Compressibility of the high-pressure rocksalt phase of ZnO

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We report the results of a combined experimental and theoretical investigation on the stability and the volume behavior under hydrostatic pressure of the rocksalt (B1) phase of ZnO. Synchrotron-radiation x-ray powder-diffraction data are obtained from 0 to 30 GPa. Static simulations of the ZnO B1 phase are performed using the \textit{ab initio} perturbed ion method and the local and nonlocal approximations to the density-functional theory. After the pressure induced transition from the wurtzite phase, we have found that a large fraction of the B1 high-pressure phase is retained when pressure is released. The metastability of this ZnO polymorph is confirmed through the theoretical evaluation of the Hessian eigenvalues of a nine-parameter potential energy surface. This allows us to treat the experimental and theoretical pressure-volume data on an equal basis. In both cases, we have obtained values of the bulk modulus in the range of 160–194 GPa. For its zero-pressure first derivative, the experimental and theoretical data yield a value of 4.4±1.0. Overall, our results show that the ZnO B1 phase is slightly more compressible than previously reported. [S0163-1829/98/07537-7]

I. INTRODUCTION

There is a current interest in the derivation of the parameters that define the equation of state (EOS) of polymorphs that are unstable under room conditions. Basically, these parameters are three zero-pressure magnitudes: the volume \((V_0)\), the bulk modulus \((B_0)\), and its first derivative \((B'_0)\). Attempts to extend standard empirical EOS, such as Birch and Vinet\(^2\) analytical functions, to this situation have been discussed previously for the high-pressure phases of KCl (Ref. 3), ZnS (Refs. 4, 5), and ZnO (Ref. 5). Very recently, Köhler \textit{et al.} have shown that the first-order H11 form of Holzapfel\(^6\) “perfectly represents” the EOS of the high-pressure cesium chloride (B2) phase of the alkali halides.\(^7\) Besides, Hofmeister has also proposed in the last year a strategy to obtain \(V_0\), \(B_0\), and \(B'_0\) from infrared measurements on the B2 phases of the alkali halides.\(^8\)

The EOS of a pressure-induced polymorph can be unambiguously determined in the laboratory when the phase transition is irreversible enough to make the high-pressure phase metastable at ambient conditions. In other cases, theoretical calculations can be a useful alternative approach. For instance, theoretical analysis can be used to select the best set of structural parameters among a series of extrapolated values to zero pressure from high-pressure data.\(^5\) Also, a theoretical scheme can give, after calibration to high-pressure polymorphs, reliable EOS parameters, as documented in the case of the orthorhombic phase of CaF\(_2\), where the results of \textit{ab initio} perturbed ion model (aiPI) were compared with synchrotron-radiation x-ray-diffraction data.\(^9\)

An important example where this structural characterization needs further investigation is ZnO. This crystal undergoes a pressure-induced phase transition from the wurtzite (B4) to the rocksalt (B1) phase in the range of 10 GPa, as reported by Bates, White, and Roy.\(^10\) These authors found that the B1 phase can be quenched, since no trend to revert to the B4 form was found even above 100 °C. However, using \textit{in situ} x-ray diffraction, Jamieson\(^11\) found this transition to be reversible at room temperature. More recently, Karzel \textit{et al.} reported a large hysteresis for this transition from x-ray and \(^{67}\)Zn-Mössbauer spectroscopy studies, although they did not give data for the B1 phase below 2.0±0.5 GPa.\(^12\) In a preliminary study,\(^13\) we have found that a large fraction of the B1 phase is retained at room temperature and ambient pressure.

The few quantum-mechanical studies directed to calculate \(V_0\), \(B_0\), and \(B'_0\) of the ZnO B1 phase leave a number of questions unsolved. Our previous \textit{aiPI} calculations\(^5\) gave a value for \(B_0\) rather different from that reported by Jaffe and Hess\(^14\) using the CRYSTAL approach.\(^15\) These two Hartree-Fock calculations did not include electron correlation, a correction that can increase the bulk modulus by more than 20%.\(^16\) In addition, Karzel \textit{et al.} complemented their experimental work with linearized-augmented plane-wave calculations in the ZnO B1 phase.\(^12\) We note, however, that they did not provide a full determination of the ZnO EOS since the value of \(V_0\) was not reported.

In this paper, we pursue a twofold objective which addresses the question of the stability of the B1 phase of ZnO at zero pressure conditions along with the accurate determination of its EOS parameters. To undertake such investigation, we have combined experimental measurements with quantum-mechanical computations. The experimental data are collected through synchrotron-radiation x-ray diffraction
at pressures obtained in a diamond-anvil cell. The computations are performed using the \( a\!i\!P\!I \) methodology including correlation energy corrections,\(^{17,18}\) and the local and nonlocal density-functional theory (DFT) approximations as implemented in the CRYSTAL program.\(^{15}\) Our experimental observations confirm that the \( B_4\!-\!B_1 \) transition takes place at 10 GPa and the high-pressure phase appears to be metastable at ambient conditions. Our calculations reveal the existence of a local minimum at the \( B_1 \) configuration in a nine-parameter energy surface. From our observed data and quantum-mechanical calculations, we report the best values of \( V_0 \), \( B_0 \), and \( B_0' \) for the EOS of the \( B_1 \) phase of ZnO.

The experimental procedure and the computational details are presented in the next section. Section III contains the results and the discussion on the metastability and the EOS. Our concluding remarks are summarized in the last section.

II. EXPERIMENTAL PROCEDURE AND COMPUTATIONAL TECHNIQUES

A. Pressure-volume measurements

X-ray powder-diffraction spectra were recorded using synchrotron radiation and the white-beam, energy-dispersive method. High pressures were obtained in a Syassen-Holzapfel type\(^{19}\) diamond-anvil cell. The sample and a small ruby chip were enclosed in a hole of diameter 0.1 mm in an inconel gasket. The incident-beam cross section was 60 \( \times \) 60 \( \mu \)m\(^2\). A 4:1 methanol:ethanol mixture was used as the pressure-transmitting medium, and the pressure was determined from the wavelength shift of the ruby \( R \) line, applying the nonlinear pressure scale of Mao et al.\(^{20}\) The uncertainty in the pressure determination is estimated to be 0.1 GPa for pressures below about 10 GPa. For higher pressures the uncertainty may be larger because of nonhydrostatic conditions as will be discussed below. The diffraction data were obtained at the UK synchrotron radiation source of the Daresbury Laboratory, and at the synchrotron radiation laboratory HASYLAB in Hamburg, Germany. All experiments were performed at room temperature. The diffraction peaks in the recorded spectra yield information on the corresponding plane spacings, and the unit-cell volume is deduced from the refined lattice constants.

B. Total energy calculations

We have followed two different theoretical approaches to compute the total energy of the ZnO \( B_1 \) phase. The first one is the \( a\!i\!P\!I \) model, a quantum-mechanical method that solves the Hartree-Fock (HF) equations of the solid in a localized Fock space by breaking the crystal wave function into local, group functions (atomic or ionic in nature).\(^{21}\) For a detailed description of the computational implementation and its latest update, we submit to Refs. 17 and 18. In the present work, we have used for Zn and O\(^{2-} \) the nearly HF multi-\( \xi \) exponential Clementi and Roetti basis sets\(^{22} \) of Zn\(^{2+} \) and O\(^{-} \). After each \( a\!i\!P\!I \) calculation, a correlation energy correction is added to the total energy of the crystal using the Coulomb-Hartree-Fock method of Clementi.\(^{23}\) These calculations are denoted as \( u\!C\!H\!F\!-\!P\!I \).

As a second theoretical tool, we have used the implementation of the DFT methodology in a periodic linear-combination-of-atomic-orbitals (LCAO) scheme included in the crystal program.\(^{15}\) In the LCAO approach, the Bloch functions are constructed as linear combinations of atom-centered Gaussian orbitals. We have considered the Perdew-Zunger parametrization\(^{24} \) of Ceperley-Alder results\(^{25} \) for the local functional calculations (referred to as LDA), and the combination of the Becke exchange functional\(^{26} \) with the Perdew-Wang correlation functional\(^{27} \) for the nonlocal case (referred to as GGA). In this work, we have selected the same basis sets for Zn and O that Jaffe and Hess used in their HF calculations of ZnO polymorphs.\(^{14} \) This fact allows a direct comparison of their results with the local and nonlocal computations performed here.
III. RESULTS AND DISCUSSION

A. Thermodynamic stability of the ZnO rocksalt phase

In our x-ray-diffraction experiments, it is observed that the hexagonal wurtzite structure of ZnO undergoes a structural phase transformation with a transition pressure $p_{tr} \approx 10$ GPa. It is around this pressure where the diffraction lines of the high-pressure rocksalt phase begin to appear. The transition is completed at about 15 GPa. The transition pressure is in good agreement with the experimental value 9.5 GPa tabulated by Phillips,28 and slightly greater than the recent value $8.7^{\pm 0.5}$ GPa obtained also in x-ray-diffraction experiments by Karzel et al.12 With respect to our theoretical predictions, we have obtained different values of the thermodynamic $p_{tr}$ depending on the computational method. It is known that under the consideration of spherical electronic densities, as in the aPI model, the $B_1$ phase of ZnO is the structure with lower energy at zero pressure (see Ref. 29). Incorporation of nonspherical deformations through GGA calculations yield a value for the thermodynamic $p_{tr}$ around 9 GPa, close to the experimental range. As expected, the LDA prediction is lower by about 5 GPa.16

The measured lattice-plane spacings as functions of pressure for the $B_1$ phase are shown in Fig. 1. Accordingly, a large fraction of the $B_1$ phase is retained when the pressure is released. In this way, the $B_1$ phase of ZnO is found to be metastable at zero pressure, and $V_0$ for this structure can be determined. The lattice parameter is $a_0 = 4.275 \pm 0.004$ Å in good agreement with the value 4.280 Å reported by Bates et al.10 and the value $4.271 \pm 0.002$ Å given by Karzel et al.12 However, in this reference the ambient pressure data are extrapolated from higher pressures, since below $\approx 2$ GPa their $B_1$ phase was completely transformed back to the wurtzite structure.

In the theoretical study of the $B_1$ phase, we have first computed the total energy of this structure at volumes ranging approximately from 0.6 to 1.2 times the experimental equilibrium value.10 The results for the uCHF-PI and GGA calculations are plotted in Fig. 2. The calculated equilibrium volumes yield the static lattice parameter of 4.058 Å at the uCHF-PI and 4.316 Å at the GGA. If we compare the present uCHF-PI with the HF-PI calculation previously published,5 we see that the stabilization effect of the correlation energy correction is greater at lower volumes leading to a decreasing of the lattice parameter with respect to the HF (PI) result, as expected. The uCHF-PI lattice parameter is now 5% smaller than the experimental value. The use of nonlocal functionals, on the other hand, produces also the expected upwards shift in the equilibrium lattice spacing (see Table I).16 The GGA equilibrium geometry is about 1% larger than the experimental value.

Once the equilibrium volume is located, we explore the response of the $B_1$ structure to general strains in two different ways. First, we seek for small perturbations from the equilibrium restricted to a primitive rhombohedral unit cell containing one ZnO molecule: the cation is at the origin of the cell, the anion at $(x, x, x)$, and the space group is $R3m$. Therefore, we allow the crystal to search lower energy situations in a three-dimensional (3D) configuration space, where the three degrees of freedom are the unit-cell length and angle ($a$ and $\alpha$), and the internal coordinate $x$. In this 3D model, the $B_1$ phase corresponds to the point for which

### Table I. Summary of EOS parameters for the $B_1$ phase of ZnO.

<table>
<thead>
<tr>
<th>Method</th>
<th>$a_0$ (Å)</th>
<th>$B_0$ (GPa)</th>
<th>$B_0'$</th>
</tr>
</thead>
<tbody>
<tr>
<td>Ref. 10 X-ray</td>
<td>4.280</td>
<td>228</td>
<td>4 (fix)</td>
</tr>
<tr>
<td>Ref. 12 X-ray</td>
<td>4.271</td>
<td>132 $^b$</td>
<td>3.8 $^b$</td>
</tr>
<tr>
<td>Ref. 5 HF-PI</td>
<td>4.225</td>
<td>203.3</td>
<td>3.6</td>
</tr>
<tr>
<td>Ref. 14 HF-CRYSTAL</td>
<td>4.294</td>
<td>205</td>
<td>4.88</td>
</tr>
<tr>
<td>This work X-ray</td>
<td>4.275</td>
<td>194 $^a$-$^b$</td>
<td>4.8 $^a$-$^b$</td>
</tr>
<tr>
<td>This work uCHF-PI</td>
<td>4.058</td>
<td>170 $^a$-$^b$</td>
<td>3.37 $^a$-$^b$-$^c$</td>
</tr>
<tr>
<td>This work GGA</td>
<td>4.316</td>
<td>176 $^a$-$^b$</td>
<td>5.39 $^a$-$^b$-$^c$</td>
</tr>
<tr>
<td>This work LDA</td>
<td>4.207</td>
<td>229 $^a$-$^b$</td>
<td>4.24 $^a$-$^b$-$^c$</td>
</tr>
</tbody>
</table>

$^a$Birch EOS fitting.
$^b$H11 EOS fitting.
$^c$Vinet EOS fitting.
In Fig. 3, we plot the energy curves for the displacement of these three parameters from their B1 values. It is found that all departures from the equilibrium are accompanied by a rising in the crystal energy, and so the B1 phase is a stable minimum with respect to these three deformations independently. It is worth noting that GGA calculations predict softer curves for α and x displacements than uCHF-PI. This effect is related to the neglect of multipolar terms in aiPI, which are stabilizing contributions for the lower symmetry situations out of the minimum. Nevertheless, both calculations predict that the softer curves are those arising from the variation of α. We are currently working in the description of a possible transition mechanism to the B1 structure from the B3 (zinc-blende) phase. This form can be described within the same 3D space (x = 1/4, α = 60°), and its structure closely resembles that of the ambient pressure polymorph, the B4 (wurtzite) phase.

We have also completed a more general and detailed investigation of the stability of the B1 phase by relaxing the symmetry constraints of the rhombohedral cell. We have considered the case of a ≠ b ≠ c, α ≠ β ≠ γ and x ≠ y ≠ z. This nine parameter (9D) model is analogous to the one used in our previous study of the microscopic description of the B1-B2 phase transition.\textsuperscript{30} In order to prove that the B1 phase of ZnO is stable, and thus a minimum in this 9D space, we determine and diagonalize its 9 × 9 Hessian matrix at the equilibrium volume. Given the high symmetry of the B1 phase, this matrix shows a blocked structure.

\[
H^{B1} = \begin{pmatrix}
H_{aa} & H_{ab} & H_{ab} & H_{aa} & H_{a\beta} & H_{a\beta} & 0 & 0 & 0 \\
H_{ab} & H_{aa} & H_{ab} & H_{a\beta} & H_{aa} & H_{a\beta} & 0 & 0 & 0 \\
H_{ab} & H_{ab} & H_{aa} & H_{a\beta} & H_{a\beta} & H_{aa} & 0 & 0 & 0 \\
H_{aa} & H_{ab} & H_{aa} & H_{a\beta} & H_{a\beta} & H_{a\beta} & 0 & 0 & 0 \\
H_{a\beta} & H_{aa} & H_{a\beta} & H_{aa} & H_{a\beta} & H_{a\beta} & 0 & 0 & 0 \\
H_{a\beta} & H_{a\beta} & H_{a\beta} & H_{aa} & H_{a\beta} & H_{a\beta} & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & H_{sx} & H_{xy} & H_{xy} \\
0 & 0 & 0 & 0 & 0 & 0 & H_{xy} & H_{xx} & H_{xy} \\
0 & 0 & 0 & 0 & 0 & 0 & H_{xy} & H_{xy} & H_{xx}
\end{pmatrix}
\]

Table I. Around 11 GPa, the methanol:ethanol mixture is known to freeze, and nonhydrostatic conditions may be introduced in the diamond-anvil cell.\textsuperscript{31} Fitting the EOS to all our experimental data in the whole observed pressure range \(p < 30\) GPa results in \(B'_0 = 6.6\), i.e., a value that is larger than expected. We believe that the freezing of the pressure transmitting medium cause the observed slope change in our experimental pressure-volume data. The previous experimental \(B_0\) and \(B'_0\) values yield a description of a less compressible ZnO B1 phase. Thus, the x-ray data by Karzel \textit{et al.}\textsuperscript{12} give a \(B_0\) value which is about 35 GPa above our experimental value. Their slightly lower value for \(B'_0\) can only explain partially this large disagreement. We should notice two facts here. First, Karzel \textit{et al.} EOS parameters come from an extrapolation to zero pressure of data in the pressure range 2–10 GPa. Second, their \(B_0\) value for the ZnO B4 phase is also ≈40 GPa greater than previously reported.\textsuperscript{32,33}

We have also generated \(p-V\) tables from the computed crystal energy (\(E_{\text{cryst}}\)) values by minimizing the static Gibbs energy (\(G = E_{\text{cryst}} + pV\)) with respect to \(V\) at selected values of the hydrostatic pressure in the range 0–30 GPa. The \(p-V\) tables were then used to generate \(B_0\) and \(B'_0\) by means, again, of Birch, Vinet, and H11 fittings. The values are included in Table I. We have obtained in the three cases very good fittings to both uCHF-PI and GGA data. The quality of

\[x = 1/2, \, \alpha = 60^\circ, \text{ and } a = a_0 / \sqrt{2}, \text{ being } a_0 \text{ the equilibrium lattice parameter previously computed.}\]
the fittings is illustrated in Fig. 2, where we plot the computed uCHF-PI and GGA $E_{\text{crv}}$ values along with the curves obtained after the integration of the corresponding Vinet EOS.

To represent the data from our experiments and computations, we have adopted the usual $V/V_0$ versus $p$ diagram, $V_0$ being the zero-pressure volume of the B1 phase (see Fig. 4). In this normalized representation, both uCHF-PI and GGA curves show very good agreement with the experimental data in the pressure region below the transition pressure (0–10 GPa). Above this value, it is found a slow progressive departure of the theory with respect to the experiment, reaching a 2% difference at the highest pressure measured (30 GPa). Due to this fact, the ZnO B1 phase is described as more compressible in our computations. This picture is consistent with the values obtained for $B_0$ and $B'_0$. In our work, experimental and theoretical $B_0$ values are in the range 160–194 GPa, and the slope change in the experimental data is due to the freezing of the pressure transmitting medium as discussed above.

The comparison with other calculations shows a better agreement than with the experiments of Karzel et al. although in general our computations yield a slightly more compressible ZnO B1 phase than previously predicted. The difference between $B_0$ values computed with and without correlation effects shows an opposite trend to the difference in the prediction of $a_0$, as mentioned previously. The same can be said for the results under DFT when comparing local and nonlocal functionals. We notice that the improvements in the theoretical approaches (i.e., inclusion of correlation energy corrections and use of nonlocal density functionals) change the predictions of $a_0$ and $B_0$ in a different direction. Thus, the uCHF-PI $B_0$ value is about 40 GPa higher than the HF-PI one, and the corresponding GGA value is lower than the LDA one by approximately 50 GPa. Finally, it is to be stressed that given the different formalisms upon which the two quantum-mechanical methodologies are based, the computed EOS parameters show an excellent agreement with each other and globally describe satisfactorily our x-ray-diffraction data.

IV. CONCLUDING REMARKS

In this paper, we have reported the results of experiments and calculations directed (i) to investigate the thermodynamic stability of the high-pressure B1 phase of ZnO, and (ii) to determine the EOS parameters of this polymorph. The synchrotron-radiation x-ray diffraction informs us of a pressure-induced transition from the B4 to the B1 phase around 10 GPa. The B1 phase is retained metastable at ambient conditions. Our quantum-mechanical calculations have confirmed the existence of a local minimum located at the B1 position on a nine-dimensional energy surface.

The experimental and calculated results have been used to obtain the $p$-$V$ behavior of the ZnO B1 phase. $B_0$ and $B'_0$ values are generated using the standard Birch and Vinet empirical EOS along with the semiempirical H11 EOS. The most interesting conclusions are: (i) our four different sets of data (one experimental and three theoretical) can be consistently represented by these EOS giving in each case very similar Birch, Vinet, and H11 EOS parameters; (ii) $B_0$ is obtained in the range 160–194 GPa in the experiments and in the theoretically improved uCHF-PI and GGA calculations; (iii) freezing of the pressure transmitting medium causes a change in slope of the experimental pressure-volume curve around 11 GPa that explains the differences with respect to the computed results; and (iv) the global comparison with previous experiments and other calculations shows that the ZnO B1 phase is slightly more compressible than reported before.

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