Lattice-matched Cu2ZnSnS4/CeO2 solar cell with open circuit voltage boost

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We report a reproducible enhancement of the open circuit voltage in Cu$_2$ZnSnS$_4$ solar cells by introduction of a very thin CeO$_2$ interlayer between the Cu$_2$ZnSnS$_4$ absorber and the conventional CdS buffer. CeO$_2$, a non-toxic earth-abundant compound, has nearly optimal band alignment with Cu$_2$ZnSnS$_4$ and the two materials are lattice-matched within 0.4\%. This makes it possible to achieve an epitaxial interface when growing CeO$_2$ by chemical bath deposition at temperatures as low as 50°C. The open circuit voltage improvement is then attributed to a decrease in the interface recombination rate through formation of a high-quality heterointerface.
Cu$_2$ZnSn(S,Se)$_4$ compounds are among the materials that are currently considered as potentially suitable for terawatt-scale solar energy production. The pure-sulfide material Cu$_2$ZnSnS$_4$ (CZTS) is particularly interesting because it avoids the relatively scarce element Se and it has a larger band gap (1.5 eV) which is appropriate both for a single-junction solar cell and for a top absorber in a tandem solar cell architecture. Even though its highest reported power conversion efficiency so far (9.1%) was achieved with a CdS heterointerface partner, or buffer layer, interface recombination is still the dominant voltage loss mechanism in the solar cell, which suggests that CdS is not the ideal partner of CZTS. Despite promising results achieved with (Zn,Cd)S and (Zn,Sn)O buffer layers, no alternative materials to date have outperformed the highest efficiency or the highest open circuit voltage achieved with a CdS heterojunction partner.

The open circuit voltage $V_{oc}$ of a solar cell limited by interface recombination can be expressed as

$$V_{oc} = \frac{E_i}{q} - \frac{kT}{q} \ln \left( \frac{qS_p N_v}{J_{sc}} \right) \quad (1)$$

for a p-type absorber with an inverted heterointerface, i.e., with holes as minority carriers at the interface. $E_i$ is the interface band gap, $q$ is the elementary charge, $k$ is Boltzmann’s constant, $T$ is temperature, $N_v$ is the effective density of states in the valence band, $J_{sc}$ is the short circuit current, and $S_p$ is the interface recombination velocity of holes. By modeling interface recombination as Shockley-Read-Hall recombination through a single interface defect level, $S_p$ can be expressed as

$$S_p = N_i \sigma_p v_{th} \quad (2)$$

where $N_i$ is the area density of interface defects, $\sigma_p$ is a cross section describing how efficiently the holes are captured by the defect, and $v_{th}$ is the thermal velocity.

Eqs. 1, 2 show that decreasing the interface defect density $N_i$ can be a way to improve the open circuit voltage in a solar cell limited by interface recombination. By analogy to high-efficiency solar cell technology based on III-V semiconductors, this can be achieved by ensuring epitaxial growth of the buffer material on the absorber material. In the absence of epitaxial growth, a large density of atomic dislocations will exist at the interface, which leads to formation of allowed electronic states within the interface band gap, thus increasing $N_i$ and enhancing interfacial Shockley-Read-Hall recombination. Epitaxial growth is facilitated by a small lattice mismatch between the two heterojunction materials. While CdS has
a reasonably small lattice mismatch with some absorber materials such as Cu(In,Ga)Se$_2$ ($\sim 1.5\%$) and Cu$_2$ZnSnSe$_4$ ($\sim 2.4\%$), its mismatch with CZTS is much larger ($\sim 7\%$). A high value of $N_i$ can therefore be anticipated at the CZTS/CdS interface.

To address this problem, we carried out an initial investigation of the non-toxic, earth abundant material CeO$_2$ as a heterojunction partner of CZTS. CeO$_2$ has a nearly perfect lattice match with CZTS (lattice constant of CZTS: 5.43 Å; lattice constant of CeO$_2$: 5.41 Å; thus a lattice mismatch under 0.4%). This opens the possibility for a high-quality epitaxial interface, which may alleviate the interface recombination problem. CeO$_2$ also has a wide indirect band gap of 3.3 eV, which minimizes parasitic light absorption. We note that, based on Eq. 1, the lattice mismatch at the CeO$_2$/CdS interface is not expected to limit $V_{oc}$, since $E_i$ is much larger at the CeO$_2$/CdS interface than at the CZTS/CeO$_2$ interface.

CZTS films were prepared by co-sputtering Cu/ZnS/SnS precursors on Mo-coated soda lime glass using a magnetron sputtering system (AJA International, Inc., model ATC-2200) as presented in detail before. Two types of elemental compositions were targeted in different CZTS precursors within the same batch: the first (C1) with Cu/Sn = 1.9 and Zn/Sn = 1.25; the second (C2) with Cu/Sn = 1.8 and Zn/Sn = 1.25. C1 and C2 were chosen because the final solar cell efficiency is very sensitive to the sample stoichiometry, especially to the Cu/Sn ratio. These precursors were then annealed using Rapid Thermal Processor (AS-One 100) in a S- and Sn-containing atmosphere at 560°C. The S atmosphere compensates for S substoichiometry in the precursors and the addition of Sn stabilizes CZTS against decomposition reactions. CeO$_2$ films were deposited on CZTS by chemical bath deposition (CBD) at 50°C in a weakly acidic solution (pH: 6) containing 10 mM Ce(CH$_3$COO)$_3$ and 5 mM KClO$_3$ in Milli-Q water under mild stirring. The expected chemical reactions resulting in thin film deposition are those presented in previous work. Cds (60 nm) was deposited by CBD with a previously described process. ZnO (60 nm) and indium tin oxide (ITO, 200 nm) layers were deposited by RF magnetron sputtering, followed by an Al contact grid. No anti-reflection coating was applied. The total area of the final cells (0.23 cm$^2$) was defined by mechanical scribing. Three different solar cell architectures were fabricated in this work as shown in Fig. 1. In the first architecture (A1), CdS is completely replaced by a stand-alone 30 nm CeO$_2$ buffer layer deposited with a 100 min process. In the second architecture (A2), a thin CeO$_2$ layer of estimated thickness between 1 and 5 nm is inserted between
FIG. 1. Schematic drawing of the A1 architecture with a 30 nm stand-alone CeO$_2$ buffer layer and the A2 architecture with a 1-5 nm CeO$_2$ interlayer between CZTS and CdS. The reference architecture is a standard CZTS solar cell structure.

CZTS and the standard CdS buffer layer using a 10-20 min deposition process. The reference architecture (Fig. 1) is a conventional CZTS solar cell structure without CeO$_2$. Only very weak n-type conductivity has been reported$^{18}$ for CeO$_2$ and, in fact, the resistivity of the films synthesized in this work was too high to be measured with conventional four-point probe apparatus. Therefore, the CeO$_2$ layer in the solar cell can be regarded as completely depleted. In the case of architecture A2, the CeO$_2$ layer is so thin that it is not expected to modify significantly the original electric field profile of the reference architecture.

Scanning electron microscope images were taken with a FEI-Nova NanoSEM 450 instrument at 2 kV beam voltage. High-resolution bright-field transmission electron microscope images (HRTEM) were taken with a FEI-Titan 80-300 TEM, at 300 kV beam voltage. X-ray photoelectron spectroscopy (XPS) was performed with a Thermo Scientific K-Alpha instrument with a monochromatized Al K$_\alpha$ x-ray source. The binding energy scale was calibrated using the adventitious C 1s peak at 284.8 eV. An electron flood source was employed to limit charging effects in CeO$_2$.$^{19}$ CeO$_2$ was found by others to be reduced to Ce$_2$O$_3$ both under prolonged x-ray exposure$^{20}$ and under ion beam sputtering. $^{19}$ Indeed, by inspection of the Ce 3d spectrum and of the valence band edge, reduction effects were evident in our films even at the lowest ion beam energy (200 eV) available in the XPS setup. Therefore no sputter cleaning was performed on our samples. XPS data analysis was performed with the Avantage 5.948 software (Thermo Scientific). Current-voltage (JV) characteristics on
finished devices were measured under AM 1.5G illumination with a solar simulator from PV Measurement and a Keithley 2400 source meter calibrated with a standard Si reference.

Before discussing solar cell results, we want to answer some basic questions about the quality of deposited CeO\(_2\) films. The first question is adhesion. By depositing CeO\(_2\) on glass/Mo substrates, it was found that the films deposited from a solution with a KClO\(_3\) concentration of 10 mM or above were easily peeled off by scotch tape. However, already at a KClO\(_3\) concentration of 5 mM, the films were strongly adherent with no peel-off by repeated scotch tape application. A KClO\(_3\) concentration of 5 mM was therefore chosen for the deposition process. The maximum CeO\(_2\) thickness on CZTS that could be achieved in a single chemical bath deposition run was about 30 nm with a 100 min process (Fig. S1(a)).

The second question is whether the deposited film indeed consists of the desired CeO\(_2\) material. The fast Fourier transform (FFT) of cross-sectional TEM images of the deposited film yields a pattern that is compatible with CeO\(_2\) and not with the main competing phase Ce\(_2\)O\(_3\) (Fig. 2(c)). However, TEM analysis involves very small regions (few nm). Raman spectroscopy was then performed over a much larger analysis area (about 2 \(\mu\)m diameter). Only one additional Raman peak at 461 cm\(^{-1}\) was revealed on CZTS/CeO\(_2\) bilayers compared to a bare CZTS spectrum (Fig. S2). The peak corresponds to the first-order-allowed Raman mode of CeO\(_2\), with a small red shift due to size effects.\(^{21}\) Finally, XPS characterization was performed over an even larger analysis area (about 400 \(\mu\)m diameter). All the XPS peaks corresponding to Ce 3d core levels were fitted and attributed to either CeO\(_2\) or Ce\(_2\)O\(_3\) according to reference spectra\(^{22}\) (Fig. S3). The fraction of CeO\(_2\) present in the deposited film is estimated as 70.4\% with this method. Hence, some Ce\(_2\)O\(_3\) inclusion should be expected. Since Ce\(_2\)O\(_3\) is not lattice-matched to CZTS and it has a band gap\(^{13}\) almost 1 eV lower than that of CeO\(_2\), we assume that Ce\(_2\)O\(_3\) inclusions promote interface recombination and reduce the open circuit voltage enhancement that could be achieved with a pure CeO\(_2\) layer.

The third question is whether the deposited film provides complete coverage of the underlying CZTS layer. SEM images of a 30 nm-thick CeO\(_2\) film (Fig. 3) show some non-uniform coverage in correspondence of CZTS grain boundaries and some smaller isolated dips in the CeO\(_2\) film profile, which may be interpreted as pinholes. The area fraction of CZTS that is covered by CeO\(_2\) can be estimated by comparing the peak intensity of Ce, Cu, Zn, and Sn core levels by XPS, as explained in the Supplementary Material. The covered area is then estimated as 62\%, 94\% and 95\% for the films deposited for 20, 60, and 90 min, with
FIG. 2. (a) Cross-sectional HRTEM image across the interface between CeO$_2$ (top) and CZTS (bottom). (b) FFT of a selected region across the interface. A reflection at 3.13 Å is visible, corresponding to the distance between (111) planes of CeO$_2$ and between the (112) planes of CZTS. (c) FFT of a selected region within CeO$_2$. The diffractogram can be indexed as a CeO$_2$ crystal along with [103] as zone axis. (d) FFT of a selected region within CZTS. Reflections corresponding to the (112) and (310) planes are visible.

estimated thicknesses 1-5, 15, and 25 nm, respectively. Thus, it seems as if the very thin films (20 min deposition time) employed in architecture A2 may have rather poor coverage. However, in that case the CeO$_2$ thickness is comparable to the XPS probing depth (about 2 nm). Therefore, part of the CZTS signal is likely to originate from CZTS buried under CeO$_2$ so that the covered fraction is in reality higher. The consequence of incomplete coverage is simply the coexistence of CZTS/CeO$_2$ and CZTS/CdS heterojunctions in parallel
FIG. 3. Cross sectional SEM image of a CeO$_2$ film grown on CZTS used for architecture A1. The film is about 30 nm thick, which is the maximum thickness that could be achieved with a single CBD process in this work. Some scattered pinholes and inhomogeneous coverage of the grain boundaries are seen.

in the solar cells.

The fourth question is whether CeO$_2$ forms an epitaxial interface with CZTS, as may be expected by their excellent lattice match. The HRTEM image in Fig. 2(a) shows that the atomic arrangement on the CZTS side continues on the CeO$_2$ side for a few nm, thus demonstrating that an epitaxial interface can be achieved despite the low deposition temperature. Analysis of the FFT pattern across the interface (Fig. 2(b)) reveals that epitaxy occurs along the CZTS(112)-CeO$_2$(111) direction, similarly to what was recently observed on a CZTS/CdS interface. A few nm away from the interface, CeO$_2$ is nanocrystalline with average grain size less than 5 nm, as inferred from TEM images (Fig. 2(a)) and by quantitative analysis of Raman peak broadening (Fig. S2). We emphasize, however, that epitaxial growth was not observed at most interface locations imaged by TEM in this work, as shown for example in Fig. S1(b).

The fifth and final question is the band alignment of CeO$_2$ with CZTS. It is well known that a moderate spike-like conduction band offset (CBO) at the heterointerface, in the 0-0.4 eV range, is optimal as it reduces interface recombination without blocking photocurrent transport. The valence band offset (VBO) was estimated by XPS as shown in Fig. 4. The experimental (bulk) band gaps of the two materials were then added to the VBO to finally obtain a CBO of $-0.12 \pm 0.20$ eV. For the CeO$_2$ indirect gap, we used a value of $3.3 \pm 0.1$ eV measured by internal photoemission yield. The extracted CBO is only slightly below the
FIG. 4. CZTS/CeO$_2$ band alignment extracted by XPS. The valence band maximum (VBM) of CZTS with respect to the Fermi level is $-0.60 \pm 0.10$ eV in a bare CZTS sample and $-0.75 \pm 0.10$ eV in a CZTS sample coated with a thin CeO$_2$ layer (20 min deposition time, red circles). From this, band bending (BB) in CZTS is estimated as $0.15 \pm 0.10$ eV. This is also confirmed from the average shift in the Cu 2p, Zn 2p, and Sn 3d core levels between the bare CZTS sample and the CeO$_2$/CZTS sample. Shifting the XPS spectrum of the bare CZTS sample by the BB (black line) allows deconvolution of the CeO$_2$ valence band signal (blue triangles), located $-2.67 \pm 0.10$ eV below the Fermi level. The shifted spectrum of bare CZTS (black line) fits well with the CeO$_2$/CZTS spectrum until the onset of the CeO$_2$ valence band, as expected.

The current-voltage characteristics of the three best solar cells fabricated with the three different architectures are shown in Fig. 5. A small improvement in open circuit voltage is achieved by architecture A2 with respect to the reference architecture. The accompanying

optimal range and it is actually more favorable than the previously measured CZTS/CdS band offset.$^{25}$ Considering the many possible sources of error in the measurement, including the use of as-deposited surfaces for analysis, this is considered a promising result. We also note that the CeO$_2$ Fermi level lies about 2.7 eV above the valence band, which indicates that CeO$_2$ is n-type in the analysis region, similarly to previous reports.$^{22,26}$
FIG. 5. Illuminated current-voltage characteristics of the best solar cells with the A1 and A2 architectures, together with the best solar cell with the reference architecture.

degradation in short circuit current will be discussed in the following. Interestingly, the statistics shown in Fig. 6 indicate that the $V_{oc}$ boost is reproducible over two separate batches (i) and (ii) regardless of precursor composition (C1 or C2). In all those cases, the open circuit voltage of the solar cells with A2 architecture is always between 20 mV and 100 mV higher than in the corresponding solar cells with the standard CdS architecture. Additional statistics in Fig. S6 confirm this trend. The highest open circuit voltage achieved in this study with the A2 architecture was 641 mV, and the highest efficiency was 6.6% (Fig. 5). Conversely, complete current blocking and no photovoltaic effect was observed in the solar cells with a stand-alone CeO$_2$ buffer layer (architecture A1 in Fig. 5).

To interpret these results, we refer to the band structure of CeO$_2$. What has been referred to as the "conduction band" in this work is a band consisting of highly localized Ce 4f states, 3.3 eV above the valence band as mentioned before. However, localization of states in this band implies that the electron effective mass is very high, and hence the electron mobility is very low. This can be inferred by a nearly complete lack of dispersion of those states in reciprocal space. For this reason, a band consisting mostly of Ce 5d states,
which lies about 6 eV above the valence band and has lighter effective masses, is often quoted as the conduction band of CeO$_2$. From our band alignment study (Fig. 4), we found that the 4f band has a nearly optimal CBO with the CZTS conduction band, while the 5d band lies at much higher energy. Then, we can attribute the complete lack of photocurrent in architecture A1 to the very poor transport properties of the 4f band. A 30 nm CeO$_2$ layer is therefore thick enough to completely impede electron transport. Instead, when CeO$_2$ is very thin as in architecture A2 (1-5 nm), short circuit current and fill factor losses are greatly diminished and in some cases eliminated (Fig. 5) and Fig. S5). This is possibly due to tunneling-based transport between the CZTS and CdS conduction bands through the thin interlayer. Residual current losses, as in Fig. 5, are probably due to thickness inhomogeneity of CeO$_2$, or to the fact that an even thinner layer is necessary. Growth of CeO$_2$ by atomic layer deposition may be beneficial in this respect.

We conclude that inserting a thin lattice-matched CeO$_2$ layer between CZTS and CdS alleviates interface recombination and results in a reproducible open circuit voltage boost in
the solar cell. This is attributed to the formation of a high-quality CZTS/CeO$_2$ heterointerface, with instances of epitaxial growth observed in some regions. We suggest that the open circuit voltage could be further improved if epitaxy could be obtained on a larger scale, if Ce$_2$O$_3$ inclusions could be decreased, and if surface coverage could be improved. Etching the CZTS surface immediately prior to CeO$_2$ deposition may facilitate formation of an epitaxial interface. It should be kept in mind that the very large electron effective mass in the CeO$_2$ conduction band puts a severe constraint on the maximum thickness of the CeO$_2$ film, which should only be a couple of nm thick in order to avoid dramatic current losses.

SUPPLEMENTARY MATERIAL

See supplementary material for additional TEM images, phase analysis of the CeO$_2$ layer, compositional analysis across the interface, estimation of CeO$_2$ coverage, and more detailed statistics on the solar cell parameters for the reference architecture and architecture A2.

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REFERENCES

1M. A. Green, Nat. Energy 1, 1 (2016).
