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FIELD AMPLIFICATION, VORTEX FORMATION, AND
ELECTRON ACCELERATION IN A PLASMA PROTOSHOCK:
EFFECT OF ASYMMETRIC DENSITY PROFILE

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Gamma ray bursts (GRBs) are thought to originate from highly relativistic jets. The
fireball model predicts internal shocks in the jets, causing magnetic field to be amplified
& particles to be accelerated. We model the effects of an asymmetric density configuration
for an internal plasma collision in a quasi-parallel magnetic field. We measured electron
acceleration & found that a tenuous population of electrons is accelerated to Lorentz
factors of $\sim 300$ - close to energy equipartition with ions. We found that the filaments
did not remain static, but were deflected by the Lorentz force & rolled up into small
vortices, which themselves merge to form a larger current vortex.

Keywords: Shocks; gamma-ray bursts; simulations.

1. Introduction

Prompt emissions of the ultrarelativistic GRBs are probably the most energetic
radiative events in the universe. Since the first observations of GRBs, thousands
have been detected. They all share a common signature of energetic non-thermal
radiation attributed to relativistic motion of electrons in strong magnetic field. The
fireball model predicts internal shocks in the jets, causing magnetic field to be
amplified & particles to be accelerated$^{1,2}$. The details of the underlying physical
mechanism that causes the prompt emissions are still unknown. In particular, how
magnetic field is spontaneously generated, & how electrons are energised to rela-
tivistic speeds & injected into the Fermi mechanism, is under debate. It is therefore
relevant to use numerical simulations to probe the behavior of magnetized shocks
to see if a robust mechanism to self-consistently amplify the field may be found.
Current models focus on exploiting the symmetry of uniform density plasma cloud collision, in this work, we extend the survey to unequal density plasma clouds.

Two plasma clouds, each consisting of ions and electrons with the mass ratio \(m_i/m_e = 250\), collide at position \(x = 0\), Figure 1 shows the simulation setup. Initial electron and ion number densities of the dense cloud both equal \(n_1\) and those of the tenuous cloud equal \(n_2 = n_1/10\). The velocity vectors of both clouds are antiparallel and aligned with \(x\). The modulus \(v_b\) of each cloud in the box frame gives the collision speed \(v_c = 2v_b/(1 + v_b^2/c^2) = 0.9c\). The dense cloud propagates to increasing values of \(x\). All species have an initial temperature of 131 keV. The magnetic energy to baryonic energy ratio, \(\sigma\), is 0.2 (0.02) in the tenuous (dense) cloud. The modulus of the convection electric field is \(|E_0| = v_bB_0/c\). Both \(\nabla \cdot B = 0\) and \(\nabla \cdot E = 0\) at \(t = 0\). The simulation box size \(L_x \times L_y = 656 \lambda_i \times 6 \lambda_i\) is resolved by \(2.8 \times 10^4 \times 256\) square grid cells. The dense (tenuous) plasma is resolved by 100 (50) particles per cell per species.

2. Results

**Vortex formation in plasma shocks:** We carried out a large scale, long term 2d relativistic PIC simulation of a plasma collision. The peak Lorentz factor of the electrons is determined, along with the orientation & the strength of the magnetic field at the cloud collision boundary. The magnetic field component orthogonal to the initial plasma flow direction is amplified to values that exceed those expected from the shock compression by over an order of magnitude. The forming shock is quasi-perpendicular due to this amplification, caused by a current sheet which develops in response to the differing deflection of the upstream electrons & ions incident on the magnetised shock transition layer. As the upstream plasma impacts on the strong magnetic field, electrons are deflected away from their original flow direction, while the ion reaction is much weaker. Current flows in the \((y,z)\) plane, which amplifies magnetic field. The electrons fall behind the ions, because their velocity along \(x\) is reduced. An \(E_x > 0\) builds up, which tries to restore the quasi-neutrality. Incoming upstream electrons are dragged by it across the magnetic field and accelerated to relativistic speeds. The distribution of electrons has a non thermal tail with a power law \(\sim 2\). A magnetic field structure resembling the cross section of a flux tube grows

Fig. 1. Left: Simulation setup. Right: formation of current vortex at collision boundary.
self-consistently in the current sheet of the shock transition layer (Fig. 1). Plasma filamentation develops behind the shock front, as well as signatures of orthogonal magnetic field striping, indicative of the filamentation instability. These magnetic fields convect away from the shock boundary & their energy density exceeds by far the thermal pressure of the plasma. Localized magnetic bubbles form. Energy equipartition between the ion, electron & magnetic energy is obtained at the shock transition layer. The electronic radiation can provide a seed photon population that can be energized by secondary processes (e.g. inverse Compton). We measured electron acceleration & found that a tenuous population of electrons is accelerated to Lorentz factors of $\sim 200$ - close to energy equipartition with ions.

We found that the filaments did not remain static, but were deflected by the Lorentz force & rolled up into small vortices, which themselves merge to form a larger current vortex. In order to check the validity of the 2d approximation we also carried out a short-term simulation in 3d & found evidence of filamentation.

**Long term 1D simulations:** We carried out long term simulations in 1D in order to verify that the acceleration & field amplification process was robust & predict the SSC emission. In 1D, the simulation timescale was increased to $919 \omega^{-1}_p$. In the 1d simulation, a reverse shock forms, which was not achieved in 2D due to the shorter timescales. The reverse shock is visible at $x = 350$ in Fig. 1. Two smaller discontinuities are seen in the ion phase space plot at $x = 400, 500$.

**Synthetic observations:** In order to analyze the data, a model of the emission from the hot spot was created using the Compton Sphere Suite$^5$. In this one-zone

![Ion Phase Space](image1)

![Electron Phase Space](image2)

Fig. 2. Upper panel: Specific x-momentum vs displacement phase space for ion clouds Lower panel: Lorentz factor vs displacement phase space for electron species
Fig. 3. Predicted synchrotron self-Compton emission, assuming Doppler factor of 450, magnetic field of 20 nT, maximum Lorentz factor of 1500 using the Compton sphere suite.

steady-state model of secondary emission processes, a sphere of a given radius is permeated by a magnetic field. Photons are produced relativistic electrons via the synchrotron process & are inverse-Compton scattered by the same electron population to γ-ray energies. The model assumes synchrotron losses dominate over SSC losses. The Thomson scattering cross section is used. The values from the numerical simulation are used, but must be scaled up from the reduced mass ratio to the true value. In Fig. 3 we plot the derived spectrum. For future work we shall include the external photon field inverse Compton losses.

Conclusions: In 2D simulations the F.I. grows despite the high temperatures & quasi-parallel background magnetic field (which tends to suppress transverse motion). The quasi-parallel field is rotated into a perpendicular field locally at the forward shock transition layer. Longer term 1D simulations show similar magnetic field amplification and electron acceleration, while also evolving for sufficient time to see a reverse shock forming. The synthetic observations show emission in the MeV range.

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