A CALIBRATION METHOD FOR PLLS BASED ON TRANSIENT RESPONSE

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ABSTRACT

A novel method to calibrate the frequency response of a Phase-Locked Loop (PLL) is presented. The method requires just an additional digital counter and an auxiliary Phase-Frequency Detector (PFD) to measure the natural frequency of the PLL. The measured value can be used to tune the PLL response to the desired value. The method is demonstrated mathematically on a typical PLL topology and it is extended to ΣΔ fractional-N PLLs. A set of simulations performed with two different simulators is used to verify the applicability of the method.

1. INTRODUCTION

Phase-Locked Loop (PLL) frequency synthesizers are building blocks of all communication systems. An accurate PLL response is required in many situations, especially when ΣΔ PLLs for direct modulation [1] are used. In these types of PLLs, the data fed into the ΣΔ modulator is often undergoing a pre-filtering process in order to cancel the low-pass PLL transfer function and thereby to extend the modulation bandwidth [2]. The pre-distortion filter presents a transfer function equal to the inverse of the PLL transfer function and it is usually implemented digitally. Consequently, a tight matching between the pre-distortion filter and the analogue PLL transfer function is necessary to avoid distortion of the transmitted data.

Especially for on-chip Voltage Controlled Oscillators (VCO), the gain \( K_{VCO} \) is typically the parameter with the poorest accuracy among the PLL analog components. However to establish an accurate PLL transfer function only the product \( K_{VCO} \times I_{CP}/C \) needs to be accurate [3]. The PLL can then be calibrated by adjusting the Charge-Pump current; the problem is how to measure the accuracy of the PLL transfer function.

A continuous calibration technique is presented in [4]. The transmitted data is digitally compared with the input data and the Charge-Pump current is then adjusted to compensate the detected error. This method offers the possibility of continuous calibration at the expense of increased circuit complexity; since the error detection is based on the cross-correlation between input and transmitted data, this approach will not work on unmodulated synthesizers. An alternative approach is found in [5], where a method based on the detection of pulse skipping is described. The presence of one or several pulse skips can be used as an indication of the bandwidth. This method requires an input frequency step large enough to push the PLL into its non-linear operating region and only offers a rough estimation of the actual PLL bandwidth.

In this paper we present a simple and novel approach that makes it possible to determine the characteristics of the PLL transfer function by simply adding a digital counter and an auxiliary Phase Frequency Detector (PFD).

The paper is organized as follows: in section 2 we present the basic idea behind the method and in section 3 we discuss its mathematical formulation. The extension of the method to ΣΔ PLLs is presented in section 4. Finally, in section 5, the results from different simulations are compared with the theory developed.

2. MEASURING SCHEME

A calibration cycle is required by this method. Consider the typical PLL topology in fig. 1: to start the calibration, the switch to \( R_{cal} \) is closed and the calibration resistor \( R_{cal} \) is connected to the resistor \( R \). By reducing the total filter resistance, the loop transfer function presents under-damped characteristics. Note that, if the loop transfer function is already designed with under-damped characteristics, then the calibration switch is not necessary. By changing the ratio \( M \) of the \( f_{ref} \) divider or of the ratio \( N \) of the \( f_{out} \) divider, a frequency step can be applied to the PLL. The natural frequency of the induced transient response can be indirectly measured by counting the UP/DOWN pulses produced by the PFD. If the counter counts 1 up for each UP pulse and counts 1 down for each DOWN pulse, then the maximum counter value is a measure of the natural frequency of the PLL transfer function. This can be seen in fig. 2, where the expected behavior of the phase error together with the counter value trajectory are presented. The Charge-Pump current can be adjusted so that the PLL natural frequency is the desired one.

For the calibration method to work properly, the leakage current in the Charge-Pump should be kept small and the trickle cur-
rent (if any) should be turned off. Any leakage or trickle currents will induce a static phase error at the PFD input. This, in turn, means an increased number of pulses in one direction (e.g., UP pulses). Consequently, the counter value is no longer an accurate representation of the natural frequency.

The auxiliary PFD in fig. 1 is required to generate stable UP/DOWN pulses for the digital counter. A possible circuit implementation that works together with a typical PFD is shown in fig. 3. The two set-reset flip-flops (SR-FF) are used to establish which one between the UP/DOWN pulses occurs first. This is necessary because the UP and DOWN pulses are simultaneously high for a length equal to the delay in the PFD reset path [5]. If the UP pulse rises before the DOWN pulse, then a logical ‘ONE’ appears at the input of the top edge-triggered resettable D flip-flop (fig. 3) and a logical ‘ZERO’ appears at the input of the bottom D flip-flop. The UP pulse delayed through a couple of inverters clocks the flip-flop and the negative transition of the REF clock resets the flip-flop. Hence the flip-flop produces an UP\_stable pulse whose length is approximately equal to the REF semiperiod.

The opposite happens if the DOWN pulse occurs before the UP pulse. If the PFD produces aligned UP and DOWN pulses (this is the case if the input phase error is smaller than the dead-zone of the PFD) then the UP\_stable and the DOWN\_stable signals are high at the same time.

3. MATHEMATICAL DERIVATION

The mathematical formulation will be based on the PLL topology of fig. 1; however, the applicability of the method extends to other topologies, as it will be demonstrated in the next section. We start by deriving the PLL loop transfer function with the aid of the linear model of fig. 4:

\[
H_{\text{loop}}(s) = \frac{I_{\text{ref}}(R \cdot C_p s + 1)K_{\text{VCO}}}{2\pi C_p s^2 N}
\]

The transfer function from phase input to phase error is given by:

\[
\frac{\Phi_{\text{err}}(s)}{\Phi_{\text{in}}(s)} = \frac{1}{1 + H_{\text{loop}}(s)} = \frac{s^2}{s^2 + 2\zeta\omega_n s + \omega_n^2}
\]

where \(\omega_n = \sqrt{\frac{L C_p K_{\text{VCO}}}{2\pi C_p N_d}}\) and \(\zeta = \frac{R}{2}\sqrt{\frac{L C_p K_{\text{VCO}}}{2\pi N}}\). The transfer function is determined by the natural frequency and the damping factor of the PLL. A unit input frequency step corresponds to an input phase ramp, with Laplace transform given by \(\Phi_{\text{in}}(s) = \frac{1}{\omega_n s}\). The Laplace transform of the phase error is then given by:

\[
\Phi_{\text{err}}(s) = \frac{s^2}{s^2 + 2\zeta\omega_n s + \omega_n^2} \cdot \frac{1}{s^2}
\]

The behavior in the time domain of equation 3 is the impulse response of a second order system:

\[
\phi_{\text{err}}(t) = \frac{1}{\omega_n \sqrt{1 - \zeta^2}} e^{-\zeta\omega_n t} \sin(\omega_0 t)
\]

where \(\omega_0 = \omega_n \sqrt{(1 - \zeta^2)}\). As long as the phase error function \(\phi_{\text{err}}(t)\) is positive, the PFD generates UP pulses. If the error becomes negative, DOWN pulses are generated. Assuming a positive frequency step, an initial sequence of UP pulses is produced by the PFD and the counter value increases monotonically. When the phase error crosses the zero-error phase, occurring at \(t_{\text{cross}} = \frac{\pi}{\omega_0}\) according to equation 4, DOWN pulses start to appear decreasing the counter value. Hence at the crossing time the counter reaches its maximum value, \(V_{\text{max}}\); the crossing point can also be expressed as \(t_{\text{cross}} = V_{\text{max}} \cdot T_{\text{ref}}\), where \(T_{\text{ref}}\) is the period of the REF clock (fig. 1). For stability reason [3], the PLL dynamics is always much slower than the REF signal: the error introduced by quantizing \(t_{\text{cross}}\) is then insignificant. \(\omega_0\) can then be approximated as:

\[
\omega_0 = \frac{\pi}{V_{\text{max}} \cdot T_{\text{ref}}}
\]
4. EXTENSION TO Σ∆ PLL TOPOLOGIES

The applicability of the measuring technique will be now demonstrated for Σ∆ fractional-N PLLs. These PLLs use high-order multi-bit Σ∆ modulators to dither the divider modulus; direct modulation is achieved by feeding the data into the modulator. Thus, to measure ωo, the frequency step is, in this case, applied to the Σ∆ modulator input.

The linear model of a Σ∆ fractional-N PLL is shown in fig. 5. A complete derivation of the linear model can be found in [6]. The Loop Filter is typically a high-order structure to attenuate the high-frequency quantization noise. In the example analyzed in this section, the Loop Filter transfer function presents 4 poles (including the Charge-Pump integration) and 1 zero. The mathematics involved in this case is lengthier, but the final transfer function can be reduced to an equivalent 2nd order equation.

We start by finding the transfer function from the Σ∆ modulator input to phase error Φerr (s). With the aid of fig. 5, considering that the effect of the Σ∆ modulator Signal Transfer Function (STF) is just adding a delay to the input data, the transfer function is given by:

\[
\Phi_{err}(s) = \frac{2\pi N + \mu_b}{s + \zeta_1} \cdot e^{-sT_{ref}} \frac{1}{1 - e^{-sT_{ref}} + H_{loop}(s)}
\]  

(6)

where \( N + \mu_b \) is the instantaneous divider ratio.

Fig. 5. Σ∆ fractional-N PLL linear model.

Indicating with \( \omega_1, \omega_4, \omega_5 \) and \( z_1 \) the high order poles and, respectively, the zero of the Loop Filter, and by setting:

\[
T_{eq}^2 = \frac{1}{\omega_1^2} + \frac{1}{\omega_4 \cdot \omega_3} + \frac{1}{\omega_4 \cdot \omega_5} + \frac{1}{\omega_3 \cdot \omega_5} + \frac{1}{\omega_3^2}
\]

\[
- \frac{1}{\omega_4 \cdot z_1} - \frac{1}{\omega_5 \cdot z_1} - \frac{1}{\omega_1 \cdot z_1}
\]

it is possible to define an equivalent natural frequency \( \omega_{n1} \) and an equivalent damping factor \( \zeta_1 \):

\[
\omega_{n1} := \sqrt{\frac{\omega_3^2}{\left[ 1 + \left( \omega_1 T_{eq} \right)^2 \right]}}
\]

(7)

\[
\zeta_1 = \frac{1}{2} \frac{\omega_{n1}}{z_1} - \frac{1}{\omega_3} - \frac{1}{\omega_4} - \frac{1}{\omega_5}
\]

(8)

After proper manipulations, eq. 6 can be reduced to the following approximated expression:

\[
\Phi_{err}(s) = \frac{G \cdot \omega_3^2}{s^2 + 2\zeta_1 \omega_{n1} s + \omega_{n1}^2}
\]

(9)

with the gain factor given by:

\[
G = \left( \frac{\omega_1 \mu_b}{\omega_o} \right)^2 \frac{\omega_3}{s T_{ref}}
\]

(10)

and it corresponds directly to equation 3. The smaller is the loop damping factor \( \zeta \), the more accurate is the approximation in equation 10. Thus, the natural frequency can be calculated as described in section 3.

However, the use of Σ∆ fractional-N PLLs introduces a new requirement for the correct applicability of the method. The input step to the modulator needs to be large enough to overcome the random effects of the modulator itself. If the input step is too small, then the UP sequence is no longer monotonic and the extracted value of \( \omega_o \) is no longer accurate. On the other hand, the equations derived so far are based on the assumption that the PLL is working in its linear region. If a large input step is applied, the PLL may be pushed out of its linear region. In this case, the previous equations are no longer valid, but, as mentioned in [7], the calibration method will still work, since the final counter value can be extracted from simulations. Alternatively, the counter can be reset whenever a pulse skip is detected: in this way the final value in the counter represents the number of UP (or DOWN) pulses occurred during the linear part of the step response. In other words, the counter starts to operate properly when the PLL leaves the frequency acquisition mode and enters the phase acquisition mode.

5. SIMULATION RESULTS

The main parameters of the simulated PLL are resumed in table 1. The Σ∆ modulator is a MASH 4th order and the parameters of the Loop Filter are presented in table 2. Based on the simulation results, the Σ∆ fractional-N PLL topology can be simulated with a linear model such as Simulink; however the system behavior has been also investigated through a Verilog implementation. The use of Verilog provides the possibility to simulate a model more close to a real PLL implementation, capable of capturing the non-linear behavior of the system [6].

As already discussed in the introduction, the VCO gain \( K_{VCO} \) is the parameter with the poorest accuracy; the Σ∆ PLL was simulated with the nominal \( K_{VCO} \) value and with a gain variation of ± 30% with respect to the nominal value. In fig. 6 the counter behavior for the 3 different \( K_{VCO} \) values is presented. It is apparent that the three curves reach different peaks according to the value of \( K_{VCO} \); as the time proceeds the effects of the Σ∆ modulator start to appear.

By substituting the values of the parameters in the equations presented in section 3, the theoretical maximum counter values for the three different VCO gains, are, respectively, 150, 123, and 106. The values extracted from the Verilog simulation of fig. 6 are 148, 122, and 105; these values closely match the predicted ones. The counter behavior for different input frequency steps is presented in fig. 7. As the step is increased, overcoming the Σ∆ modulator noise, the measured values match very well the predicted ones. In the same figure the results for both simulators, Verilog and Simulink, are presented. Notice that the Simulink curves are
very close to the curves obtained with Verilog, which confirms that
the linear simulator describes accurately the transient behavior of
the ΣΔ PLL even for fairly large input frequency steps.

As previously discussed, the presence of a leakage current will
result in an average phase error different from zero. If the leakage
current is too large, even in lock condition of the PLL, the PFD will
only produce UP pulses (for a negative leakage current) or DOWN
pulses (for a positive leakage current). The simulations show that
the calibration method is very robust to leakage currents: a
±1% leakage current will produce less than ±6% deviation from the
nominal ω0. By observing the difference between the maximum
number of UP (DOWN) pulses and the maximum number of fol-
lowing DOWN (UP) pulses, it is actually possible to obtain the
polarity and a magnitude estimation of the static phase offset. In
fact, with zero static phase offset the length of the two sequence
would be the same; if a positive phase offset is present, the phase
error curve for a positive frequency step is shifted up with respect
to the zero offset curve: in this case the monotonic sequence of
UP pulses will be longer than the following sequence of DOWN
pulses.

### 6. CONCLUSIONS

A new method to calibrate the PLL transfer-function has been pre-
icted. The implementation does not require any additional ana-
logue component. The only extra circuitry necessary is an auxil-
ary PFD and a digital counter. This new approach does not offer
continuous calibration and it requires a calibration cycle, but it is
very simple and virtually no extra silicon area and no extra power
consumption is required. Moreover, this technique works for both
linear and non-linear PLL frequency step responses; also, it can be
used to estimate the static phase offset. The mathematical formul-
ation of the method has been verified with simulations based on a
ΣΔ fractional-N PLL topology, run both on Verilog and Simulink.
Results from both simulations closely match the theoretical values.

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