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Full-Wave Analysis of a Two Slot Microstrip Filter Using a New Algorithm for Computation of the Spectral Integrals

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Abstract—An integral equation is formulated in the spectral domain for a two slot microstrip filter using the exact Green’s function for the grounded dielectric substrate. Using a moment method (MM) procedure, the integral equation has been discretized. The elements of the impedance matrix and the excitation vector are given by two dimensional Sommerfeld type integrals in closed form. An efficient and accurate numerical integration scheme for computation of the elements is presented. The $S$ parameters obtained from the MM procedure have been found to be in excellent agreement with measurements.

I. INTRODUCTION

In the present communication we consider the microstrip filter problem shown in Fig. 1. An integral equation for the filter problem is derived using the exact Green’s function for the geometry. Hence, both free space radiation and surface waves are included in the solution. Using basis functions modeling the incident, reflected and transmitted currents on the filter [1], the integral equation is solved numerically for the scattering parameters using a Galerkin type MM procedure.

The elements of the impedance matrix and source vector obtained from the MM procedure are given by two dimensional integrals over the spectral coordinates $k_x$ and $k_y$.

Instead of integrating these integrals in the polar coordinate system [1], an efficient and accurate scheme is applied in which the integrals are integrated in the cartesian $k_x$ and $k_y$ coordinates using a deformation technique in which the integration contour follows a path parallel to the imaginary axis in the complex plane for one of the integration variables [2] while the other integration variable is integrated along the real axis using the weighted average algorithm [3]. In this communication no approximation for the Green’s function is used nor are the integration limits truncated.

In order to test the integration algorithm, the microstrip filter problem was chosen since the space radiation and surface wave radiation at resonance were expected to be substantial, hence the transfer function of the filter would be suited for testing against precision measurements using a HP8510 network analyzer.

II. THEORY

In Fig. 1, the outline of the microstrip filter is shown. The filter is in the form of an infinite microstrip line in which two narrow gaps are cut. The substrate has thickness $d$, relative permittivity $\varepsilon_r$, and is assumed to be infinitely wide in the $x$ and $y$ directions. The filter lines are assumed to be infinitely thin and perfectly conducting. Since the width of the filter lines are narrow in terms of wavelength, we consider only $x$-directed surface currents. [4]

Green’s Function for the Grounded Dielectric Slab:
Using the spectral domain Green’s function, we obtain the $x$-directed electric field at $(x, y, d)$ from an $x$-directed current distribution [5], [6]:

$$E_{xx}(x, y, d) = \frac{1}{4\pi^2} \int \int_{-\infty}^{\infty} \hat{G}_{xx}(k_x, k_y) \hat{K}_x(k_x, k_y) e^{-j k_x x} e^{-j k_y y} dk_x dk_y$$

(1)

where

$$\hat{K}_x(k_x, k_y) = \int \int_{-\infty}^{\infty} K_x(x, y, d) e^{-j k_x x} e^{-j k_y y} dx dy$$

(2)

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$K_x(x, y, d)$ is the $x$-directed electric surface current distribution on the microstrip filter lines.

$$
\hat{G}_{xx} = -\frac{j \pi}{k_0} (\varepsilon_r k_0^2 - k_2^2) k_{22} \cos(k_{12} d) + j(k_0^2 - k_2^2) k_{21} \sin(k_{12} d)
$$

$$
\cdot \sin(k_{12} d), k_0 = \omega \sqrt{\mu / \varepsilon}
$$

$$
T_m = \varepsilon_r k_0^2 \cos(k_{12} d) + j k_{21} \sin(k_{12} d)
$$

$$
T_x = k_{21} \cos(k_{12} d) + j k_{22} \sin(k_{12} d)
$$

$$
k_{21}^2 = \varepsilon_r k_0^2 - \lambda^2 \quad \text{Im}(k_{11}) < 0
$$

$$
k_{22}^2 = k_0^2 - \lambda^2 \quad \text{Im}(k_{22}) < 0
$$

Because of the $x$-dependence of the $x$-directed current distribution (1) is reduced to (12), \[1\]:

$$
E_{xx}(x, y, d) = \frac{1}{4\pi^2} \int_{-\infty}^{\infty} \hat{G}_{xx}(k_x, k_y) \hat{K}_x(k_y) e^{jk_y y} dk_y.
$$

Applying Galerkin’s method of moments on (12) yields

$$
\begin{bmatrix}
Z_{11} & Z_{12} & Z_{13} \\
Z_{21} & Z_{22} & Z_{23} \\
Z_{31} & Z_{32} & Z_{33}
\end{bmatrix}
\begin{bmatrix}
C_1 \\
C_2 \\
C_3
\end{bmatrix} = \begin{bmatrix}
0 \\
0
\end{bmatrix}
$$

where

$$
Z_{pq} = \frac{1}{2\pi^2} \int_{0}^{\infty} \hat{G}_{xx}(k_x, k_y) \hat{f}_{pq}(k_y) \hat{f}_{pq}(k_y) dk_y.
$$

Infinite Microstrip Line

An accurate solution of the microstrip filter problem requires that the effective propagation constant $k_e$ for the current distribution on the half infinite microstrip lines is computed accurately. The effective propagation constant is computed from the corresponding infinite microstrip line problem where we have assumed the traveling-wave form $e^{-jk_e x}$ \[1\] for the $x$-dependence of the $x$-directed current distribution.

Since the main computational effort is spent computing the impedance matrix elements obtained from the Galerkin procedure, a judicious choice of basis functions is important to obtain a convergent solution using a minimal number of basis functions. The $y$-dependence of the current distribution is expanded in a set of entire domain Maxwellian basis functions (10) which are also used as test functions \[7, p. 342\]. The Maxwellian functions satisfy the proper edge condition for the surface current and are even functions in the transverse coordinate $y$, as required by the symmetry of the geometry.

The method of solution involves formulating an integral equation for the $x$-directed electric field on the microstrip line which is discretized using the Galerkin’s method of moment procedure. The linear system obtained is subsequently solved for the unknown expansion coefficients:

$$
K_x(x, y) = K_x(x) K_y(y)
$$

$$
K_x(y) = \sum_{q=1}^{N_y} C_q f_{pq}(y)
$$

$$
K_x(x) = e^{-jk_e x}
$$

where

$$
f_{pq}(y) = \frac{\cos \left( y(q-1) \frac{2\pi}{w} \right)}{\sqrt{1 - \left( y \frac{2\pi}{w} \right)^2}}; \quad q = 1, 2, \ldots
$$

In the spectral domain, the basis functions are given by

$$
\hat{f}_{pq}(k_y) = \frac{w\pi}{4} \left( J_0 \left( k_y \frac{w}{2} + (q-1)\pi \right) + J_0 \left( k_y \frac{w}{2} - (q-1)\pi \right) \right); \quad q = 1, 2, \ldots
$$

$$
f_s(u) = \begin{cases}
\sin u & -M\pi < u < 0 \\
0 & \text{otherwise}
\end{cases}
$$

and

$$
K_{tr} = \left( -T f_{u1}(x) - jT f_{u2}(x) \right) f(y)
$$

where

$$
f_{u1}(x) = f_u \left( k_e(x - L - 2G) - \frac{\pi}{2} \right)
$$

$$
f_{u2}(x) = f_u \left( k_e(x - L - 2G) \right)
$$

$$
f_u(u) = \begin{cases}
\sin u & 0 < u < M\pi \\
0 & \text{otherwise}
\end{cases}
$$

$$
f(y) = \sum_{q=1}^{N_y} C_q f_{pq}(y), C_1 = 1.
$$
Whether \( M \) in (19) and (22) is an integer value or not is of no consequence for the numerical convergence of the corresponding integrals obtained later from the Galerkin procedure when computed by the integration scheme described in this paper. However, \( M \) is chosen as an integer since this yields a simpler expression in the spectral domain. Numerical tests indicate that the solution to the filter problem is almost independent of \( M \) when the traveling wave modes are longer than five wavelengths \( (M = 10) \). For the computations in the present paper, we have used \( M = 15 \).

Subsectional piecewise sinusoidal (PWS) modes have been used for modeling the \( x \)-dependence of the current in the vicinity of the gaps as shown in Fig. 2. The PWS modes are defined

\[
\tilde{f}_{mn}(x) = \begin{cases} 
\sin k_m(h_n - |x - x_n|) & [x - x_n] \leq h_n \\
0 & [x - x_n] > h_n
\end{cases}
\]

(24)

where \( 2h_n \) is the width and \( x_n \) is the position of the centre of the \( n \)th PWS function. The \( y \)-dependence of the \( N_x, x \)-directed PWS modes are expanded into a sum of \( N_y \) Maxwellian modes (10).

Substituting the true current with an expansion using the basis functions described, the \( x \)-directed E-field at the upper surface of the dielectric layer may be written in the form of an integral equation (25):

\[
E_x(x, y, d) = \int_{-\infty}^{\infty} \tilde{G}_{xx}(\vec{k} \cdot \vec{e}_x) \left( \tilde{K}_{inc} + \tilde{K}_{ref} + \tilde{K}_{tr} + \sum_{q=1}^{N_y} \sum_{n=1}^{N_x} \tilde{K}_{nq} \tilde{f}_{x_{n}} \tilde{f}_{y_{q}} \right) e^{ik_x x} e^{ik_y y} dy \, dx,
\]

(25)

Equation (25) is tested with \( (N_x + 2)N_y \) PWS test functions (10), (24). Since the traveling wave modes are obtained from the solution of the infinite microstrip line problem, the boundary condition for the tangential electric field is automatically satisfied away from the gaps. Hence, it is only necessary to test (25) in the gap regions.

The PWS test functions are located as the PWS basis functions, except for the two additional test functions necessary to obtain the same number of equations as unknowns. The two additional test PWS functions are located immediately after the last PWS basis function at each side of the two gaps. The \( y \)-dependence \( K_s(y) \) of these two test functions are assumed identical to the traveling wave modes.

The impedance matrix elements can then be defined using the following indices:

- \( n \): Index of PWS basis functions \( x \)-dependent part
- \( q \): Index of PWS basis functions \( y \)-dependent part
- \( m \): Index of PWS test functions \( z \)-dependent part
- \( p \): Index of PWS test functions \( y \)-dependent part
- \( t_l \): Additional test functions on the left microstrip line
- \( t_r \): Additional test functions on the right microstrip line.

The testing procedure results in a matrix equation for the unknown coefficients \( R, T, K_{11}, \ldots, K_{N, N} \).

A simple example with two PWS modes on the centre strip (and none on each of the half infinite lines), each expanded in two Maxwellian modes for the \( y \)-dependence, yields matrix equation (31), which is shown at the bottom of the next page.

\[
Z_{mnqp} = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \tilde{G}_{xx}(\vec{k} \cdot \vec{e}_x) \tilde{f}_{x_{n}} \tilde{f}_{x_{m}} \tilde{f}_{y_{q}} \tilde{f}_{y_{p}} dy \, dx
\]

(26)

\[
Z_{mpsi} = \sum_{q=1}^{N_y} \int_{-\infty}^{\infty} \tilde{G}_{xx}(\vec{k} \cdot \vec{e}_x) \tilde{f}_{x_{i}} \tilde{f}_{x_{m}} \tilde{f}_{y_{q}} \tilde{f}_{y_{p}} dy \, dx \quad i \in \{1, 2\}
\]

(27)

\[
Z_{mpui} = \sum_{q=1}^{N_y} \int_{-\infty}^{\infty} \tilde{G}_{xx}(\vec{k} \cdot \vec{e}_x) \tilde{f}_{x_{i}} \tilde{f}_{x_{m}} \tilde{f}_{y_{q}} \tilde{f}_{y_{p}} dy \, dx \quad i \in \{1, 2\}
\]

(28)

\[
Z_{tinq} = \sum_{p=1}^{N_y} \sum_{q=1}^{N_y} \int_{-\infty}^{\infty} \tilde{G}_{xx}(\vec{k} \cdot \vec{e}_x) \tilde{f}_{x_{n}} \tilde{f}_{x_{q}} \tilde{f}_{y_{q}} \tilde{f}_{y_{p}} dy \, dx \quad t \in \{1, 2\}
\]

(29)

\[
Z_{tiij} = \sum_{q=1}^{N_y} \int_{-\infty}^{\infty} \tilde{G}_{xx}(\vec{k} \cdot \vec{e}_x) \tilde{f}_{x_{i}} \tilde{f}_{x_{j}} \tilde{f}_{y_{q}} \tilde{f}_{y_{p}} dy \, dx \\
= \sum_{q=1}^{N_y} \sum_{p=1}^{N_y} C_q C_p \int_{-\infty}^{\infty} \tilde{G}_{xx}(\vec{k} \cdot \vec{e}_x) \tilde{f}_{x_{i}} \tilde{f}_{x_{j}} \tilde{f}_{x_{p}} \tilde{f}_{y_{q}} dy \, dx
\]

(30)
III. COMPUTATION OF THE Z AND THE V ELEMENTS

The main computational effort when solving the filter problem is the numerical computation of the individual elements in the impedance matrix \( Z \) and the excitation vector \( V \). The elements are given by two dimensional Sommerfeld type integrals given in closed form in (26) to (30), which is shown at the bottom of the previous page. The integrands contain poles corresponding to the generation of surface waves. Furthermore, the integrands are rapidly oscillating as \( k_x^2 + k_y^2 \to \infty \), especially as the distance between basis and test functions increases.

Usually, the spectral domain integration intervals in (26) to (30) are transformed into finite and semi-infinite intervals by changing the cartesian integration variables \((k_x, k_y)\) into the polar integration variables \( (\lambda, \phi) \) [3], [8], [9].

One difficulty when using the polar integration scheme is to compute the \( \phi \) integration accurately since the oscillation frequency of the integrand is directly proportional to the integration variable \( \lambda \). Attempting to truncate the semi-infinite integration interval for \( \lambda \) in order to set an upper bound for the oscillations in \( \phi \) reduces the obtainable numerical accuracy of the integral which is slowly convergent.

In this work, the numerical computation of the integrals is made in the cartesian coordinates \((k_x, k_y)\) without truncation of the infinite intervals.

Since \( \tilde{G}_x \) and \( \tilde{G}_y \) are real functions and \( \tilde{G}_{xx} \) is an even function of \( k_x \) and of \( k_y \) the integrals in (26) to (30) can be written as in (32):

\[
Z = \int_{-\infty}^{\infty} S(k_y) \tilde{f}_a(k_y) \tilde{f}_b^*(k_y) dk_y
\]

where

\[
S(k_y) = \int_{-\infty}^{\infty} \tilde{G}_{xx}(k_x, k_y) \tilde{f}_a(k_x) \tilde{f}_b^*(k_x) dk_x
\]

\[
= \int_{-\infty}^{\infty} \tilde{G}_{xx}(k_x, k_y)(\tilde{f}_a(k_x) \tilde{f}_b^*(k_x) + \tilde{f}_a^*(k_x) \tilde{f}_b(k_x)) dk_x.
\]

The singularities of the Green’s function are located in the range \( k_y^2 < k_x^2 + k_y^2 < k_0^2 \) [11] and correspond to transverse magnetic surface waves \( T_m(\lambda) = 0 \) and/or transverse electric surface waves \( T_e(\lambda) = 0 \). In the \((k_x, k_y)\) plane, the poles are located on concentric circles. Hence, when integrating along a contour parallel to the \( k_y \) axis, the position of singular points (if any) depends on the value of \( k_x \). In order to avoid the difficulty of \( k_x \)-dependent singular points when integrating parallel to the \( k_y(k_x) \) axis, the integration path has been deformed into the complex \( k_y(k_x) \) plane in the range from 0 to \( (1 + \sqrt{c})k_0 \) using the Cauchy integration theorem [12]. The deformed integration path used for both the \( k_x \) and the \( k_y \) integration variables is shown in Fig. 3.

Alternative integration contours have been investigated. Rixberg [9] uses the upper half of an ellipse for the integration path in the \( \lambda \) plane. In [2a], a triangle is used for the integration contour in both the \( k_x \) and the \( k_y \) planes.

The integration path cannot be arbitrarily deformed into the first quadrant for the \( k_x \) variable when using the approach described in the present paper since the integration contour chosen later for the \( k_y \) variable introduces branch cuts in this quadrant. Furthermore, care must be taken when choosing the path for numerical reasons. If the path is too close to the singular points on the real axis, the values of the integrand become too large for accurate numerical computation. On the other hand, if the integration path is too far from the real axis, the integrand increases exponentially in amplitude.

It is necessary to integrate the three linear parts of which the deformed contour is composed independently since the path itself is not differentiable. In spite of this, we found from numerical experiments that the integration contour in Fig. 3 was more efficient in terms of integration points for a given accuracy than the elliptic path in the complex \( k_x \) and \( k_y \) planes.

Furthermore, we found that the optimum values for \( \text{Im}(k_x) \) and \( \text{Im}(k_y) \) to be about 0.1\( k_0 \), which is similar to the value choose by Newman and Forrai [10] for the deformation into the complex \( \lambda \) plane.

\[
\begin{align*}
Z_{1111} & = Z_{1112} = Z_{1121} = Z_{1211} = (Z_{1111} - jZ_{1112}) = Z_{1121} = jZ_{1221} = Z_{1211} = (Z_{1111} - jZ_{1112}) = jZ_{1221} = Z_{1221} = (Z_{1111} - jZ_{1112})
\end{align*}
\]
It is important to note that the integrand must be analytical when the Cauchy integration theorem is used. Hence, the complex conjugated terms in (32) and (33) must be computed as if \( k_x \) and \( k_y \) were purely real because the \( k_x^* \) and \( k_y^* \) terms are not analytical.

The \( k_y \) integration over the interval from \((\sqrt{\tau}+1)k_0\) to \(k_{\text{asymp}} = 45/W\) is made using the Gauss-Legendre integration algorithm. For \( k_y \) larger than \( k_{\text{asymp}} \), the product between the basis and test functions in (32) is written as a sum of functions (34) composed of an oscillating and a nonoscillating part given in closed form using the Jacobi asymptotic expansion [13] for the zero order Bessel functions.

\[
\tilde{f}_n(k_y) \tilde{f}_n^*(k_y) + \tilde{f}_n^*(k_y) \tilde{f}_n(k_y) = \sum_i \tilde{A}_i(k_y) \cos(\omega_i k_y). \quad (34)
\]

For the Maxwellian basis and test functions used for the transverse dependence of the microstrip line current, the right hand side of (34) contains two terms. One term has twice the oscillation frequency of the Bessel functions while the other term becomes nonoscillating. The integration of the nonoscillating part is made using the Gauss-Legendre integration algorithm after applying the variable transformation (44) yielding a finite integration interval can be applied without introducing numerical difficulties.

\[
\begin{aligned}
Z &= \int_{0}^{(\sqrt{\tau}+1)k_0} S(k_x) \left( \tilde{f}_n(k_x) \tilde{f}_n^*(k_x) + \tilde{f}_n^*(k_x) \tilde{f}_n(k_x) \right) dk_x \\
&\quad + \sum_i Z_i,
\end{aligned}
\]

where

\[
Z_i = \int_{(\sqrt{\tau}+1)k_0}^{\infty} S(k_x) \tilde{A}_i(k_x) e^{j\omega_i k_x} dk_x. \quad (39)
\]

It is important that all the integrands are analytical and that the integration contours do not cross any branch cuts. These conditions are met for the integrand in (39). In Fig. 4, it is shown how the deformation and the transformation is made for \( \omega_i > 0 \).

We have proved that the integration contour along \( C^{\pm}_{(x)} \) does not contribute to \( Z_i \) (39). Hence the integrals along \( C^{+}_{y} \) and \( C^{-}_{y} \) are identical.

\[
Z_i = \int_{C^{+}_{y}} \tilde{S}(k_x) \tilde{A}_i(k_x) e^{j\omega_i k_x} dk_x. \quad (40)
\]

We define

\[
\tilde{F}_i(k_x) = \tilde{S}(k_x) \tilde{A}_i(k_x). \quad (41)
\]

The path \( C^{+}_{y} \) is given by

\[
k_x = A + j \tau \quad \quad \frac{dk_x}{dt} = j \quad t \in [0, \infty) \quad \text{where} \quad A = (\sqrt{\tau}+1)k_0. \quad (42)
\]

From (40), (41) and (42) we obtain

\[
Z_i = \int_{0}^{\infty} \tilde{F}_i(A + j \tau) e^{j(A+t)\omega_i} j dt = j e^{jA\omega_i} \int_{0}^{\infty} \tilde{F}_i(A + j t) e^{-j\omega_i t} dt \quad \omega_i > 0. \quad (43)
\]

Since the integrand in (43) is decreasing exponentially, the variable transformation (44) yielding a finite integration interval can be applied without introducing numerical difficulties.

\[
t = -\frac{1}{g} \log z; \quad \frac{dz}{dt} = -\frac{1}{g}z; \quad g > 0. \quad (44)
\]

Using (44) upon (43) we obtain

\[
\begin{aligned}
Z_i &= j e^{jA\omega_i} \int_{0}^{1} \tilde{F}_i \left( A - \frac{j}{g} \log z \right) e^{j\omega_i t} g^{-1} dt \\
&\quad \quad \omega_i > 0. 
\end{aligned}
\]

(45)

For \( \omega_i < 0 \) a similar procedure is used. For \( \omega_i = 0 \) the variable transformation (35) is applied for the \( k_x \) variable.
In summary, see (46), which is at the bottom of this page. In (44) we have introduced a constant \( g > 0 \). The value of (46) is independent of \( g \). However, the location and the number of integration points necessary for the numerical computation of the integral in (46), given a certain accuracy, is dependent upon the actual value of \( g \). Hence, the optimum, in terms of the computational effort, is to minimize the sum of the number of integration points used with the selected values of \( g \).

In all the integrals for the filter, we only use \( \tilde{G}_{xx} \) from the dyadic Green’s function, however the integration scheme described in this paper can be applied for any of the remaining elements in the dyadic Green’s function.

Since \( \tilde{G}_{xx} \) is an even function of \( k_{z1} \), one may ignore the branch cut due to \( k_{z1} = \sqrt{\varepsilon \varepsilon' - k_x^2 - k_y^2} \). The branch cut due to \( k_{z2} = \sqrt{k_o^2 - k_x^2 - k_y^2} \), which applies for \( \tilde{G}_{xx} \) is shown in Fig. 5. From Fig. 5 we appreciate that either the \( k_x \) or the \( k_y \) integration must be performed along a contour parallel to the real axis.

For the \( k_x \) integration, \((25 + 25 + 25 + 12 + 2 \cdot 50) = 187 \) integrations points are used. The first 75 points are used for the deformation contours above the surface wave poles. The 12 integration points are used for the nonoscillating part in (46). The remaining \( 2 \cdot 50 \) points are used for the deformation into the first and fourth quadrants of the complex plane.

The accuracy obtained for the elements of \( Z \) and \( V \) are 4 to 5 significant digits regardless of the distance between the basis and test functions. The accuracy of the current vector (31) is also 4 to 5 digit.

\[
Z_i(k_y) = \begin{cases} 
\frac{j}{g} e^{j\omega_i} \int_0^1 \tilde{S}(A - j \frac{1}{g} \log z, k_y) \tilde{A}_i(A - j \frac{1}{g} \log z) z^{(\omega_i/g) - 1} dz & \text{for } \omega_i > 0 \\
\int_0^{1/j} \tilde{S}(1/v, k_y) \tilde{A}_i(1/v) \frac{1}{v^2} dv & \text{for } \omega_i = 0 \\
\frac{-j}{g} e^{j\omega_i} \int_0^1 \tilde{S}(A + j \frac{1}{g} \log z, k_y) \tilde{A}_i(A + j \frac{1}{g} \log z) z^{(-\omega_i/g) - 1} dz & \text{for } \omega_i < 0 
\end{cases}
\]

(46)
Fig. 6. Current distribution at first resonance. Input parameters: Substrate: RT/duroid 5870, d = 0.7874 mm, ε = 2.33 Dimension of the filter: W = 2.30 mm, G = 0.3 mm, L = 10 mm. Frequency = 9.66 GHz. Basis functions: (41, 41, 41)*3.

Fig. 7. Current distribution at second resonance. Input parameters are the same as used in Fig. 6 except for the frequency. Frequency = 19.32 GHz.

IV. RESULTS

Using the described procedure for computation of the elements in (31), the current distribution on the filter has been determined from solving the linear equations. In Fig. 5, 6 and 7, the current distribution at the first and second resonance frequency are shown. The value of the current distributions at the edges are not shown due to the edge singularity. The number of basis functions used for modeling the filter currents are written using the notation (x, y, z)*v. The x, y and z denote the number of PWS basis functions used for modeling the current on the input, centre and output microstrip line respectively. v denotes the number of Maxwellian basis function used for the y-dependence. Regardless of the number of PWS basis functions, their base covers the area indicated in Fig. 2.

In Figs. 6 and 7, a standing wave at the semi-infinite input microstrip line and a transmitted wave at the semi-infinite output microstrip line can be observed. The excited currents on the centre microstrip line demonstrates the first and second resonance. Since the x directed current is singular at the edges parallel to the x axis (10), the current shown in Figs. 5, 6 and 7 does not include the current value at the edges.

In Fig. 8, the amplitude of the S11 and S21 parameters for the microstrip filter. Input parameters: Substrate: RT/duroid 5870, d = 0.7874 mm, ε = 2.33 Dimension of microstrip filter: W = 2.30 mm, G = 0.3 mm, L = 10 mm.

V. CONCLUSION

A full-wave analysis has been presented for the problem of microstrip gap discontinuities. The S11 and S21 parameters of the two gap microstrip filter have been computed numerically and compared with measurements. The agreement is excellent. Plots of the current distribution and the parameters have been presented. The spectral integral has been computed using a new integration algorithm which is both fast and accurate. The method is efficient in terms of the number of integration points needed for convergence. The basis and test functions are
Fig. 9. Measured and computed amplitude of the $S_{11}$ and $S_{21}$ parameters. (a) Amplitude plot of the measured and computed $S$-parameters and power loss. (b) Measured and computed $S$-parameters in Smith chart. Input parameters: Substrate: RT/duroid 5870. $d = 0.7874$ mm, $t_7 = 2.33$ mm. Dimension of the filter: $CV = 2.30$ mm, $G = 0.3$ mm, $L = 10$ mm. Frequency sweep: 9.2 GHz to 10.6 GHz.

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