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Compression limits in cascaded quadratic soliton compressors
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Abstract
Cascaded quadratic soliton compressors generate under optimal conditions few-cycle pulses. Using theory and numerical simulations in a nonlinear crystal suitable for high-energy pulse compression, we address the limits to the compression quality and efficiency.

Introduction
Soliton compressors are attractive because only a single nonlinear medium is needed to achieve many-fold pulse compression. In cascaded quadratic soliton compressors (CQSCs) soliton compression of high-energy fs pulses is possible, and few-cycle pulses can be reached in the near-infrared [1–6]. The cascaded quadratic nonlinearity is achieved by phase-mismatched second-harmonic generation (SHG), where the fundamental wave (FW) experiences a strong nonlinear phase shift from the cyclic energy transfer to the second harmonic (SH). Unique for the CQSC is that solitons exist with normal dispersion, because the effective cubic nonlinearity induced by the cascaded SHG can be made self-defocusing [1]. Consequently the CQSC can compress arbitrarily high-energy pulses, and soliton compression may occur even in the near-infrared where the absence of anomalous dispersion prevents traditional soliton compressors to work. It is well known that in the stationary regime clean compression is possible in the SHG by increasing the oscillatory nature of the FW [3, 5], where $d_{eff}$ is the effective quadratic nonlinearity, and $n_{Kerr}$ is the cubic (Kerr) nonlinear refractive index. Time and the propagation coordinate are normalized to the FW input duration $T_0$ and the FW GVD length $L_{D,1} = T_0^2/|k_1|^2$, and $U_1 = E_1/E_0$ is scaled to the peak input electric field. This generalized nonlinear Schrödinger equation (NLSE) shows that the cascaded quadratic nonlinearity imposes a temporal nonlocal response on the FW, with the nonlocal response functions

$$R_+(\tau) = \frac{\tau_0}{2\tau_0^2} \exp(-is_n\tau/\tau_0) \exp(-|\tau|/\tau_b) \quad (2)$$

$$R_-(\tau) = \frac{\tau_0}{2\tau_0^2} \exp(-is_n\tau/\tau_0) \sin(|\tau|/\tau_b) \quad (3)$$

where the dimensionless nonlocal time scales $\tau_0$ and $\tau_b$ depend on $k_1$ and $\tau_0$ depend on the FW and SH dispersion. According to the nonlinear theory the GVM effects dominate for $\Delta k < d_{g2}^{-1}/2k_2^{(2)}$ : this nonstationary regime is controlled by the oscillatory response function $R_-$. For $\Delta k > d_{g2}^{-1}/2k_2^{(2)}$ the cascaded nonlinearities dominate: this stationary regime is controlled by the local response function $R_+$.

Weakly nonlocal limit
In the weakly nonlocal limit, where the nonlocal response is much faster than the response of $U_1^2$, Eq. (1) can be approximated as [6]

$$[i\partial_\tau - \frac{1}{2} \partial^2_\tau] U_1 - N_{SHG}^2 U_1|U_1|^2 = N_{Kerr}^2 [is_n\tau R_{SHG}|U_1|^2 \partial_\tau U_1 + \frac{1}{2}n_2 s_R U_1^3 \rho(\tau, U_1)] \quad (4)$$

where $s_R = +1$ ($s_R = -1$) in the stationary (non-)stationary regime, and $s_n = sgn(d_{g2}k_2^{(2)})$. The LHS is an NLSE supporting solitons if the effective soliton order $N_{eff} = (N_{SHG}^2 - N_{Kerr}^2)^{1/2}$ is above unity. $N_{eff}$ also controls the compressor performance through the NLSE-like scaling laws [5]. The RHS gathers two detrimental terms: (1) A GVM-induced Raman-like perturbation with a characteristic dimensionless time $\tau_{R,SHG} \equiv 2|d_{g2}|/\Delta k T_0$, (2) A GVM-induced term $U_1^3 \rho(\tau, U_1)$ containing oscillatory components with periods dictated by $\tau_0$ and $\tau_b$, which explains the trailing oscillations often observed in the nonstationary regime, see Fig. 1(a) for $\Delta k = 30$ mm$^{-1}$, $\rho(\tau, U_1)$ is caused by the oscillatory nature of $|R_-|$, so it appears only in the nonstationary regime. The RHS of Eq. (4) holds another insight: for a given $\Delta k$, increasing $N_{SHG}$ by increasing...
the intensity does not necessarily lead to better compression because the Raman-like term, which causes strong pulse asymmetry and soliton splitting, scales as $N_{\text{SHG}}^2$. Similarly in the nonstationary regime, the detrimental oscillatory term $U_1^* \rho(t, U_1)$ also scales as $N_{\text{SHG}}^2$.

**Numerical results and discussion**

Fig. 1 shows numerics all having the same soliton order $N_{\text{eff}} = 8$. Thus, the 200 fs input pulse should in all cases be compressed to 6.0 fs [5]. This is indeed observed in the stationary regime for $\Delta k = 50 \text{ mm}^{-1}$. For larger $\Delta k$, still in the stationary regime, Kerr XPM gradually degrades compression. For smaller $\Delta k$ the transition to the nonstationary regime is approached ($\Delta k = 43 \text{ mm}^{-1}$), where pulse compression is limited by the nonlocal time scale $t_b = \tau_b T_0$. In the nonstationary regime ($\Delta k < 42$), pulse compression degrades due to increasing Raman-like effects ($\tau_{R, \text{SHG}} \propto \Delta k^{-1}$), and to slow trailing oscillations (evident for $\Delta k = 30 \text{ mm}^{-1}$), caused by the GVM-induced oscillatory term $\rho$ in Eq. (4). All simulations have a FW peak around 3 $\mu$m, which is a dispersive wave phase-matched to the FW soliton, causing the fast trailing oscillations for $\Delta k = 50, 43, 41 \text{ mm}^{-1}$; these prevent reaching single-cycle pulses for larger $N_{\text{eff}}$. In the nonstationary regime a distinct red-shifted peak appears in the SH spectrum at a frequency $\Omega_+$ determined by the nonlocal theory. In turn, close to the transition ($\Delta k = 41 \text{ mm}^{-1}$) the FW has a corresponding spectral hole at $\Omega_+$, while further away ($\Delta k = 30 \text{ mm}^{-1}$) it becomes a spectral peak. We show in Fig. 1(d) the red-shifted holes/peaks found numerically versus $\Delta k$, with an impressive agreement with the nonlocal theory.

**Conclusions**

In summary the compression limits in the nonstationary regime are the GVM-induced Raman-like effects and oscillatory components. In the stationary regime the GVM-induced Raman-like effects, nonlocal effects, competing cubic nonlinearities and XPM effects, and dispersive waves, which only exist when taking into account higher-order dispersion, all limit compression.

**References**


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**Fig. 1:** Soliton compression with $N_{\text{eff}} = 8$ of a 200 fs FWHM $\lambda_1 = 1064$ nm pulse in a BBO crystal. (a) FW time plot, (b) the FW and (c) SH spectra at the optimal compression point. (d) The red-shifted spectral peaks in the nonstationary regime from numerics (symbols) and nonlocal theory (lines). The full coupled SHG equations are used, including self-steepening on all nonlinear terms and higher-order dispersion.