchange Current for Hydrogen Evolution

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A density functional theory database of hydrogen chemisorption energies on close packed surfaces of a number of transition and noble metals is presented. The bond energies are used to understand the trends in the exchange current for hydrogen evolution. A volcano curve is obtained when measured exchange currents are plotted as a function of the calculated hydrogen adsorption energies and a simple kinetic model is developed to understand the origin of the volcano. The volcano curve is also consistent with Pt being the most efficient electrocatalyst for hydrogen evolution.

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The hydrogen evolution reaction is one of the most important electrochemical reactions, whereby protons from solution combine with electrons at an electrode to form first hydrogen atoms chemisorbed at the electrolyte surfaces and then H2 gas. The ability of a given metal to catalyze the hydrogen evolution reaction is usually measured by the exchange current density, which is the rate of hydrogen evolution per surface area at the electrode potential, where the reaction is in equilibrium.

Different materials have widely different exchange current densities. In this paper, we investigate the origin of these variations. We calculate the chemisorption energy for hydrogen for a number of metals. We also collect a number of measured exchange current densities for the same metals and show a volcano curve when the measured currents are plotted as a function of the calculated hydrogen adsorption energies. We show that we are able to understand the variation in the exchange current semi-quantitatively using a simple kinetic model. We also show that the fact that the Pt group metals are best for hydrogen evolution follows directly from the thermodynamic and electronic structure of the reaction. Finally, we discuss the relationship of the volcano curve presented here to other volcano curves introduced previously.1,2

Database for Hydrogen Chemisorption Energies

The starting point for our discussion is a consistent set of hydrocarbon chemisorption energies. We have obtained these from density functional theory (DFT) calculations. The electronic structure problem has been solved using a plane wave pseudopotential implementation,3,4 employing the ultrasoft pseudopotentials of Vanderbilt5 to represent the ionic cores. All calculations were performed with the RPBE exchange-correlation functional6 on periodically repeated metal slabs.

The surface coverage of the adsorbates was 1/4 or 1 monolayer. In both cases the integral adsorption energy is defined by

\[ \Delta E_{\text{H}} = \frac{1}{n} \left( E(\text{surf} + n\text{H}) - E(\text{surf}) - \frac{n}{2} E(\text{H}_2) \right) \]  

where \( n \) is the number of H atoms in the calculation. We have used a \( 2 \times 2 \) surface cell for all calculations, and for this unit cell, \( n = 1 \) corresponds to a coverage of 1/4, while \( n = 4 \) corresponds to a coverage of 1. Unless otherwise noted, all calculations were done on a three layer fcc(111) slab at the RPBE lattice constant separated by five equivalent layers of vacuum. Adding a fourth layer changes the hydrogen binding energy on Pt(111) by 0.01 eV. The bottom two layers were fixed, and the top layer was allowed to relax. The adsorbate was in the fcc site. A \( 4 \times 4 \times 1 \) Monkhorst-Pack k-point setup was used, with maximum symmetry applied to reduce the number of k-points in the calculations. The dipole correction was used in all cases. Spin-polarization did not affect the trends presented here, and was not included. The plane wave cutoff was 350 eV for H, and 450 eV for O adsorption calculations.

Experimental

In Table I we include a collection of experimental data for exchange currents for hydrogen evolution. We have included both the polycrystalline data as well as data from single crystal experiments. There are some variations from one measurement to the next on the same metal, but clearly the trends are the same no matter which measurement is used.

Volcano Curve

We are now in a position to generate a plot of the experimental exchange currents as a function of the calculated hydrogen chemisorption energies. We use the low coverage chemisorption energies for all metals. We will return to the effect of coverage later. The result is shown in Fig. 1. Clearly, the data order quite nicely into a volcano. Note that all the reactive metals are on the left “leg” of the volcano and the unreactive ones, Cu, Ag, and Au, are on the right.

Kinetics and the Origin of the Volcano

The total hydrogen evolution reaction can be written

\[ \text{H}^+ + e^- \rightarrow \frac{1}{2}\text{H}_2 \]  

It takes place at an electrode supplying the electrons, and providing an intermediate state of the process:

\[ \text{H}^+ + e^- + * \rightarrow \text{H}^* \]  

\[ 2\text{H}^* \rightarrow \text{H}_2 \]

where the * denotes a site on the surface (so an * by itself denotes a free site and \( \text{H}^* \) denotes a hydrogen atom adsorbed on the surface). The final hydrogen evolution step may also be

\[ \text{H}^+ + e^- + \text{H}^* \rightarrow \text{H}_2 \]

We now consider hydrogen evolution at a metal electrode, which is part of an electrochemical cell. During hydrogen evolution a current \( i \) will be running

\[ i = -er \]  

where \( r = r_\text{f} - r_\text{ns} \) is the net rate of (Eq. 2).

The exchange current is the forward (and backward) rate when Eq. 2 is in equilibrium.
1/2 H₂ is the adsorbed state is small meaning that the entropy of adsorption of barriers other than those coming from the thermochemistry. This contributions to the variation in the rate constants from any reaction contributions to these trends originating directly from the variations trends in exchange rates from one metal to the next and on the potential.

than the other metals for hydrogen evolution at the equilibrium potential over Pt the reaction is essentially thermo-neutral where as this is not for H/Cu

representative for all the metals we study here. This means that gas phase at standard conditions.

The free energy diagrams for the reaction at equilibrium over Pt(111), Ni(111), Mo(110), and Au(111) are shown in Fig. 2. At equilibrium the free energy per H atom (the chemical potential) of the initial and final states of reaction (Eq. 2) are the same. In Fig. 1 and 2 we have also included the free energy of the adsorbed state calculated as

$$D_{H} = \Delta E_{H} + \Delta E_{ZPE} - T \Delta S_{H}$$

where $\Delta E_{H}$ is the hydrogen chemisorption energy from Table I and $\Delta E_{ZPE}$ is the difference in zero point energy between the adsorbed and the gas phase. We exploit the fact that the vibrational entropy in the adsorbed state is small meaning that the entropy of adsorption of 1/2 H₂ is $\Delta S_{H} \approx -1/2 S_{H}^{0}$, where $S_{H}^{0}$ is the entropy of H₂ in the gas phase at standard conditions. $\Delta E_{ZPE}$ is calculated to be 0.04 eV for H/Cu(111), see the Appendix,¹ and we take this value to be representative for all the metals we study here. This means that $\Delta G_{H^+} = \Delta E_{H} + 0.24$ eV.

In Fig. 2 we have chosen to show data for low coverage (1/4 monolayer). As indicated in Table I, trends will not change if we include the possibility of a higher coverage. The figure indicates that over Pt the reaction is essentially thermo-neutral where as this is not the case for metals on either side of the volcano, e.g., Ni, Au, and Mo. This in itself indicates why Pt is a much better electrocatalyst than the other metals for hydrogen evolution at the equilibrium potential.

We now turn to the kinetic modelling. We will concentrate on the trends in exchange rates from one metal to the next and on the contribution to these trends originating directly from the variations in the thermochemistry, Fig. 2 and Table II. We therefore neglect contributions to the variation in the rate constants from any reaction barriers other than those coming from the thermochemistry. This means that the variations from one metal to the next in the rate constants for getting out of the chemisorption state, H⁰ in Fig. 2, are the same. Combined with the fact that at equilibrium the rates involved in the forward and backward steps must be the same, it means that the exchange rates of the two consecutive steps are the same. We choose here to focus on the rate of the proton reduction, Eq. 3. The forward rate of Eq. 3 can be written

$$r_1 = k_1(1 - c_H^+)$$

where $k_1$ is the rate constant, 0 is the H coverage (measured in fraction of a monolayer with respect to the number of metal atoms in the surface layer), and $c_H^+$ is the concentration of protons in the electrolyte. The coverage of hydrogen can be calculated from Eq. 4, and hence

$$i_0 = i_{-1/4} = \frac{v}{i_{-1}} = -e r_1$$

Table I. Calculated hydrogen adsorption energies ($\Delta E_{H}$) and measured exchange currents ($i_0$) on different transition and noble metals. The hydrogen chemisorption energies are calculated both for low (1/4 monolayer) and high (1 monolayer) coverage. The references to the experimental data are shown in the right column.

<table>
<thead>
<tr>
<th>Metal</th>
<th>$\Delta E_{H}$ (eV)</th>
<th>$\Delta E_{H}$ (eV)</th>
<th>$\Delta E_{H}$ (eV)</th>
<th>$\Delta E_{H}$ (eV)</th>
<th>Log(i₀/A cm⁻²)</th>
<th>Ref.</th>
</tr>
</thead>
<tbody>
<tr>
<td>Au</td>
<td>0.21</td>
<td>0.37</td>
<td>-6.6</td>
<td>-6.8</td>
<td>8</td>
<td></td>
</tr>
<tr>
<td>Au(111)</td>
<td>0.27</td>
<td>0.38</td>
<td>-6.8</td>
<td>-5.4</td>
<td>0.33</td>
<td></td>
</tr>
<tr>
<td>Ag</td>
<td>0.27</td>
<td>0.34</td>
<td>-5.0</td>
<td>-7.85</td>
<td>1</td>
<td></td>
</tr>
<tr>
<td>Pd</td>
<td>-0.38</td>
<td>-0.33</td>
<td>-3</td>
<td>-3</td>
<td>9</td>
<td></td>
</tr>
<tr>
<td>Pt</td>
<td>-0.33</td>
<td>-0.27</td>
<td>-3.1</td>
<td>-2.63</td>
<td>9</td>
<td></td>
</tr>
<tr>
<td>Pt(111)</td>
<td>0.34</td>
<td>0.34</td>
<td>3.34</td>
<td>12</td>
<td>9</td>
<td></td>
</tr>
<tr>
<td>Rh</td>
<td>-0.34</td>
<td>-0.30</td>
<td>-3.6</td>
<td>-3.22</td>
<td>9</td>
<td></td>
</tr>
<tr>
<td>Ir</td>
<td>-0.21</td>
<td>-0.16</td>
<td>-3.7</td>
<td>-3.46</td>
<td>9</td>
<td></td>
</tr>
<tr>
<td>Ni</td>
<td>-0.51</td>
<td>-0.47</td>
<td>-5.2</td>
<td>-5.21</td>
<td>9</td>
<td></td>
</tr>
<tr>
<td>W</td>
<td>-0.67</td>
<td>-0.83</td>
<td>-5.9</td>
<td>-5.9</td>
<td>9</td>
<td></td>
</tr>
<tr>
<td>Co</td>
<td>-0.51</td>
<td>-0.49</td>
<td>-5.32</td>
<td>-5.32</td>
<td>9</td>
<td></td>
</tr>
<tr>
<td>Cu</td>
<td>-0.05</td>
<td>0.03</td>
<td>-5.37</td>
<td>-5.37</td>
<td>9</td>
<td></td>
</tr>
<tr>
<td>Mo</td>
<td>-0.61</td>
<td>-0.77</td>
<td>-7.01</td>
<td>-7.01</td>
<td>9</td>
<td></td>
</tr>
<tr>
<td>Re</td>
<td>-0.56</td>
<td>-0.45</td>
<td>-2.87</td>
<td>-2.87</td>
<td>9</td>
<td></td>
</tr>
<tr>
<td>Nb</td>
<td>-0.80</td>
<td>-0.80</td>
<td>-6.8</td>
<td>-6.8</td>
<td>9</td>
<td></td>
</tr>
</tbody>
</table>

¹ bcc(110)
To be activated by at least D.

Clearly, a simple model, including in particular the position of the maximum.

The exchange current is

\[ i_0 = -e k_0 \frac{1}{1 + \exp(-\Delta G_{H^+}/kT)} \]  

[12]

For the other case where the proton transfer is endothermic (∆GH⁺ > 0, e.g., Au in Fig. 2), we would expect the proton transfer to be activated by at least ∆GH⁺. Under the same assumptions as above, the rate constant for this case is

\[ k_1 = k_0 \exp(-\Delta G_{H^+}/kT) \]  

[13]

and the exchange current is

\[ i_0 = -e k_0 \frac{1}{1 + \exp(-\Delta G_{H^+}/kT)} \exp(-\Delta G_{H^+}/kT) \]  

[14]

Figure 3 suggests the following physical picture. At equilibrium, the low coverage state is fully occupied. The addition of a new proton and electron must add a hydrogen atom in the high coverage state. Both states can turn over, but the overall activation energy is smaller for the high coverage state, making this the most likely state. Both states can turn over, but the overall activation energy is smaller for the high coverage state, making this the most likely state.

Water Dissociation on the Surface

One question in connection with the experimental data in Table I and Fig. 1 is whether reaction products from the adsorption and
dissociation of water might cover some of the metals. To investigate this we have first calculated the effect of molecular water on the adsorption. The result is that adding a water overlayer changes the adsorption energy of hydrogen on Pt(111) by less than 0.02 eV. Even though the effect could be a little larger for the more reactive metals it seems justified to neglect it in the present trend study.

In addition we have calculated the reaction energies for the reactions

\[ H_2O + \ast \rightarrow OH^\ast + 1/2H_2 \quad [16] \]

and

\[ H_2O + \ast \rightarrow O^\ast + H_2 \quad [17] \]

The results are included in Table II. In general one should also consider the possibility of forming solvated H\(^+\) and electrons at the electrode rather than H\(_2\), but for the exchange current at standard conditions this makes no difference since the solvated hydrogen is in equilibrium with H\(_2\)(gas). The values in Table II are relative to gas phase water. We also include the adsorption free energies, calculated using the zero point energies and entropies collected in the Appendix. The result is that of the metals considered in Table II, only Mo and W should be covered by surface oxygen under the conditions of interest here. This means that the measured exchange currents for these metals in Fig. 1 should not be taken to be representative of the metal, but rather of the surface oxide.

Conclusions

The results presented above clearly demonstrate the correlation between the hydrogen chemisorption energies and the exchange current for hydrogen evolution. The volcano curves in Fig. 1 also indicate that Pt is a better electrocatalyst than other metals for hydrogen evolution, primarily because the evolution reaction is thermo-neutral on Pt at the equilibrium potential. These findings can potentially be useful in predicting bimetallic electrocatalysts for hydrogen evolution, as well as for hydrogen electro-oxidation for fuel cell applications.

Acknowledgments

Useful discussions with J. Meier and A. K. Pal are appreciated. This work was funded in part by U.S. Army Research Office (ARO), USA, under grant DAAD19-02-1-0311 (U.S. and S.P.). The DFT calculations have been performed with support from the Danish Center for Scientific Computing through grant no. HDW-1101-05. A.L. acknowledges financial support from the European Union through contract no. ENK5-CT-2001-00572.

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### Table A-1. Zero point energy corrections and entropic contributions to the free energies.

<table>
<thead>
<tr>
<th></th>
<th>( T) (eV)</th>
<th>( T\Delta S) (eV)</th>
<th>( E_{ZPE}) (eV)</th>
<th>( \Delta E_{ZPE}) (eV)</th>
</tr>
</thead>
<tbody>
<tr>
<td>( H_2O) (0.035 bar)</td>
<td>0.67</td>
<td>0</td>
<td>0.56</td>
<td>0</td>
</tr>
<tr>
<td>( ^*OH + 1/2H_2)</td>
<td>0.20</td>
<td>-0.47</td>
<td>0.44</td>
<td>-0.12</td>
</tr>
<tr>
<td>( ^*O + H_2)</td>
<td>0.41</td>
<td>-0.27</td>
<td>0.34</td>
<td>-0.22</td>
</tr>
<tr>
<td>( 1/2 O_2 + H_2)</td>
<td>0.73</td>
<td>0.05</td>
<td>0.32</td>
<td>-0.24</td>
</tr>
<tr>
<td>( H_2)</td>
<td>0.41</td>
<td>...</td>
<td>0.27</td>
<td>...</td>
</tr>
<tr>
<td>( 1/2 O_2)</td>
<td>0.32</td>
<td>...</td>
<td>0.05</td>
<td>...</td>
</tr>
<tr>
<td>( O^*)</td>
<td>0</td>
<td>...</td>
<td>0.07</td>
<td>...</td>
</tr>
<tr>
<td>( OH^*)</td>
<td>0</td>
<td>...</td>
<td>0.30</td>
<td>...</td>
</tr>
<tr>
<td>( H^*)</td>
<td>0</td>
<td>...</td>
<td>0.17</td>
<td>...</td>
</tr>
</tbody>
</table>

Appendix

In Table A-1 we show the entropies and zero point energies (ZPEs) used in the construction of the free energies in Fig. 2 and 3 and in Table II. The gas phase values are from Ref. 15, while the values for the adsorbed species are taken from DFT calculations.

References